

Superconducting circuits I

Demonstration of conditional gate operation using superconducting charge qubits

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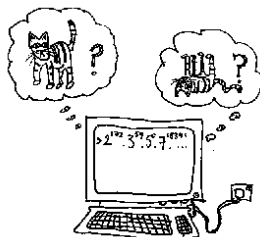
Demonstration of controlled-NOT quantum gates on a pair of superconducting quantum bits

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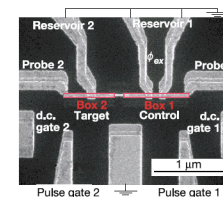
Outline

- Motivation

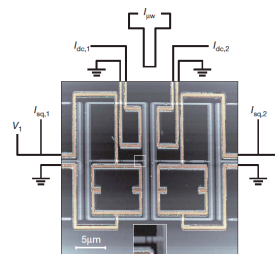


- C-NOT gate

- First realization of C-NOT gate with CPBs



- C-Not gate with flux qubits



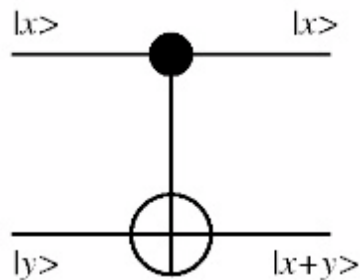
- Comparison and Summary

Motivation

- Approaching the goal of quantum computation
- Fulfilling DiVincenzo criteria
- Superconducting quantum bits as building blocks for a quantum computer
- High fidelity gate operation

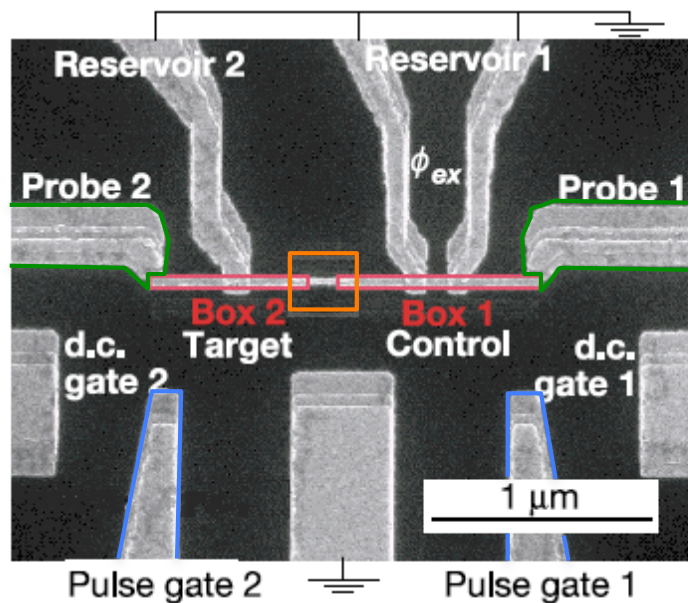
c-NOT gate

- Single qubit operation & c-NOT gate form a universal set of gates (\rightarrow *any computation can be done using these gates*)
- Definition:
„The target qubit is flipped if and only if the control qubit is in a given state“



$$\text{CNOT} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}.$$

The sample



Requirements:

- 2 two-level systems
 - capacitive coupling between qubits
- four-level system
- gates for independent control
 - probes for read-out
 - control bit tuned with B-field

Hamiltonian describing the system:

$$H = \sum_{n_1, n_2=0,1} E_{n_1 n_2} |n_1, n_2\rangle \langle n_1, n_2| - \frac{E_{J1}}{2} \sum_{n_2=0,1} (|0\rangle\langle 1| + |1\rangle\langle 0|) \otimes |n_2\rangle \langle n_2| - \frac{E_{J2}}{2} \sum_{n_1=0,1} |n_1\rangle \langle n_1| \otimes (|0\rangle\langle 1| + |1\rangle\langle 0|)$$

electrostatics

Josephson coupling of box 1

Josephson coupling of box 2

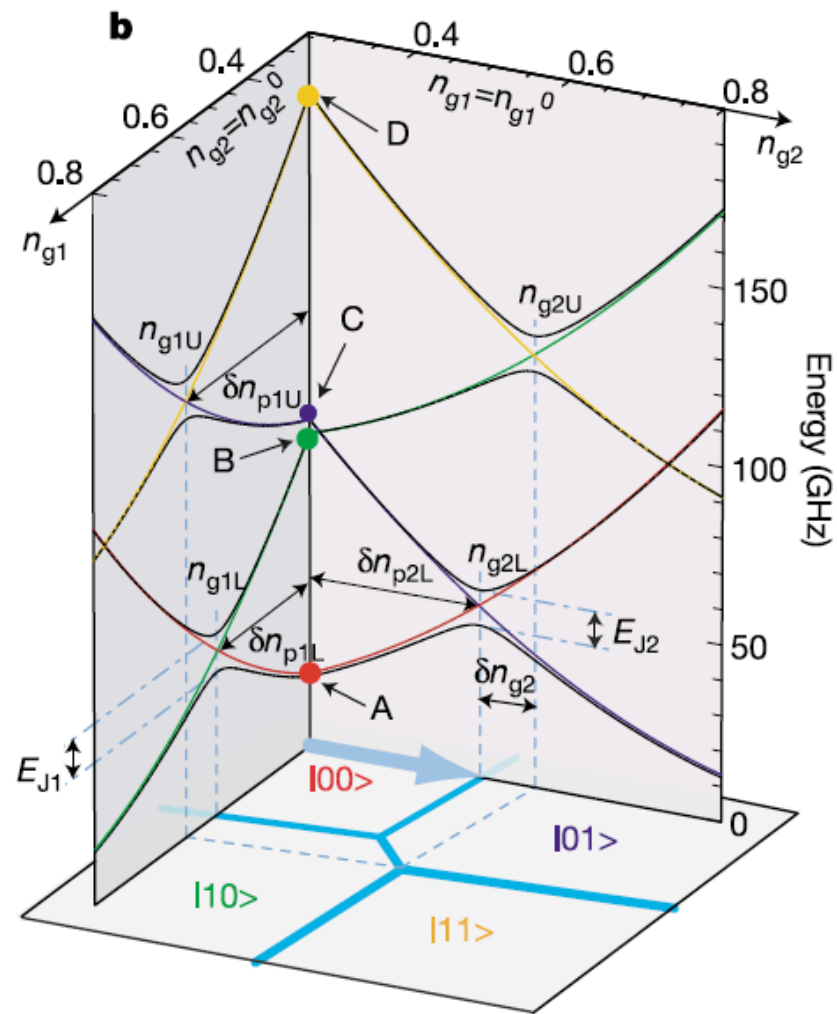
Energy band diagram

Eigenenergies of system at constant values for n_{g1} and n_{g2}

(number of excess Cooper pairs on respective box)

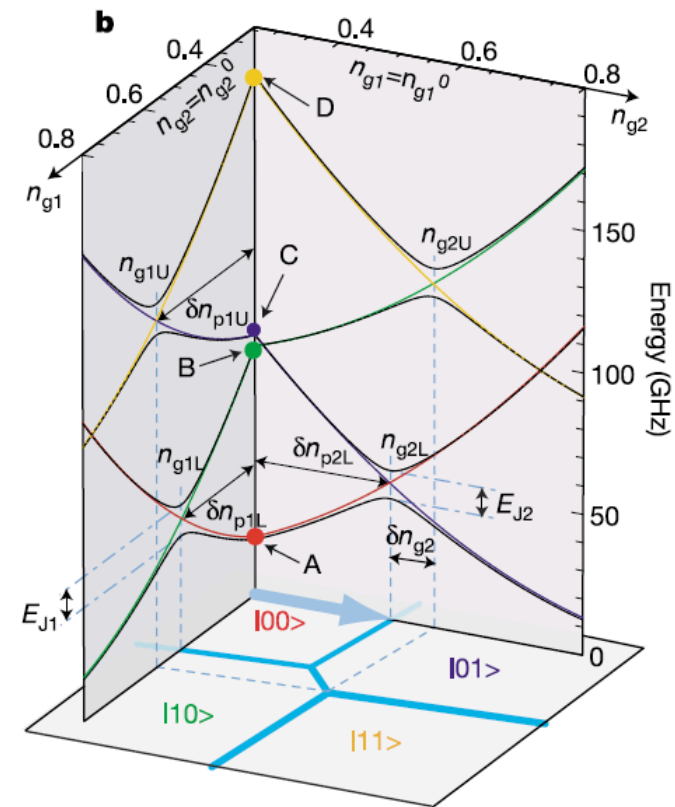
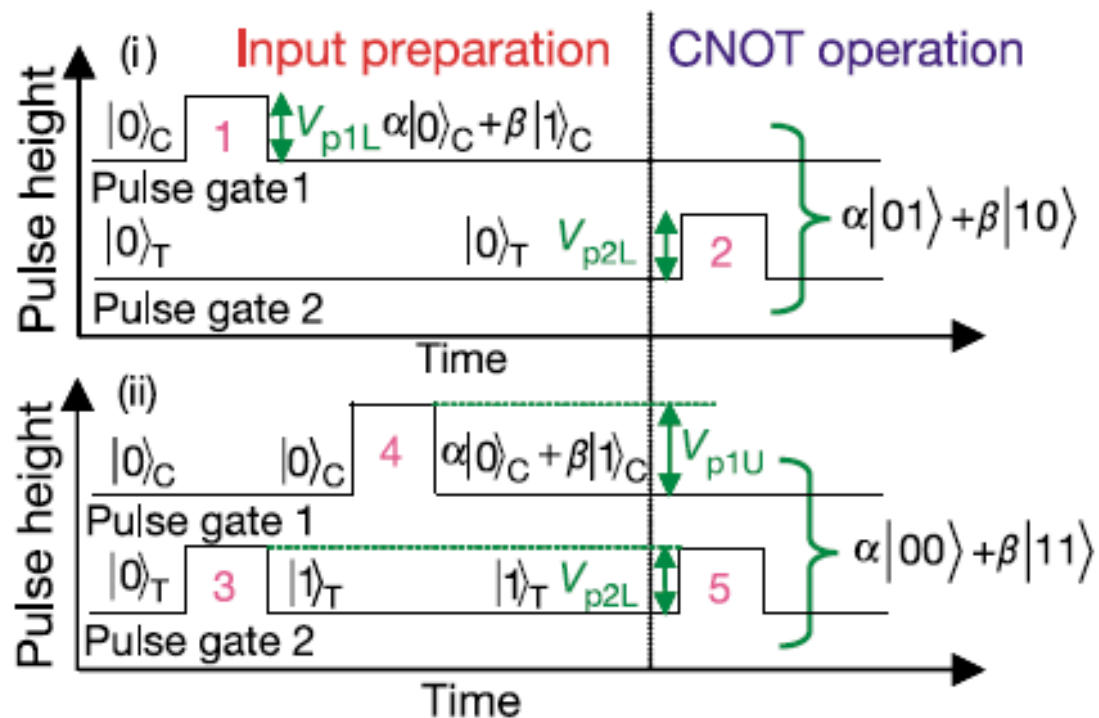
Determining possible outcomes of different pulse schemes

→ Controlled gate operation can be implemented



Pulse scheme and C-NOT operation

1. Preparation of specific input state
2. Applying c-NOT operation



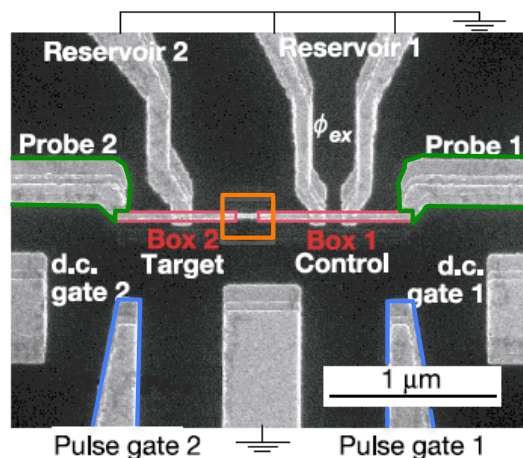
→ **Creation of entangled states**

Tunability and read-out

Periodic modulation of E_{J1} due to SQUID-geometry: $E_{J1} = E_{J1max} \left| \cos \left(\pi \frac{\phi_{ex}}{\phi_0} \right) \right|$

Recording JQP current through probe 1 and 2 (I is proportional to n_g)

→ read-out of state

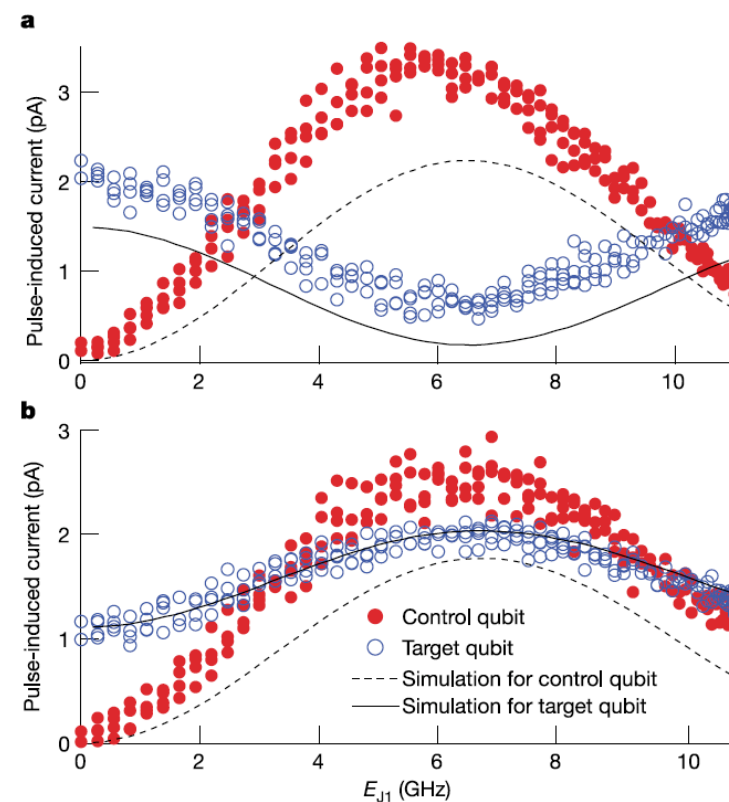


$$\alpha |01\rangle + \beta |10\rangle$$

pulse scheme (i)

$$\alpha |00\rangle + \beta |11\rangle$$

pulse scheme (ii)



Simulation: Time evolution of density matrix

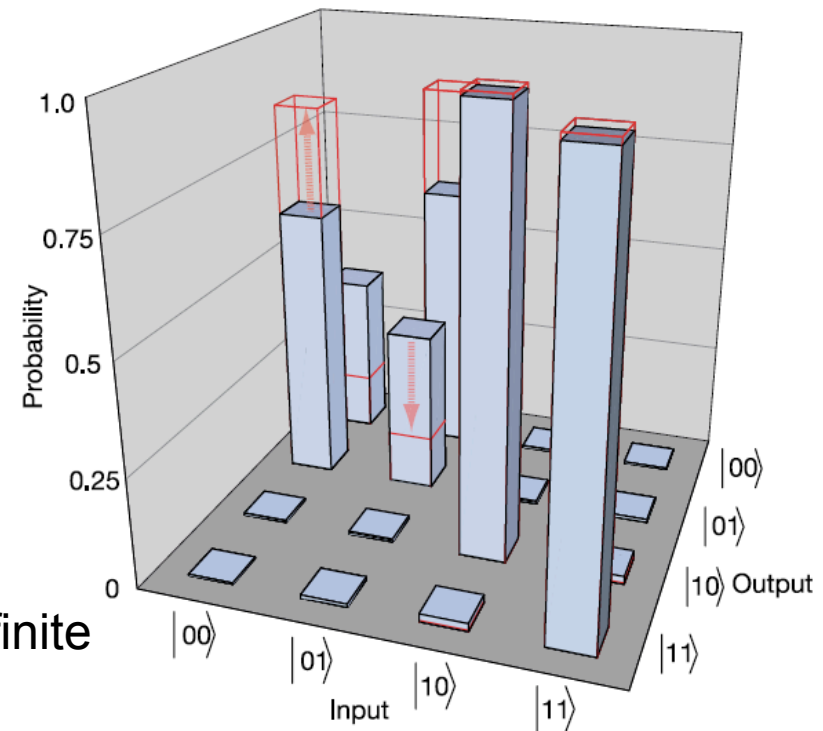
Truth table

Read-out method does not allow for individual measurement of the four states

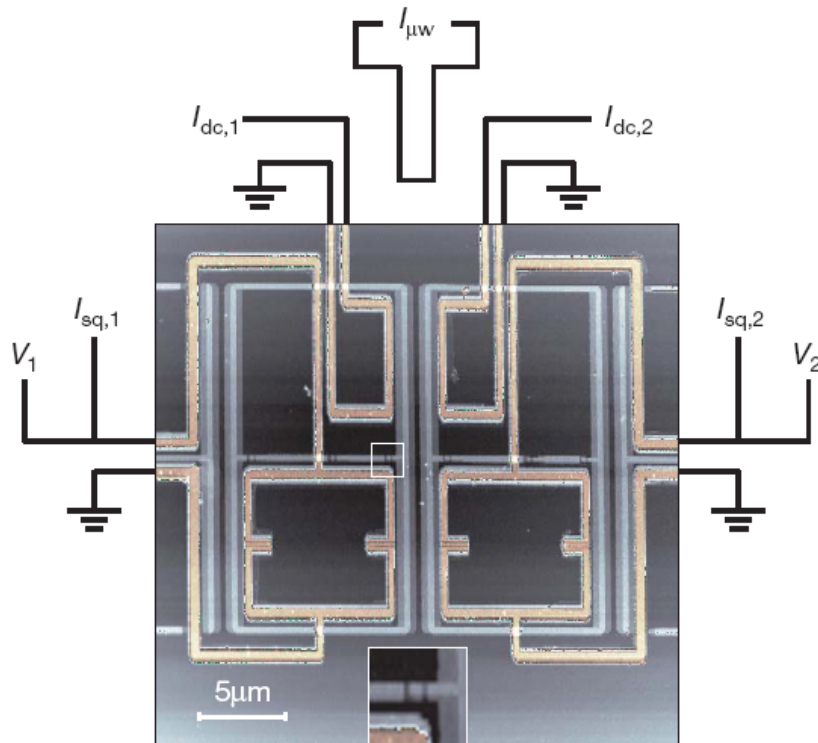
→ Calculation of time evolution of four perfect input states under gate operation pulse

Deviation from expected accuracy due to finite rise/fall time of pulse

→ Improving pulse shape and the coupling E_m



Coupled qubits set-up



Requirements:

- single pair of coupled flux qubits
- Inductive coupling between qubits
- Either qubit can be control or target qubit due to symmetry
- 4 level system

Set-up:

- two '8' shape flux qubits consisting of a superconducting loop interrupted by 3 Josephson junctions
- Two SQUIDs used as switching quantum state detectors

Hamiltonian describing the system:

$$H = H_1 + H_2 + H_{12} = -\frac{1}{2} (\epsilon_1 \sigma_z^1 + \Delta_1 \sigma_x^1 + \epsilon_2 \sigma_z^2 + \Delta_2 \sigma_x^2) + J \sigma_z^1 \sigma_z^2$$

Operation of the coupled-qubit device

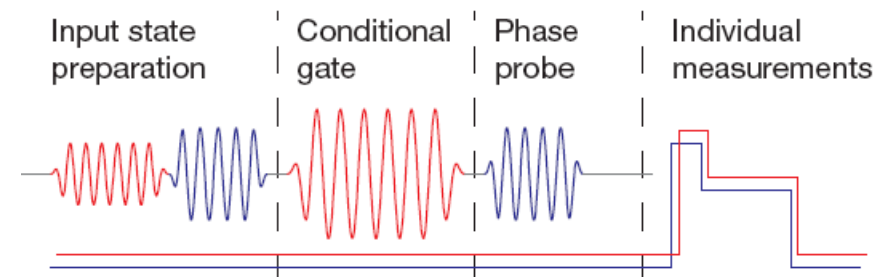
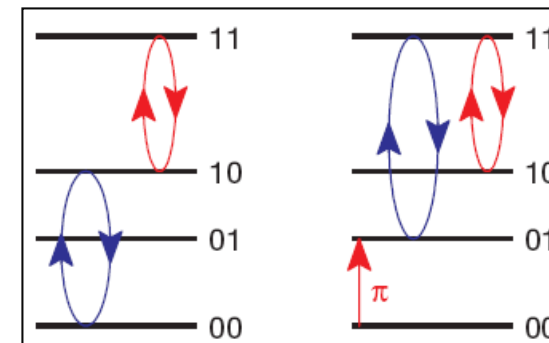
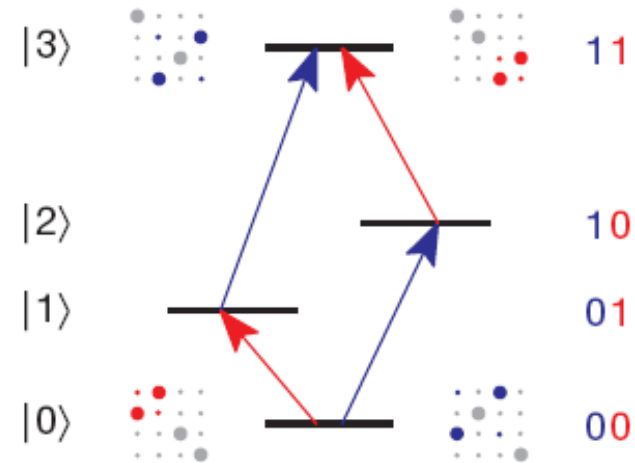
Energy level diagram

- four resonance frequencies
- A resonant microwave pulse induces rotations in the computational basis

$$|0_C 0_T\rangle, |0_C 1_T\rangle, |1_C 0_T\rangle, |1_C 1_T\rangle$$

Sequence of operations

- Initial ground state $|0_C 0_T\rangle$
- Preparation of input states
 1. $\frac{1}{\sqrt{2}}(|00\rangle + |10\rangle)$
 2. $\frac{1}{\sqrt{2}}(|01\rangle + |11\rangle)$
- Gate operation by applying a pulse
- Analysis of the resulting density matrix with probe pulses
- Simultaneous and independent determination of the two qubit state
- Repetition of N times \Rightarrow state counts
 $N_{00}, N_{01}, N_{10}, N_{11}$



Tunability and read-out

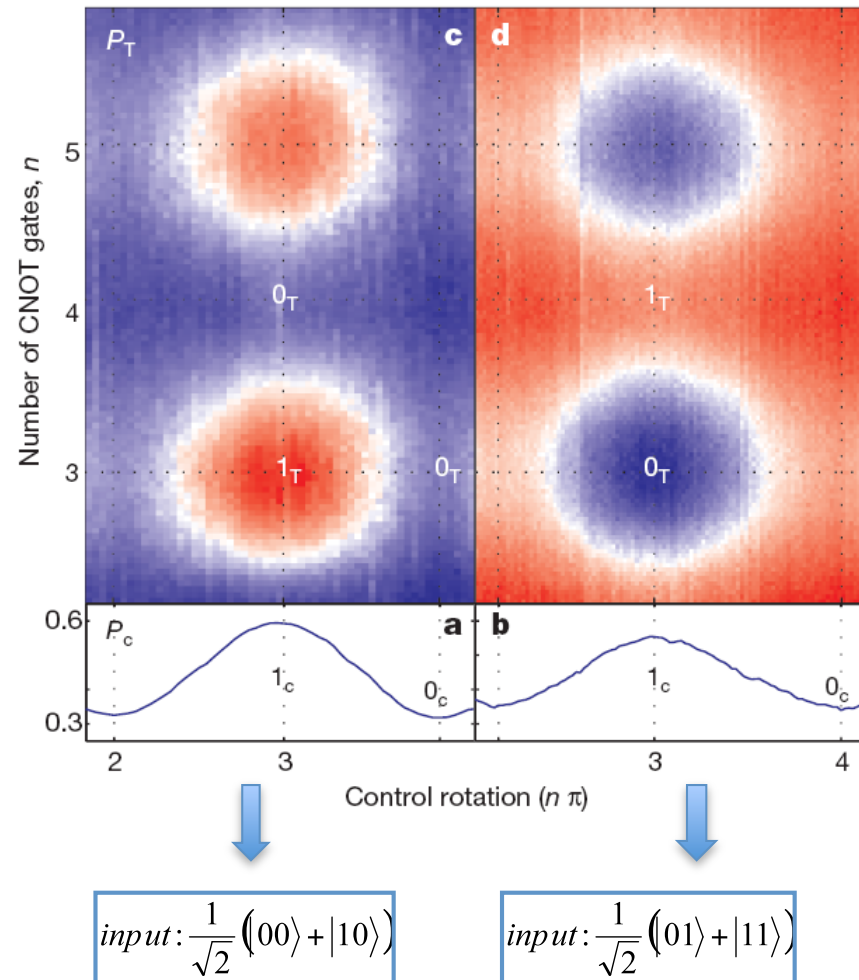
- Measured joint probabilities

$$P_{00}, P_{01}, P_{10}, P_{11}$$

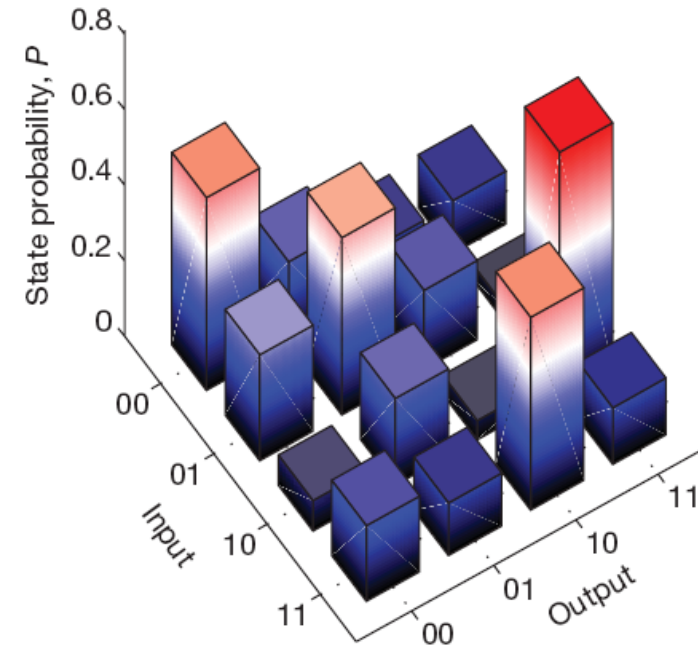
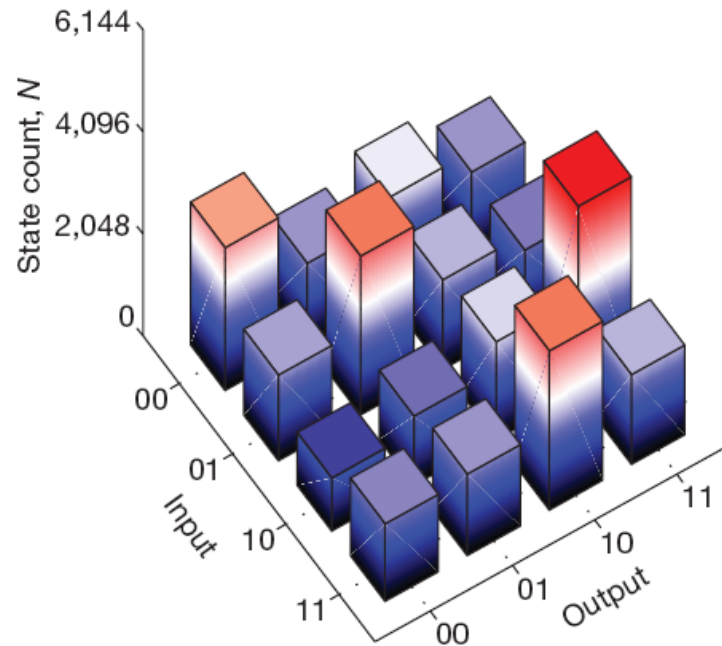
$$P_C = P_{10} + P_{11}$$

$$P_T = P_{01} + P_{11}$$

- Odd numbers of π rotations and C-NOT gates flips the target qubit
- It is a 1c-controlled gate



Truth table-Correction



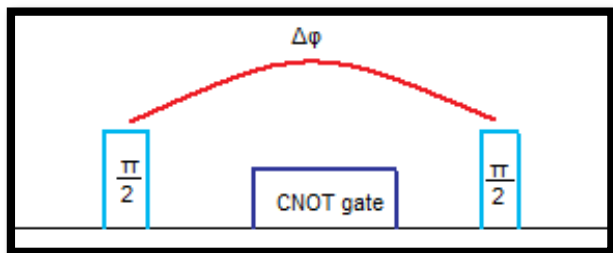
- Input states $0_C 0_T, 0_C 1_T \Rightarrow$ remain unaffected
 $1_C 0_T, 1_C 1_T \Rightarrow$ target qubit inverted


Corrected truth table

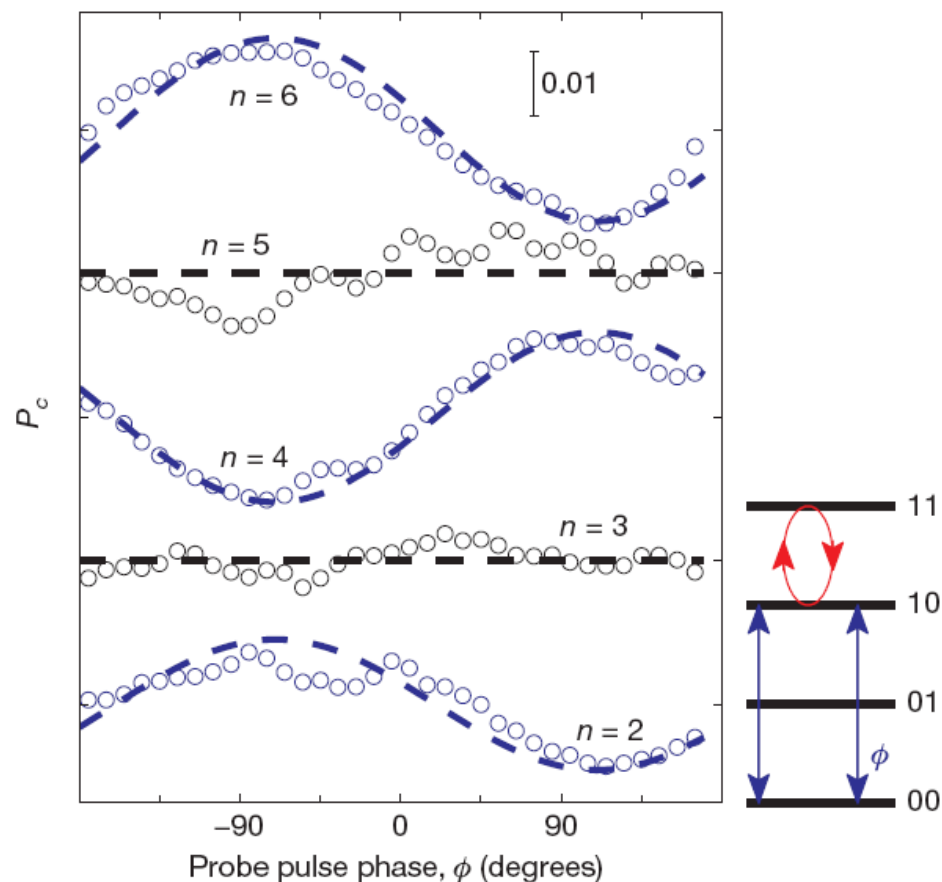
- Correction with conditional spectroscopy measurements
- New $F=0.4$

Phase Factor

- Ramsey-like interference experiment on n consecutive CNOT gates
- Starting at superspositions instead of starting at eigenstates
- Additional $\pi/2$ pulse after the gate with phase difference $\Delta\phi$ to the one before the gate



- Even number of gates
 Phase gate



Conclusions

⊕ Two different implementations of CNOT gate

⊕ Main differences:

- | | |
|-------------------------|------------------------|
| • Charge qubits | • Flux qubits |
| • Capacitive coupling | • Inductive coupling |
| • Zero-controlled gate | • One-controlled gate |
| • Simulated truth table | • Measured truth table |
| • Phase unknown | • Phase determination |

➡ Two qubit algorithms and solid-state qubit entanglement is possible

Summary-Outlook

- Superconducting qubits are among the most promising candidates for quantum computation
- Obstacles to overcome:
 - ➡ increasing decoherence time
 - ➡ improvement of read-out fidelity
 - ➡ implementing error correction methods