

Experimental implementation of Grover's algorithm with transmon qubit architecture

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What is Grover's algorithm?

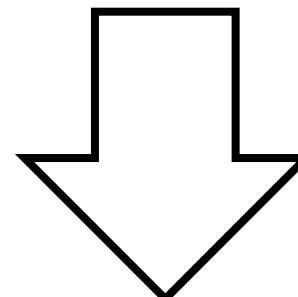
- Quantum search algorithm
- Task: In a search space of dimension N , find those $0 < M < N$ elements displaying some given characteristics (being in some given states).

Classical search (random guess)	Grover's algorithm
<ul style="list-style-type: none">• Guess randomly the solution• Control whether the guess is actually a solution	<ul style="list-style-type: none">• Apply an ORACLE, which <i>marks</i> the solution• Decode the marked solution, in order to <i>recognize</i> it
$O(N)$ steps $O(N)$ bits needed	$O(\sqrt{N}) \times n$ steps $O(\log(N))$ qubits needed

The oracle

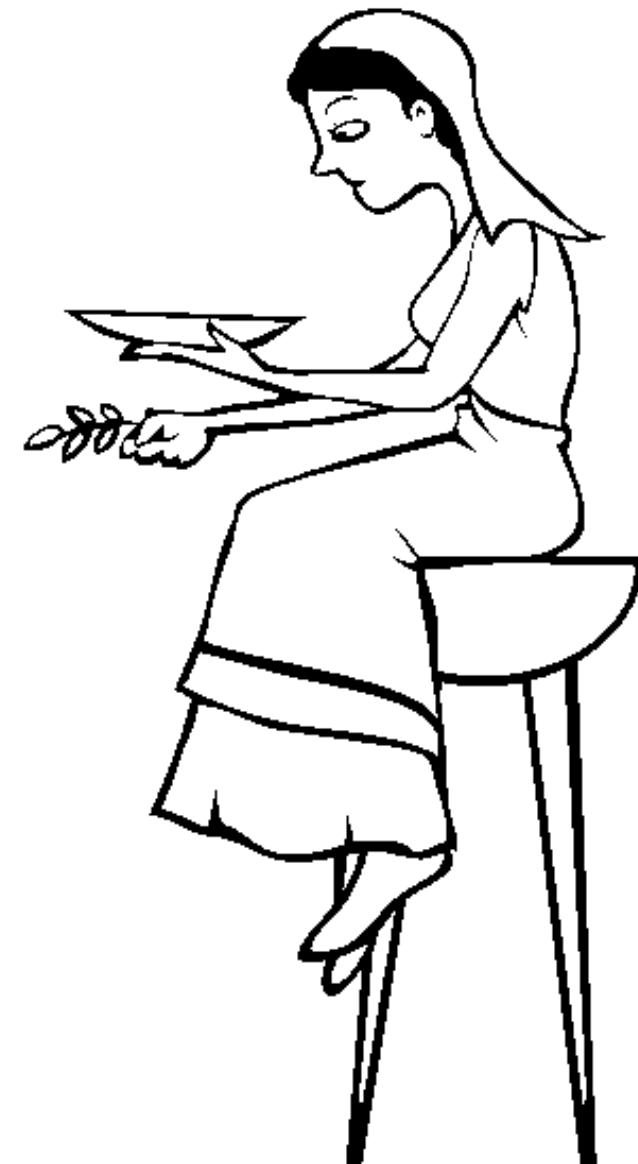
- The oracle MARKS the correct solution

$$f(x) = \begin{cases} 0 & x \text{ is not solution} \\ 1 & x \text{ is solution} \end{cases}$$



Dilution operator
(interpreter)

- Solution is more recognizable



Grover's algorithm

Procedure

- Preparation of the state

$$|\psi\rangle = \frac{1}{\sqrt{N}} \sum_x^N |x\rangle$$

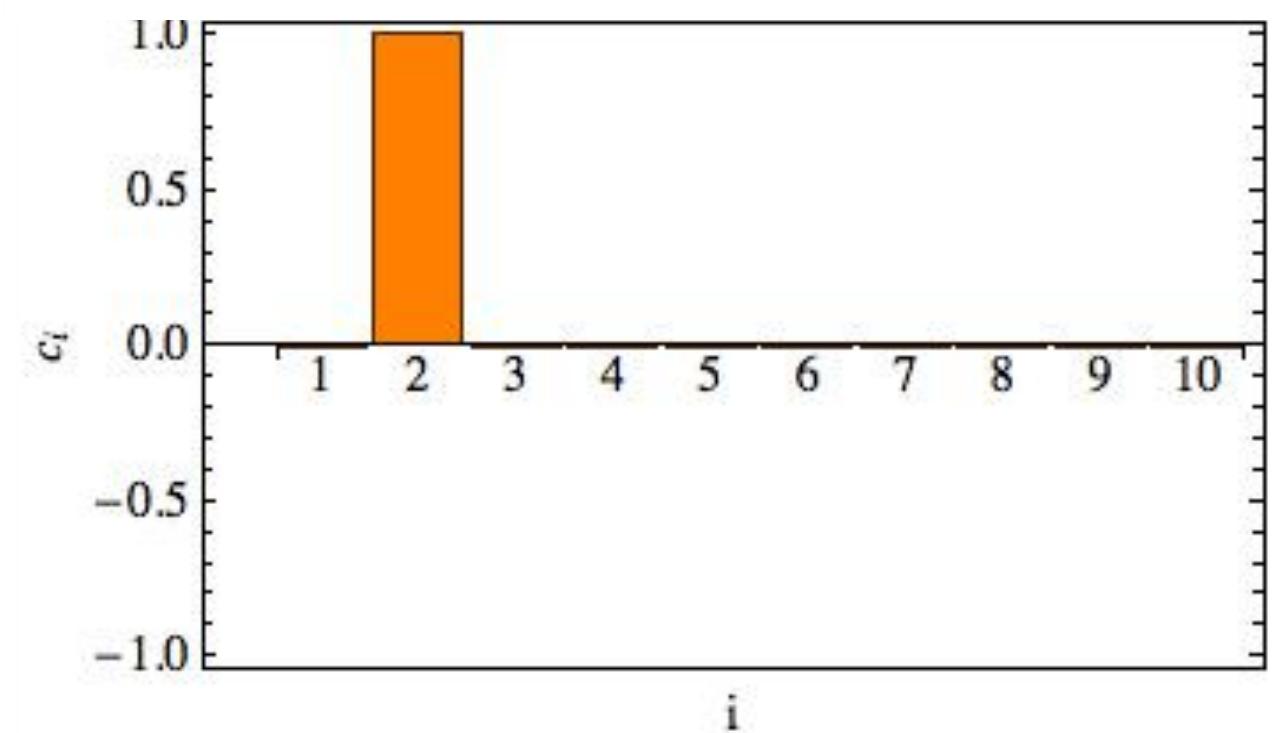
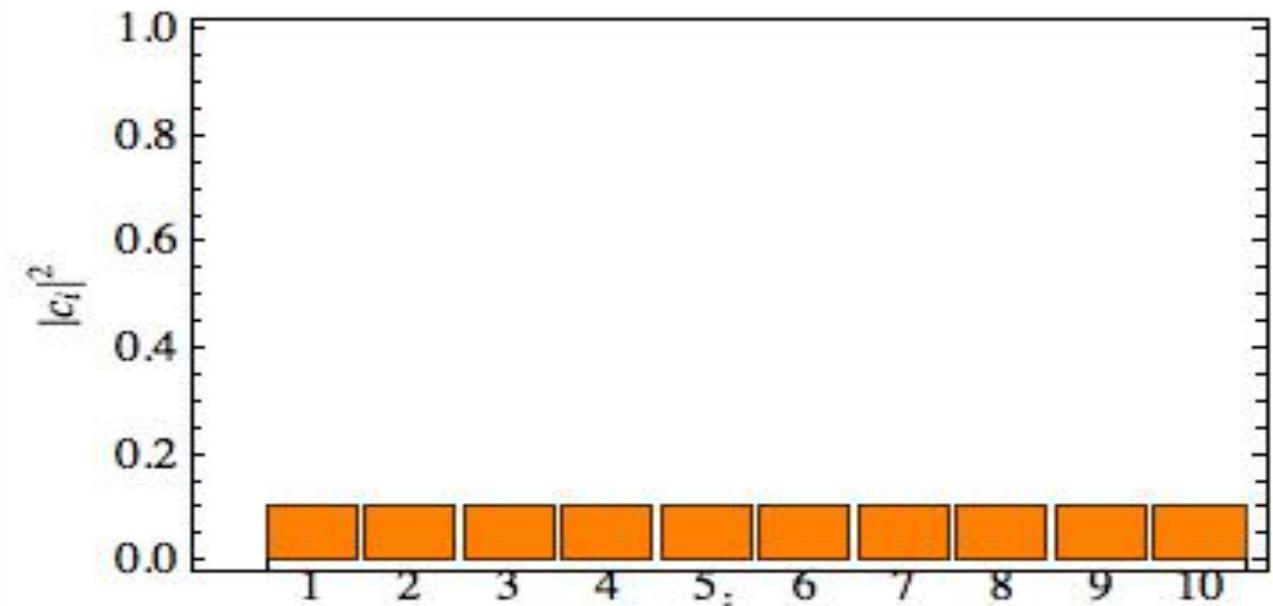
- Oracle application

$$|x\rangle \xrightarrow{O} (-1)^{f(x)} |x\rangle$$

- Dilution of the solution

$$2|\psi\rangle\langle\psi| - I$$

- Readout



Geometric visualization

- Preparation of the state

$$|\psi\rangle = \frac{1}{\sqrt{N}} \sum_x^N |x\rangle$$

$$|\bar{t}\rangle = \frac{1}{\sqrt{N-M}} \sum'_x |x\rangle$$

$$|t\rangle = \frac{1}{\sqrt{M}} \sum''_x |x\rangle G|\psi\rangle$$

$$|\psi\rangle = \sqrt{\frac{N-M}{N}} |\alpha\rangle + \sqrt{\frac{M}{N}} |\beta\rangle$$

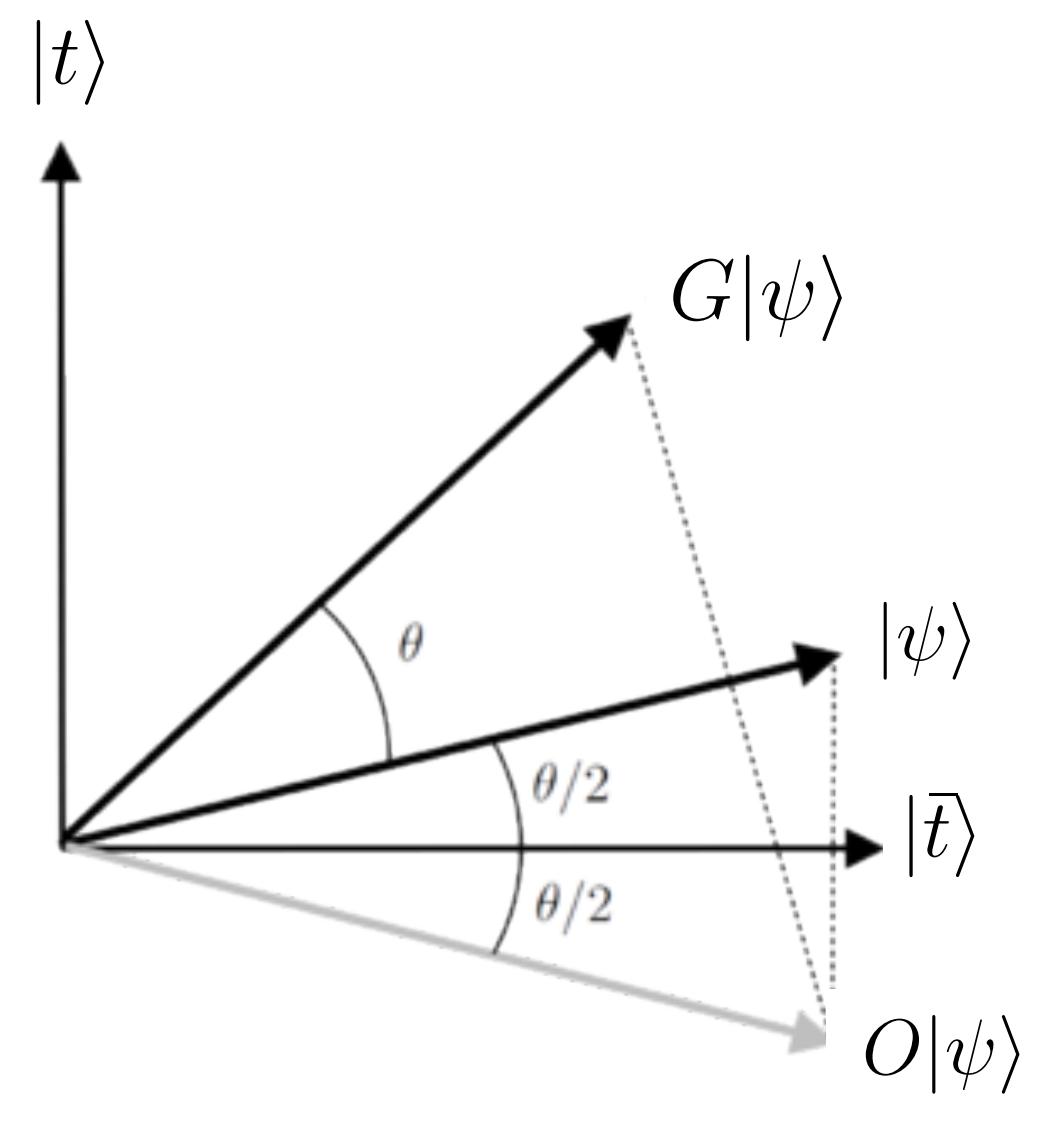
Geometric visualization

$$|\psi\rangle = \cos(\theta/2)|\bar{t}\rangle + \sin(\theta/2)|t\rangle$$

$$O|\psi\rangle = \cos(\theta/2)|\bar{t}\rangle - \sin(\theta/2)|t\rangle$$

$$2|\psi\rangle\langle\psi| - I$$

$$G|\psi\rangle = \cos\left(\frac{3\theta}{2}\right)|\bar{t}\rangle + \sin\left(\frac{3\theta}{2}\right)|t\rangle$$



Grover's algorithm

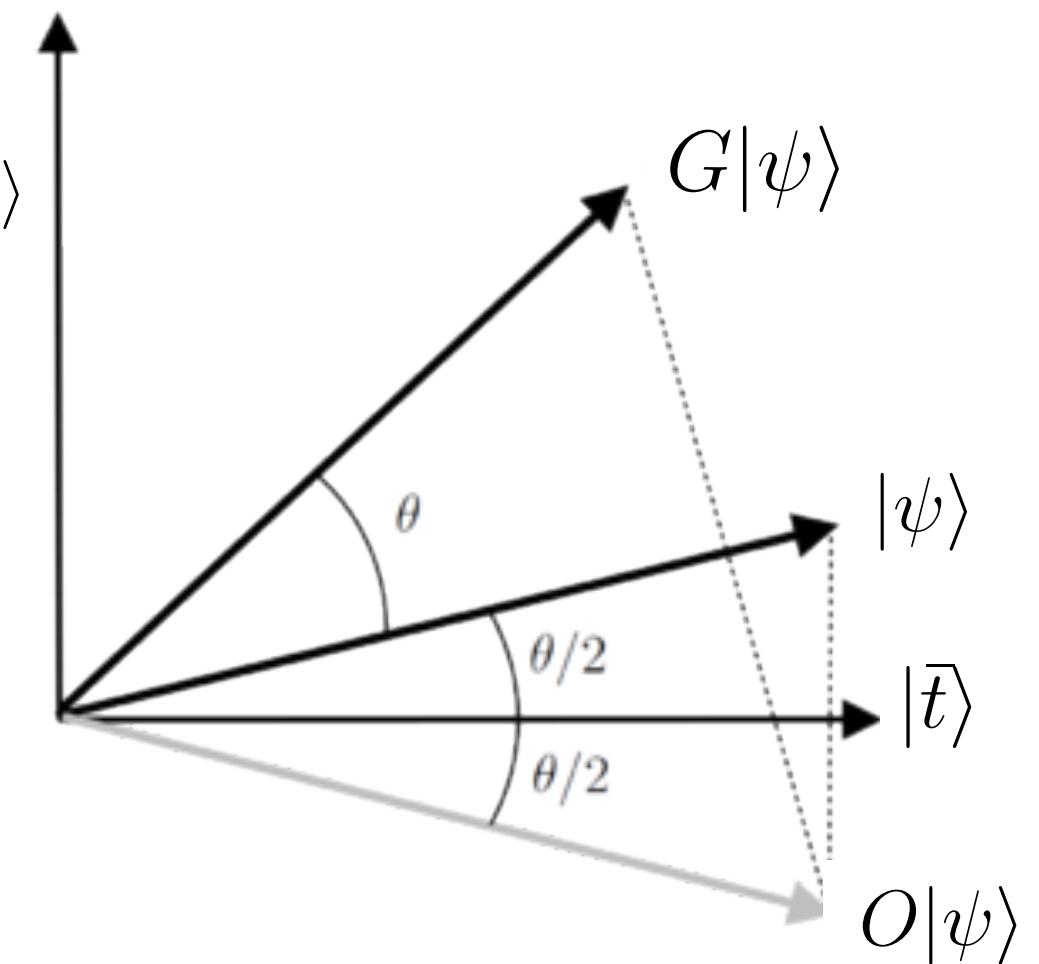
Performance

- Every application of the algorithm is a $|t\rangle$ rotation of θ

$$G^k|\psi\rangle = \cos\left(\frac{(2k+1)\theta}{2}\right)|\bar{t}\rangle + \sin\left(\frac{(2k+1)\theta}{2}\right)|t\rangle$$

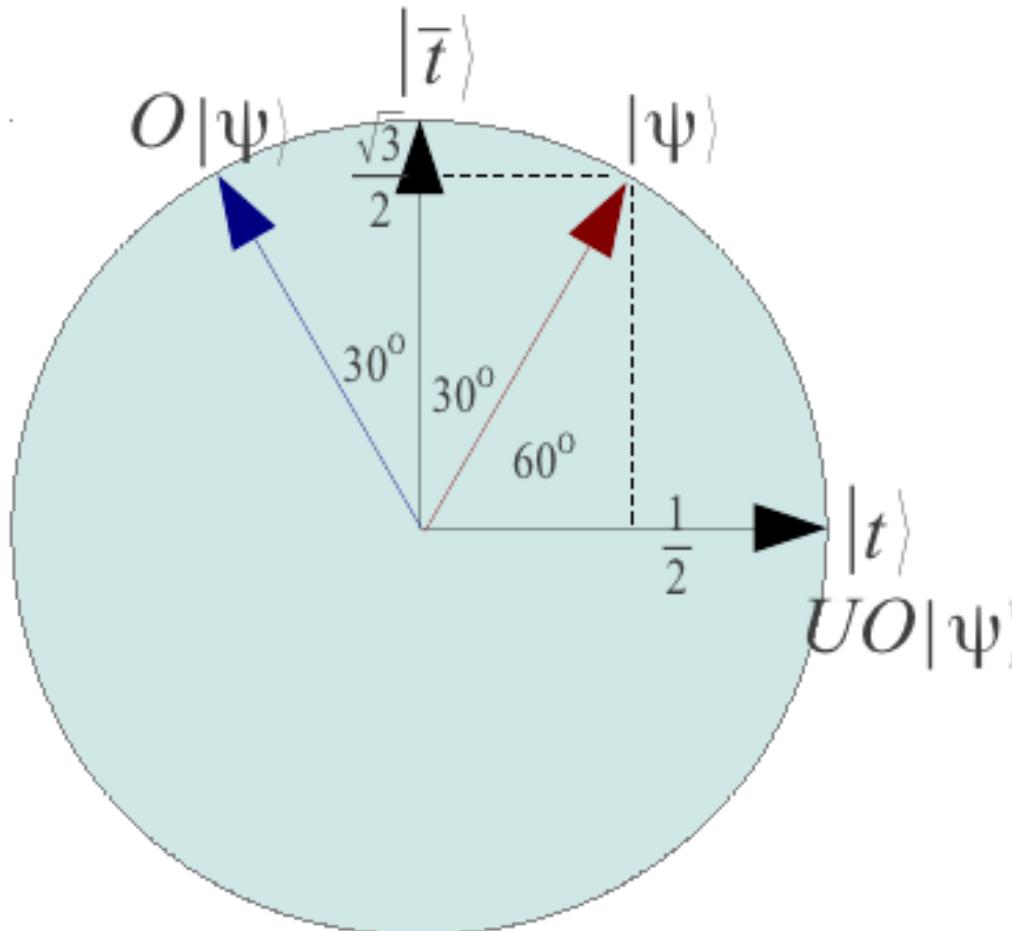
- The Ideal number of rotations is:

$$R = CI \left[\frac{\arccos \sqrt{M/N}}{\theta} \right]$$



Grover's algorithm

2 qubits



N=4

Oracle marks one state M=1

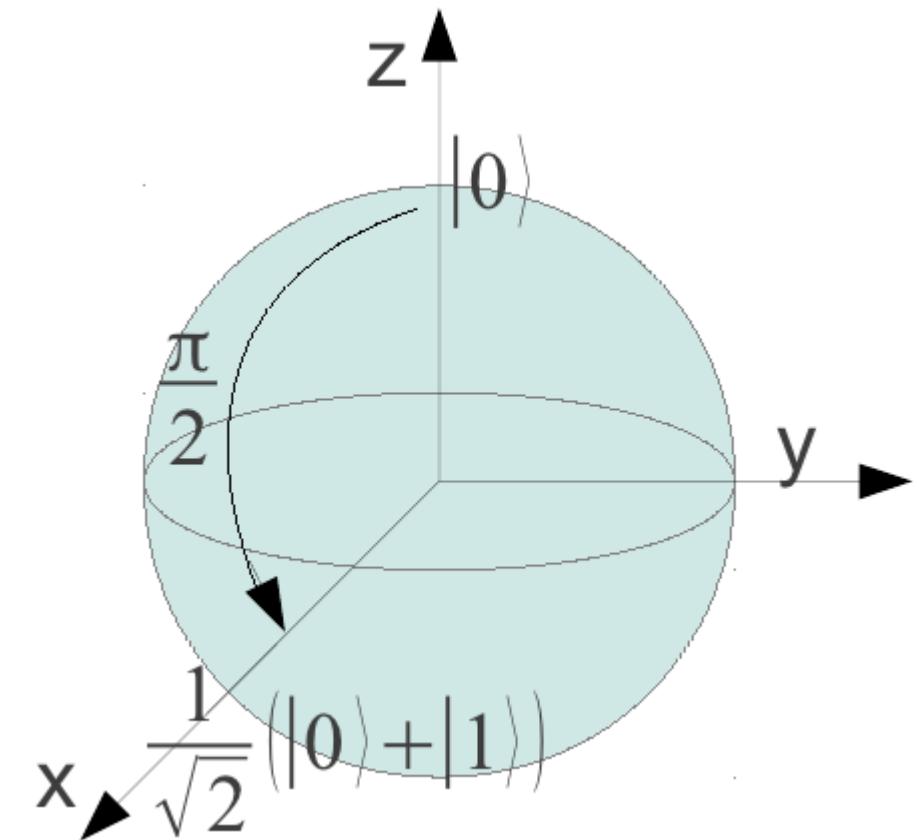
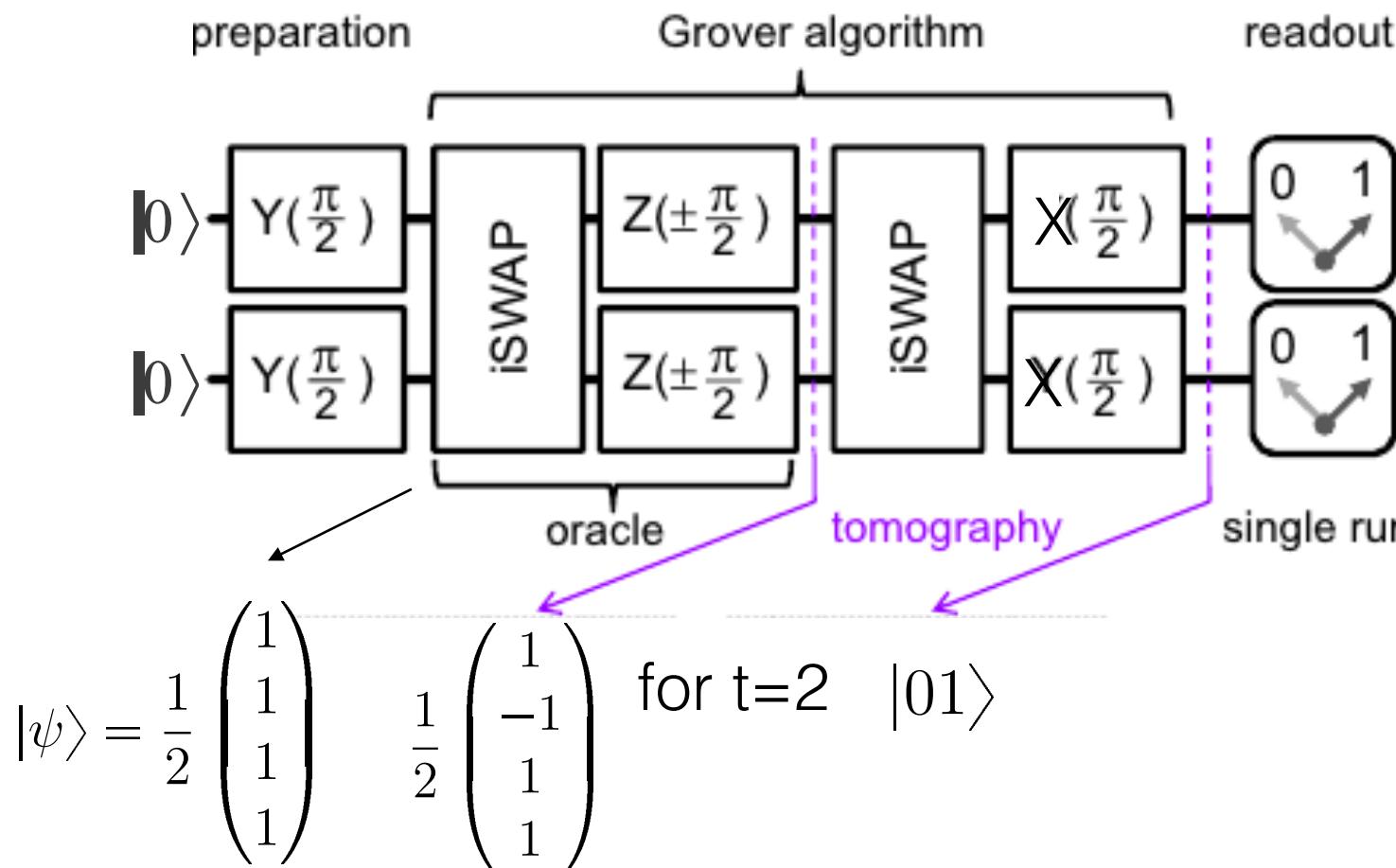
$$\begin{aligned} |\psi\rangle &= \frac{1}{2} (|00\rangle + |01\rangle + |10\rangle + |11\rangle) \\ &= \frac{1}{2}|t\rangle + \frac{\sqrt{3}}{2} \frac{1}{\sqrt{3}} (|\bar{t}_1\rangle + |\bar{t}_2\rangle + |\bar{t}_3\rangle) \\ &= \frac{1}{2}|t\rangle + \frac{\sqrt{3}}{2}|\bar{t}\rangle \end{aligned}$$

After a single run and a projection measurement will get target state with probability 1!

Grover's algorithm

Circuit

Preparation



$$|\psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) = \frac{1}{2} \begin{pmatrix} 1 \\ 1 \\ 1 \\ 1 \end{pmatrix}$$

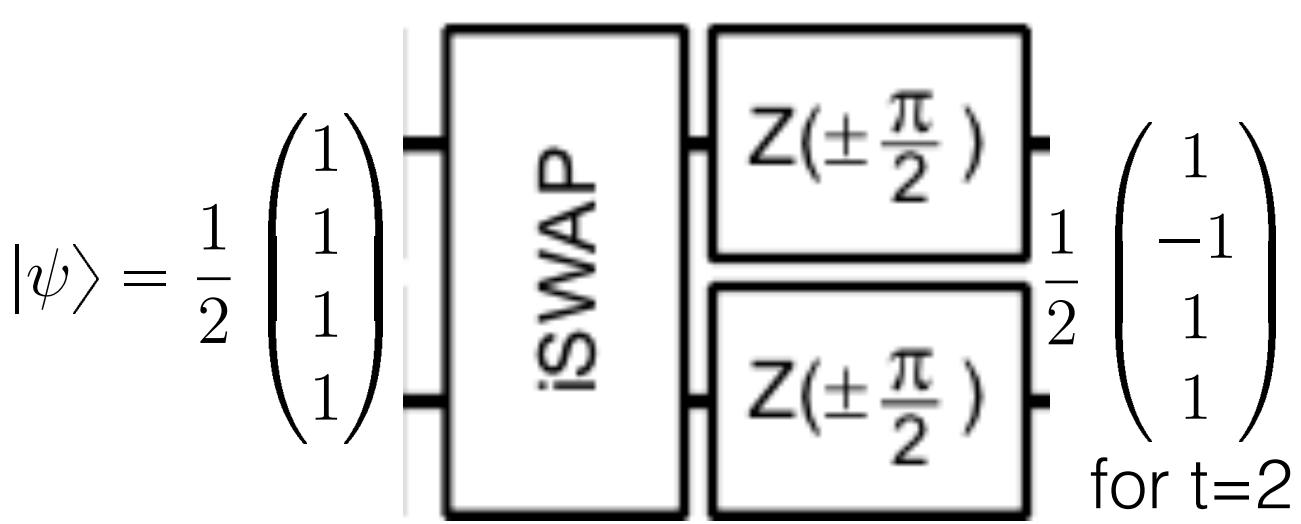
Grover's algorithm

Oracle

Will get 2 cases:

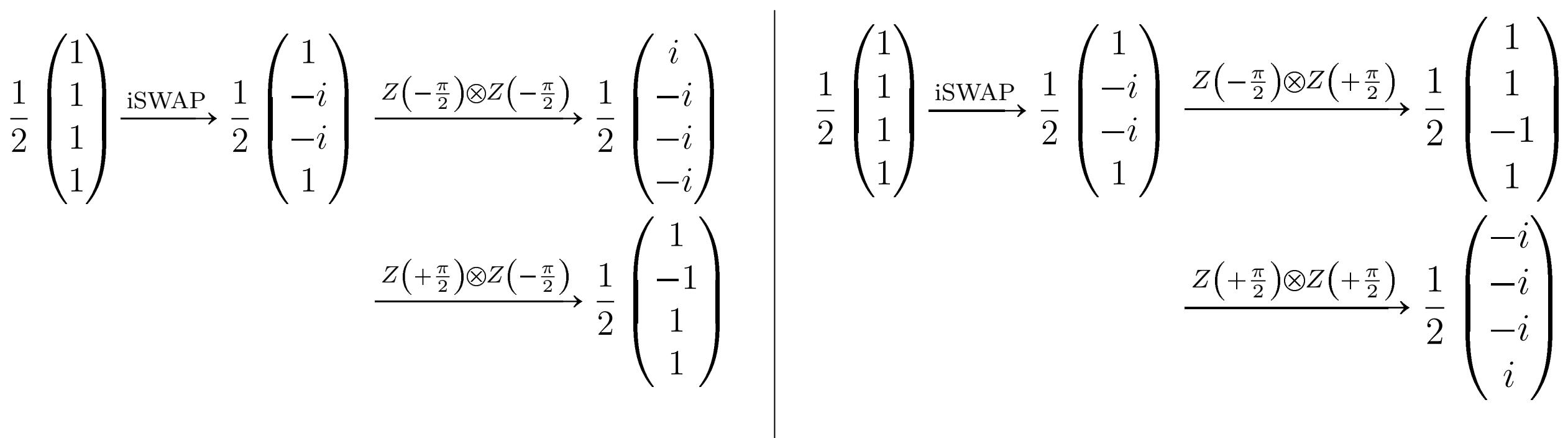
$$e^{\pm i\pi} = \pm i$$

$$e^{\pm i0} = 1$$



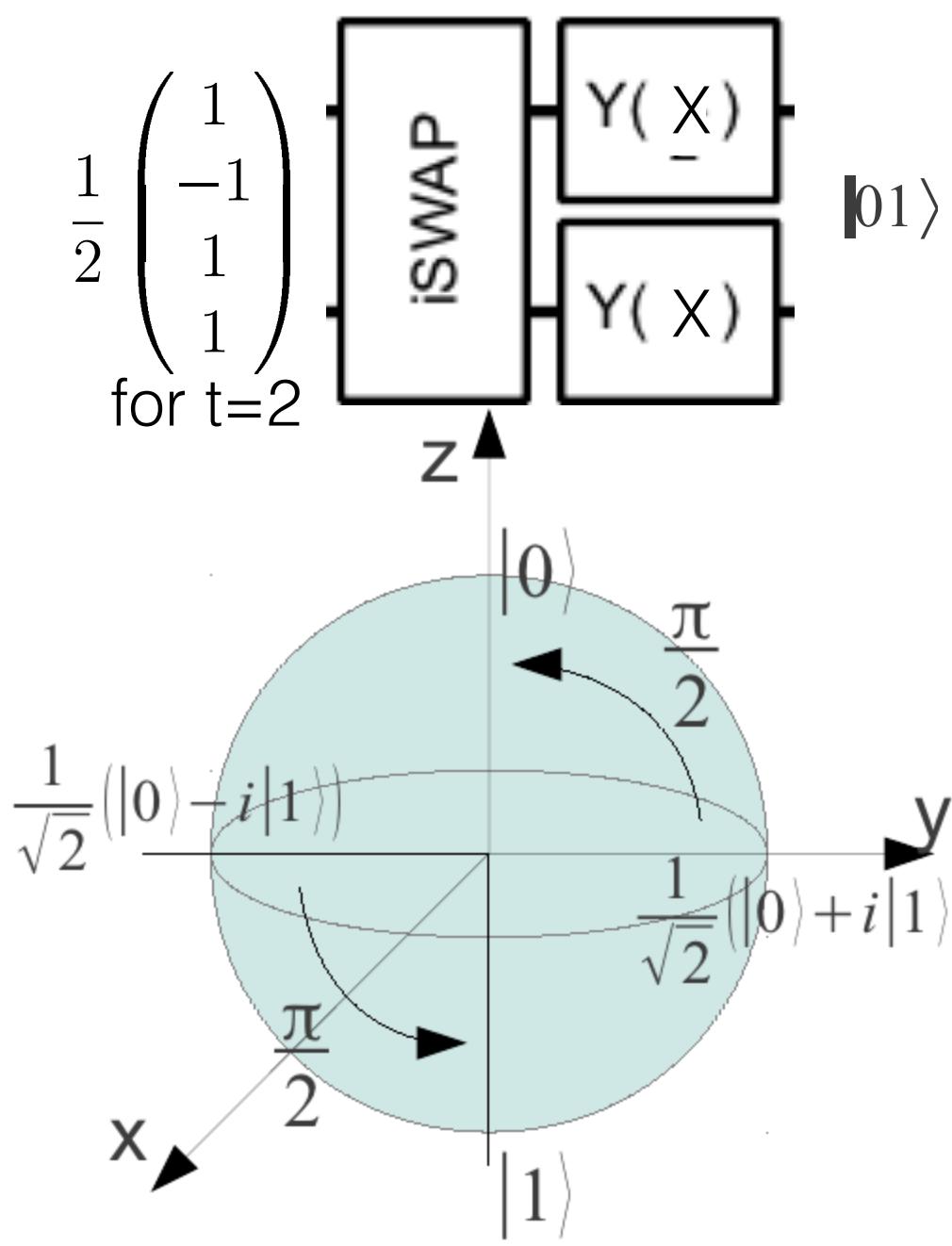
$$Z(\theta) = e^{-i\frac{\theta}{2}Z} = \begin{pmatrix} e^{-i\frac{\theta}{2}} & 0 \\ 0 & e^{i\frac{\theta}{2}} \end{pmatrix}$$

$$Z(\theta) \otimes Z(\phi) = \begin{pmatrix} e^{i\frac{-\theta-\phi}{2}} & 0 & 0 & 0 \\ 0 & e^{i\frac{-\theta+\phi}{2}} & 0 & 0 \\ 0 & 0 & e^{i\frac{\theta-\phi}{2}} & 0 \\ 0 & 0 & 0 & e^{i\frac{\theta+\phi}{2}} \end{pmatrix}$$



Grover's algorithm

Decoding



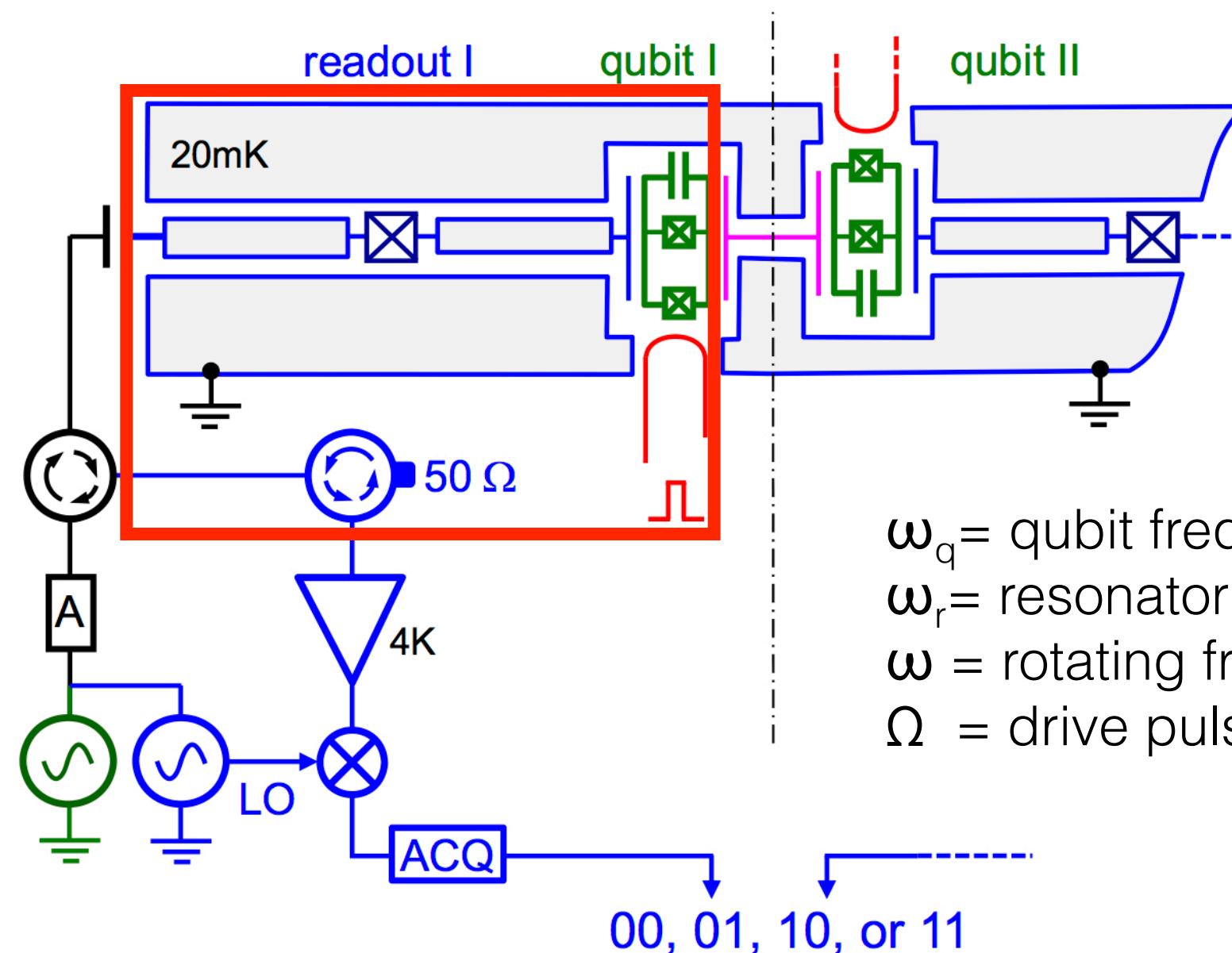
$$\frac{1}{2} \begin{pmatrix} -1 \\ 1 \\ 1 \\ 1 \end{pmatrix} \xrightarrow{\text{iSWAP}} \frac{1}{2} \begin{pmatrix} -1 \\ -i \\ -i \\ 1 \end{pmatrix} = -\frac{1}{\sqrt{2}} (|0\rangle + i|1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + i|1\rangle) = |\circlearrowleft\rangle \otimes |\circlearrowleft\rangle$$

$$\frac{1}{2} \begin{pmatrix} 1 \\ -1 \\ 1 \\ 1 \end{pmatrix} \xrightarrow{\text{iSWAP}} \frac{1}{2} \begin{pmatrix} 1 \\ -i \\ i \\ 1 \end{pmatrix} = |\circlearrowleft\rangle \otimes |\circlearrowright\rangle$$

$$\frac{1}{2} \begin{pmatrix} 1 \\ 1 \\ -1 \\ 1 \end{pmatrix} \xrightarrow{\text{iSWAP}} \frac{1}{2} \begin{pmatrix} 1 \\ i \\ -i \\ 1 \end{pmatrix} = |\circlearrowright\rangle \otimes |\circlearrowleft\rangle$$

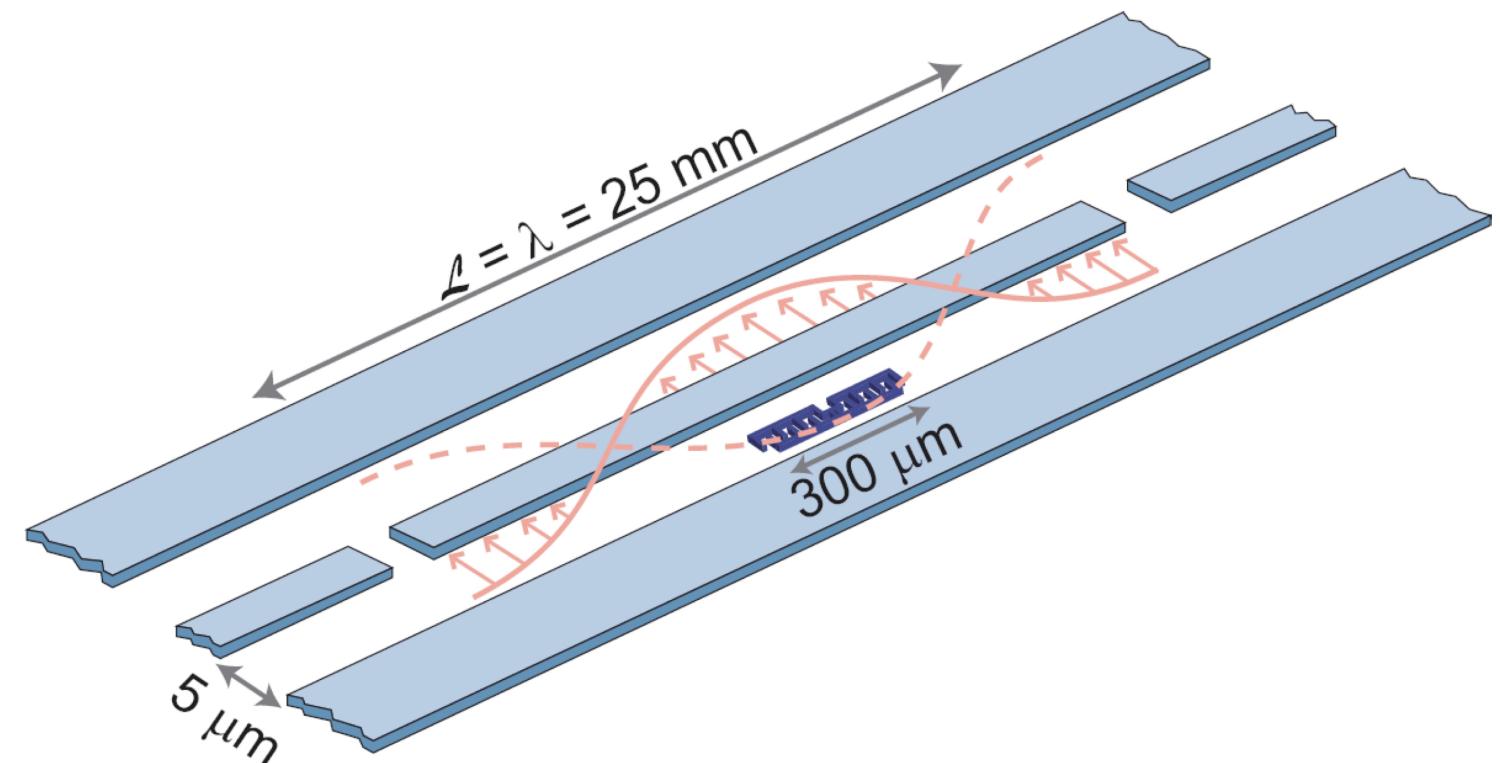
$$\frac{1}{2} \begin{pmatrix} 1 \\ 1 \\ 1 \\ -1 \end{pmatrix} \xrightarrow{\text{iSWAP}} \frac{1}{2} \begin{pmatrix} 1 \\ -i \\ -i \\ -1 \end{pmatrix} = |\circlearrowright\rangle \otimes |\circlearrowright\rangle$$

Experimental setup



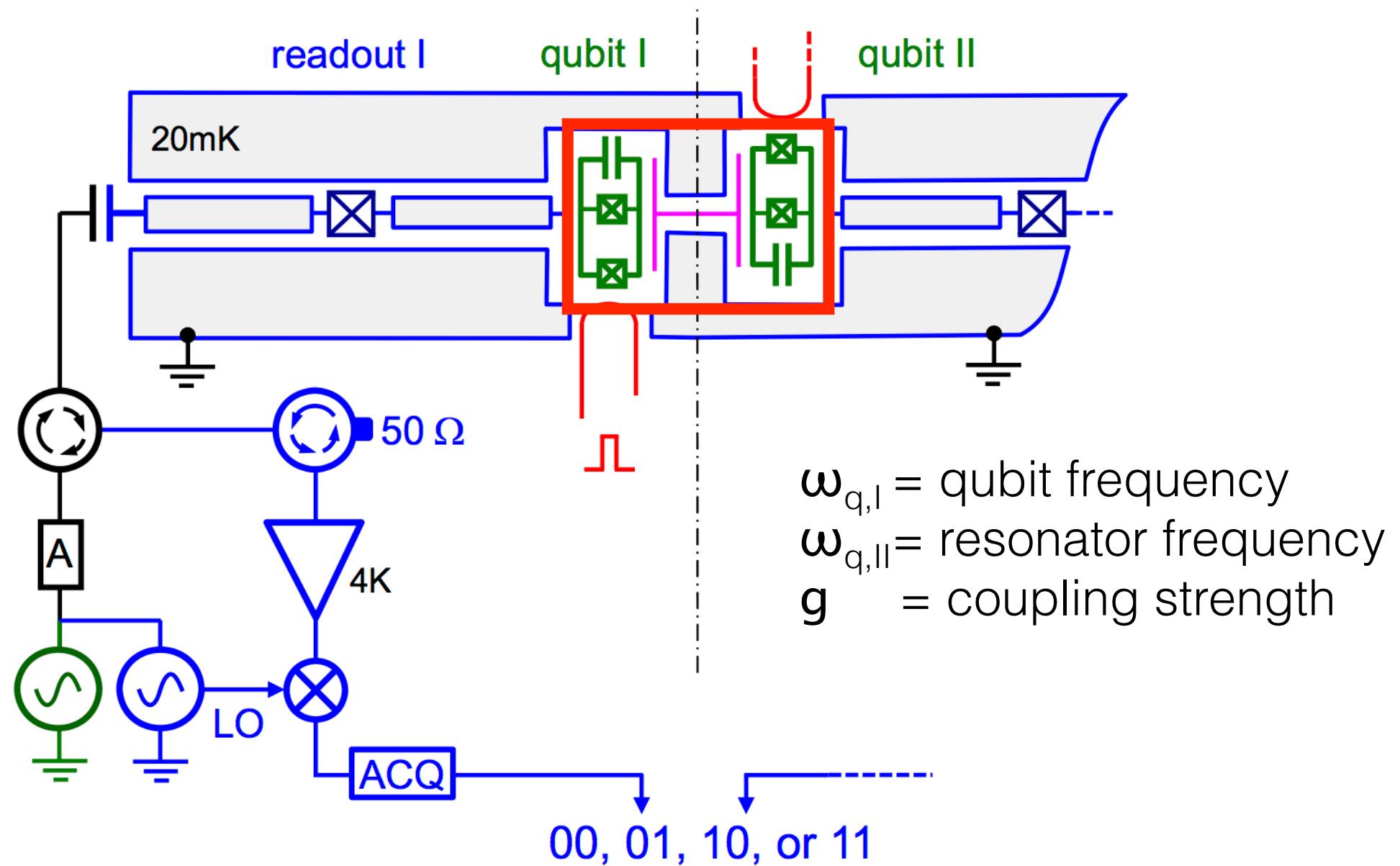
Single qubit manipulation

- Qubit frequency control via flux bias
- Rotations around z axis: detuning Δ
- Rotations around x and y axes: resonant pulses Ω



$$H_{rot} = \frac{\omega_q - \omega}{2} \sigma_z + \frac{\Omega}{2} (\cos \phi \sigma_x + \sin \phi \sigma_y) \equiv \frac{\delta}{2} \sigma_z + \frac{\Omega_x}{2} \sigma_x + \frac{\Omega_y}{2} \sigma_y.$$

Experimental setup



Qubit capacitive coupling

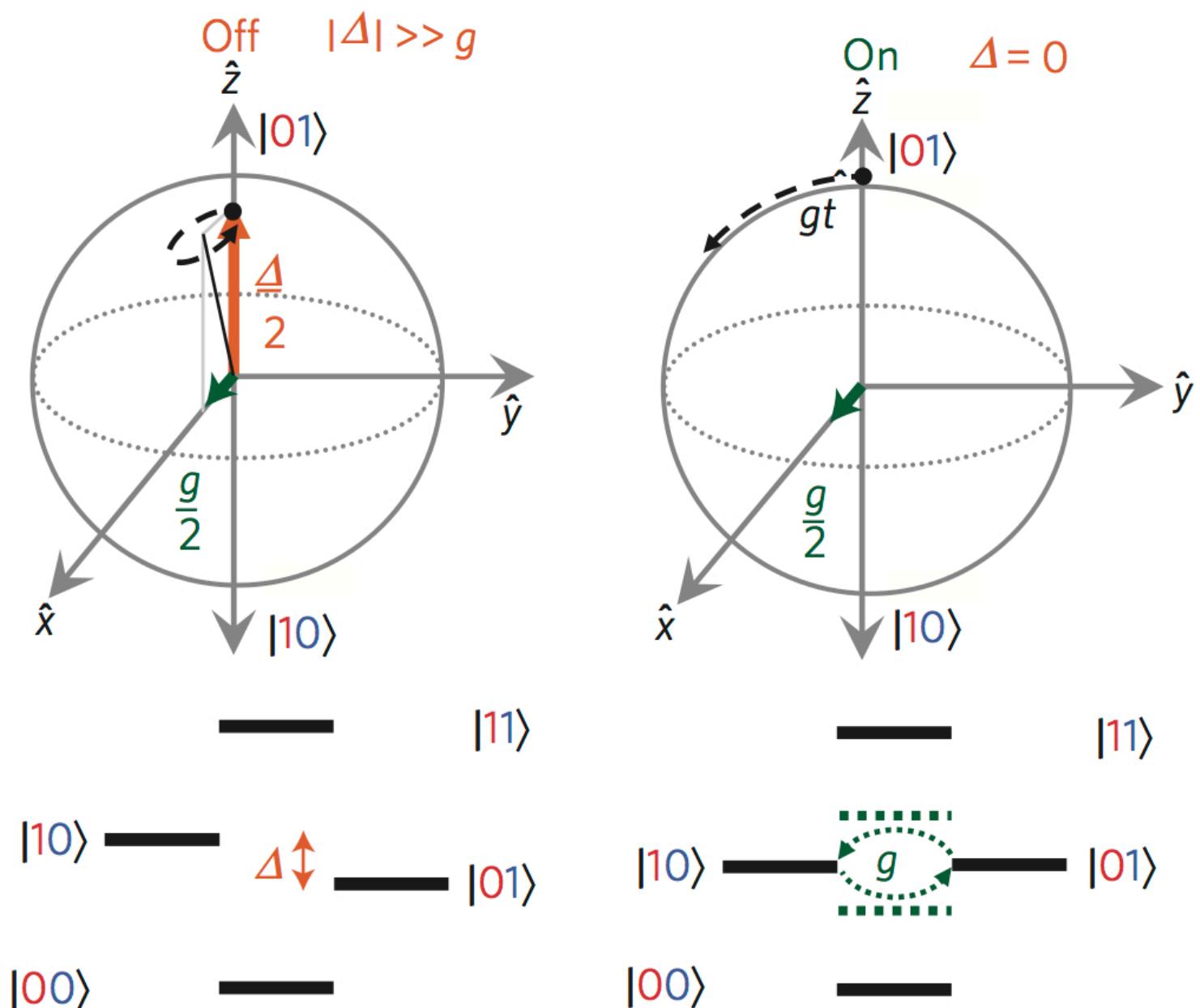
In the rotating frame

$$\omega = \omega_{q,\parallel}$$

the coupling Hamiltonian is:

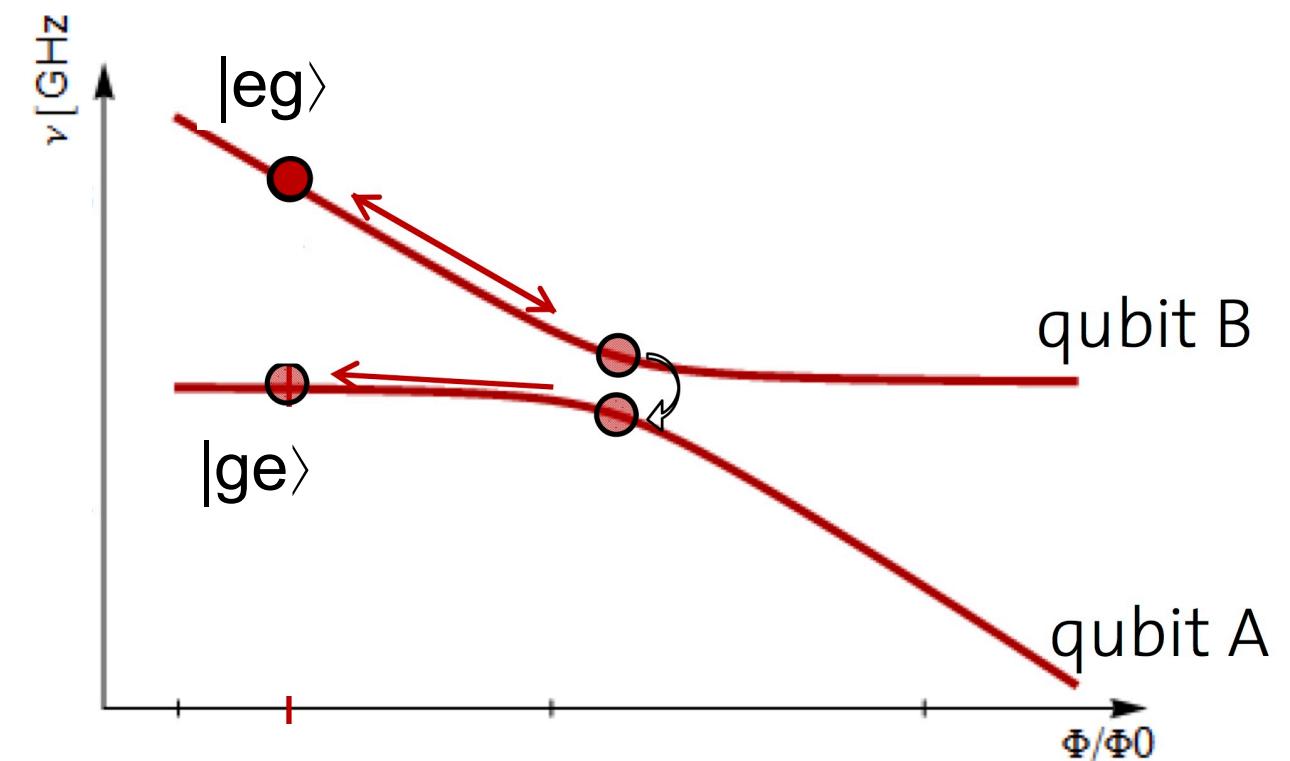
$$H_{tot} = \hbar(\omega_{q,I}\sigma_z^I + \omega_{q,II}\sigma_z^{II} + H_{int})$$

$$H_{int} = g(|10\rangle\langle 01| + |01\rangle\langle 10|)$$



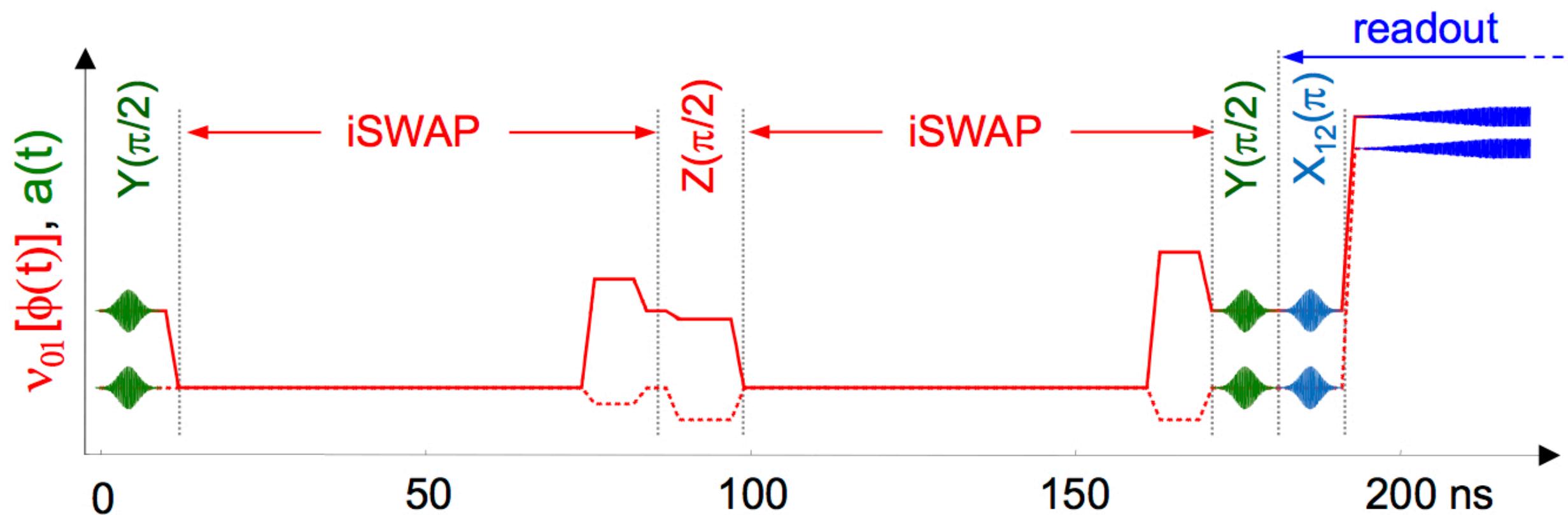
iSWAP gate

- Controlled interaction between $\langle 10 |$ and $\langle 01 |$
- By letting the two states interact for $t = \pi/g$ we obtain an iSWAP gate!

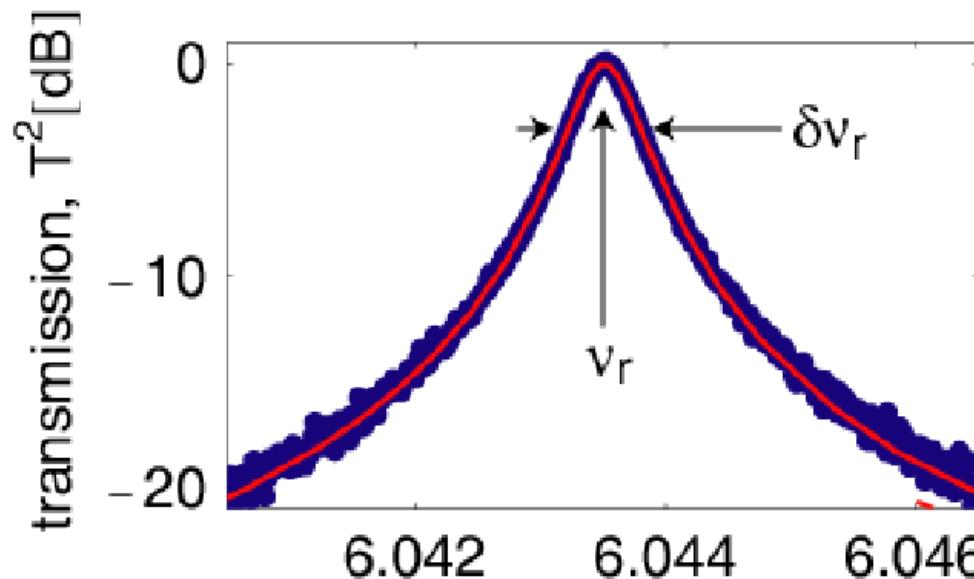


$$U_{\text{int}} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos(gt/2) & -i\sin(gt/2) & 0 \\ 0 & -i\sin(gt/2) & \cos(gt/2) & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

Pulse sequence



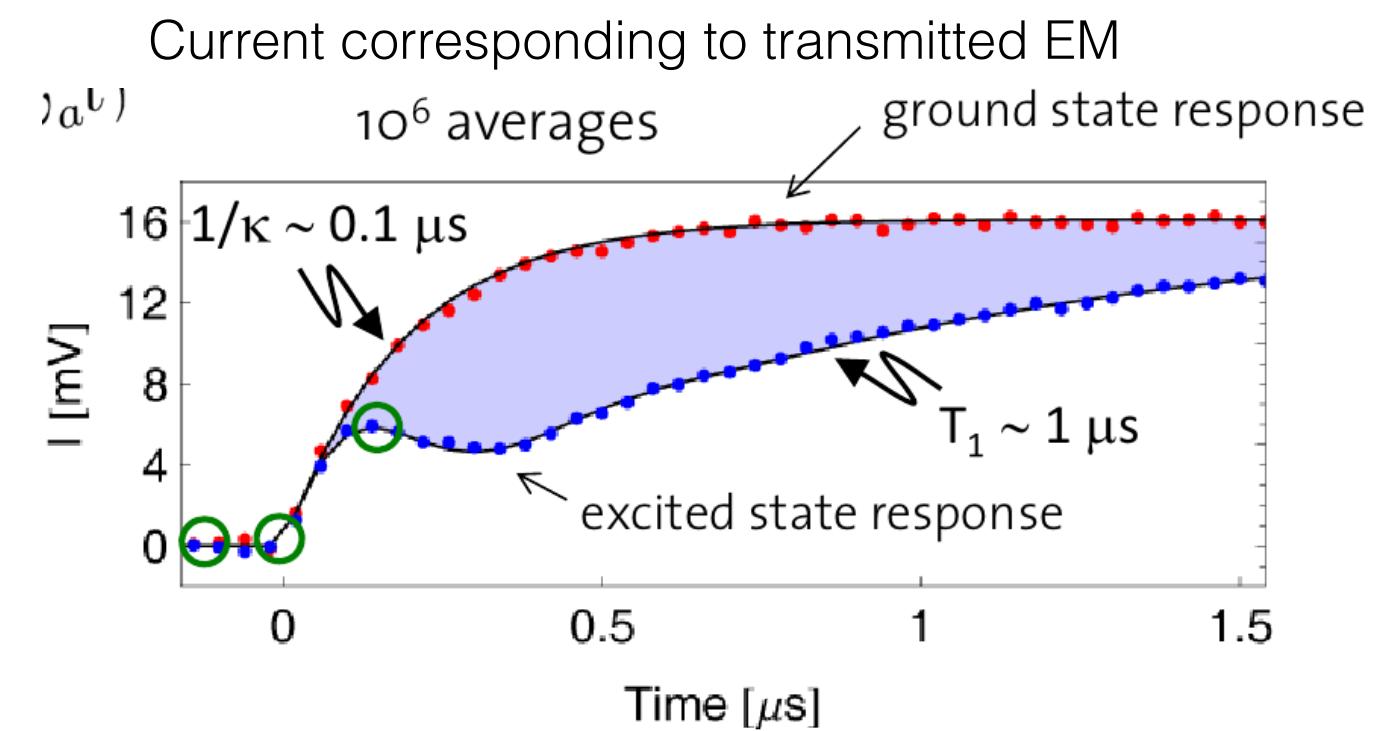
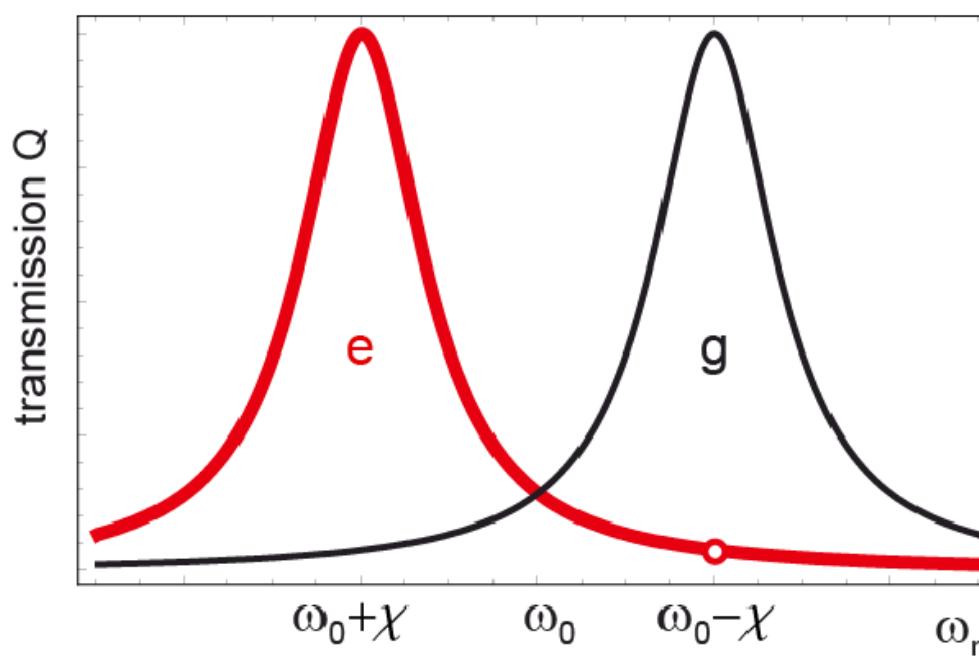
Linear transmission line and single-shot measurement



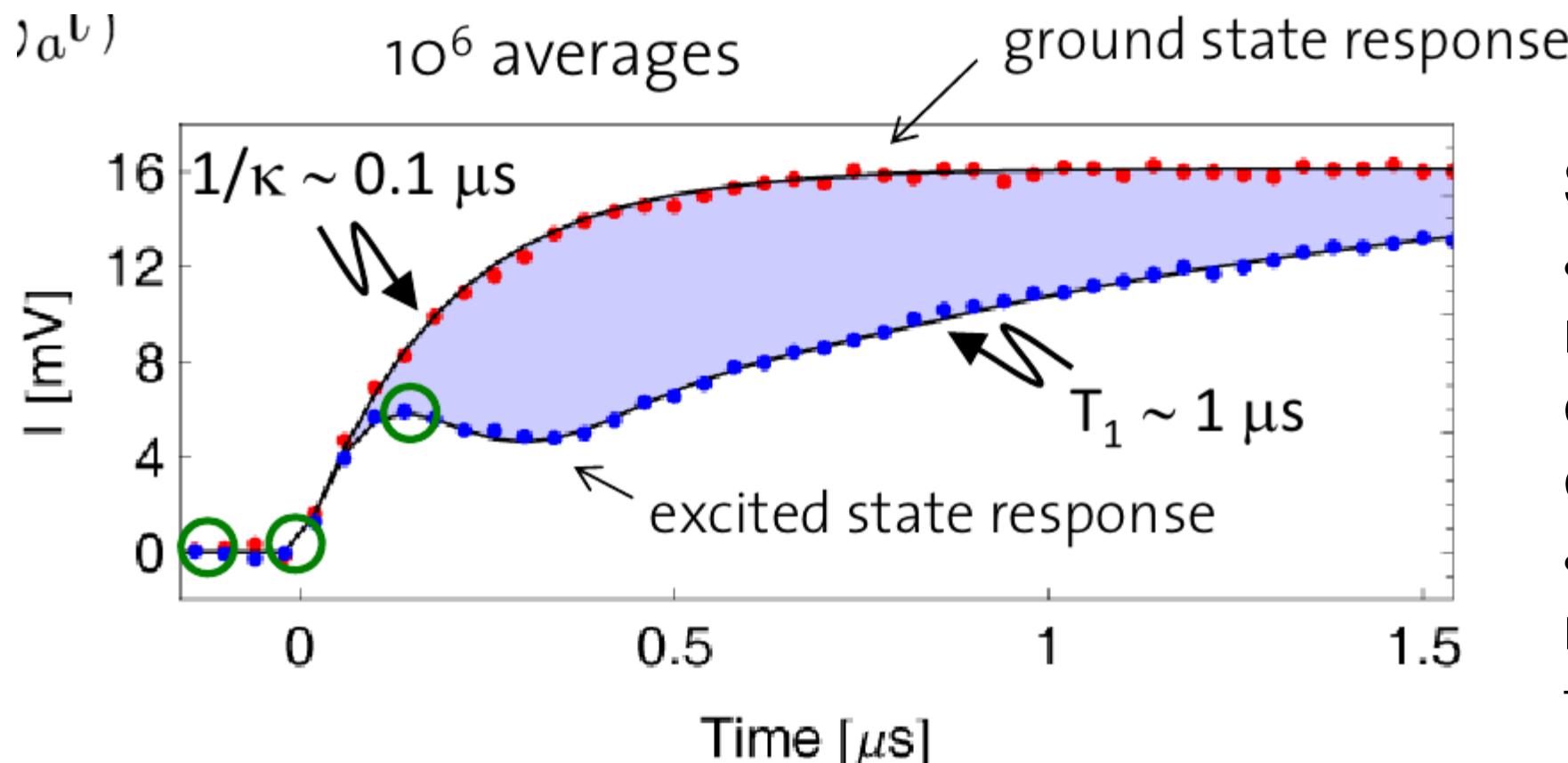
Quality factor of empty linear transmission line $Q = \frac{v_r}{\Delta v_r}$

With resonant frequency v_r

Presence of a transmon in $|g\rangle$ state shifts the resonant frequency of the transmission line $\omega_0 = 2\pi v_r \rightarrow \omega_0 - \chi$
If microwave at $\omega_0 - \chi$ full transmission for $|g\rangle$ state, partial for $|e\rangle$ state



Linear transmission line and single-shot measurement

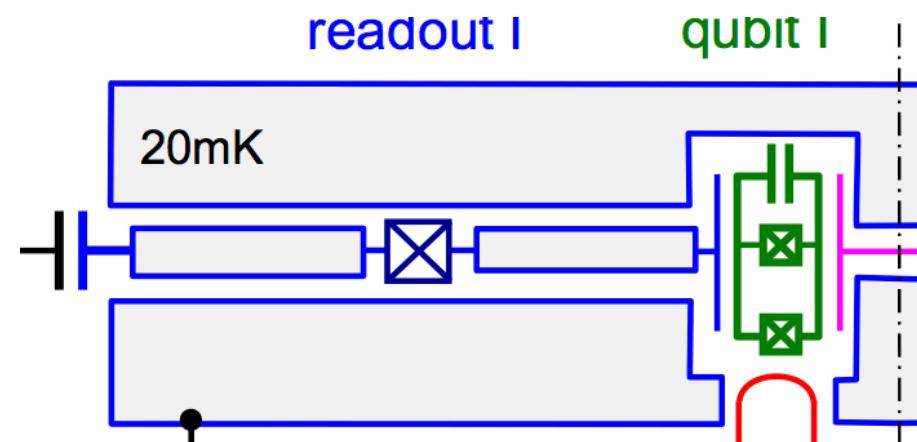


- Single shot
 - If there was no noise, would get either blue or red curve
 - Real curve so noisy that cannot tell whether $|g\rangle$ or $|e\rangle$

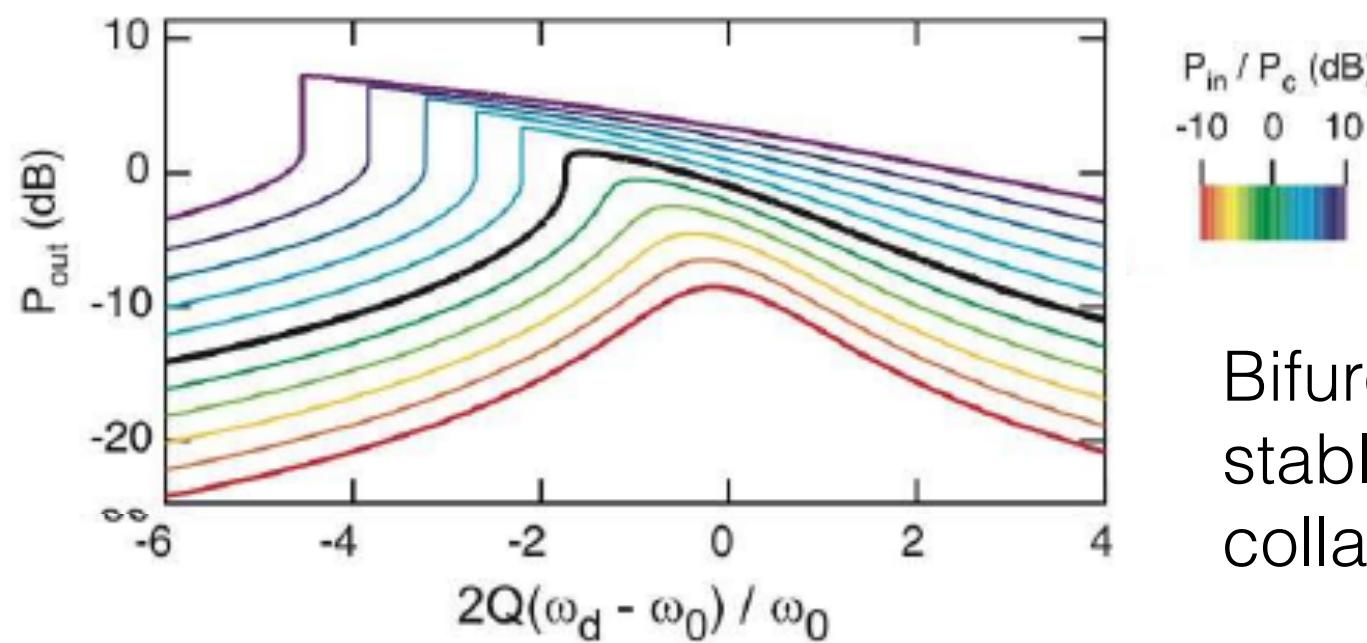
Cannot do single-shot readout

We need an amplifier which increases the area between $|g\rangle$ and $|e\rangle$ curves, but does not amplify the noise

Josephson Bifurcation Amplifier (JBA)



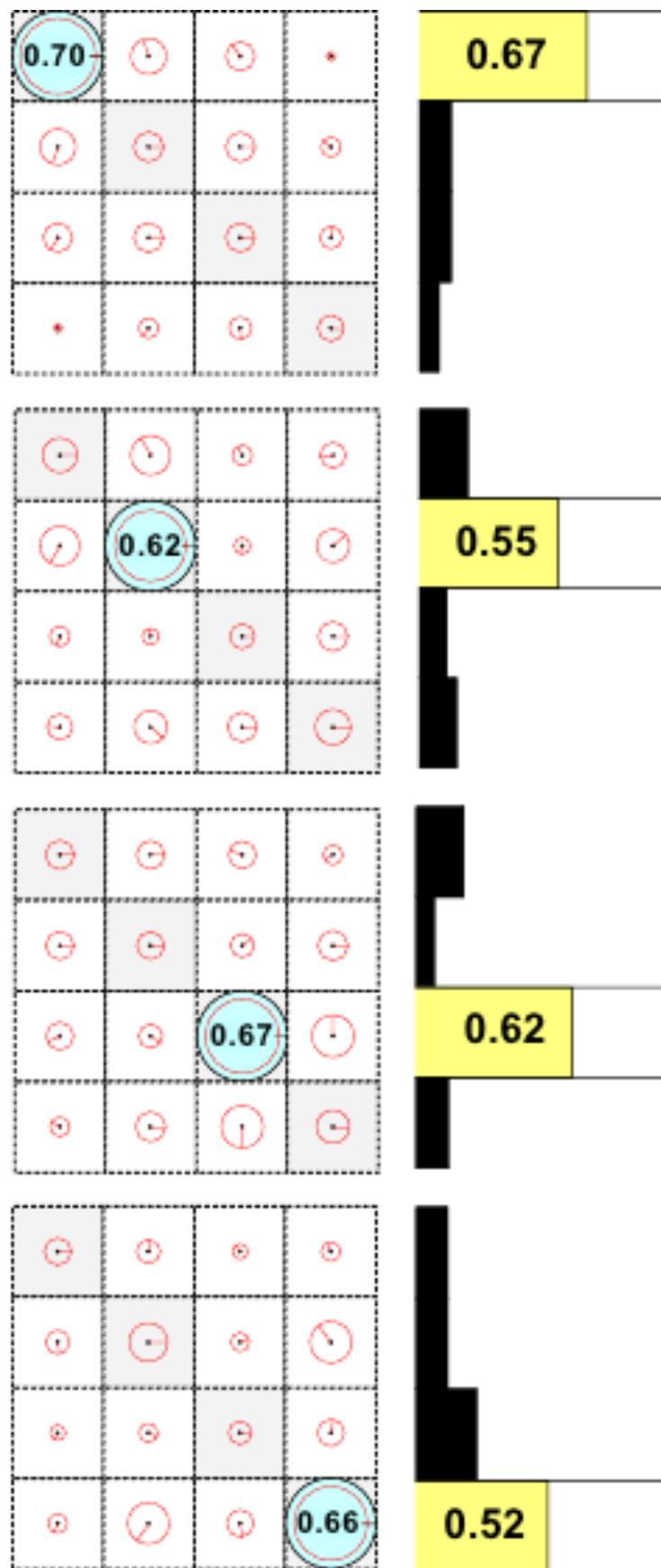
Nonlinear transmission line due to
Josephson junction
Resonant frequency ω_0



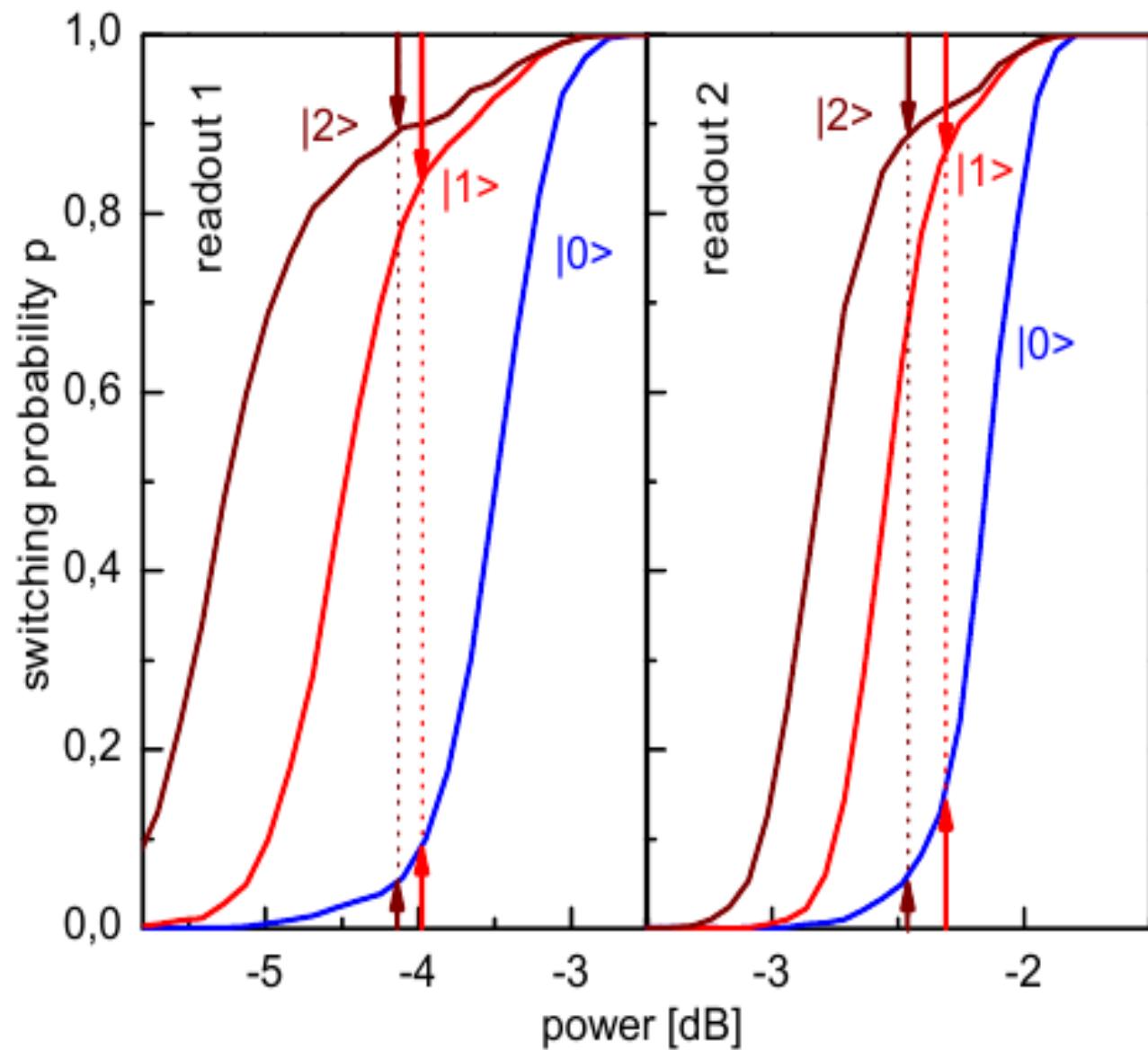
At $P_{\text{in}} = P_c$ max. slope
diverges

Bifurcation: at the correct $(P_{\text{in}}, \omega_d)$ two
stable solutions, can map the
collapsed state of the qubit to them

Josephson Bifurcation Amplifier (JBA) - Excite $|1\rangle \rightarrow |2\rangle$ - Choose power corresponding to the biggest difference of switching probabilities Errors - Nonzero probability of incorrect mapping - Crosstalk Switching probability p : probability that the JBA changes to the **higher-amplitude** solution



Errors



- Nonzero probability of incorrect mapping
- Crosstalk

Conclusions

- Gate operations of Grover algorithm successfully implemented with capacitively coupled transmon qubits
- Arrive at the target state with probability 0.62 – 0.77 (tomography)
- Single-shot readout with JBA (no quantum speed-up without it)
- Measure the target state in single shot with prob 0.52 – 0.67 (higher than 0.25 classically)

Sources

- Dewes, A; Lauro, R; Ong, FR; et al.,
“Demonstrating quantum speed-up in a superconducting two-qubit processor”,
arXiv:1109.6735 (2011)
- Bialczak, RC; Ansmann, M; Hofheinz, M; et al.,
“Quantum process tomography of a universal entangling gate implemented with Josephson phase qubits”,
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Thank you for your
attention

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