

1.11. Quantum Teleportation

Dienstag, 06. März 2012
18:57

(Bennett et al 1993)

Task: Transfer unknown quantum state $|ψ\rangle$ from Alice to Bob at a distance

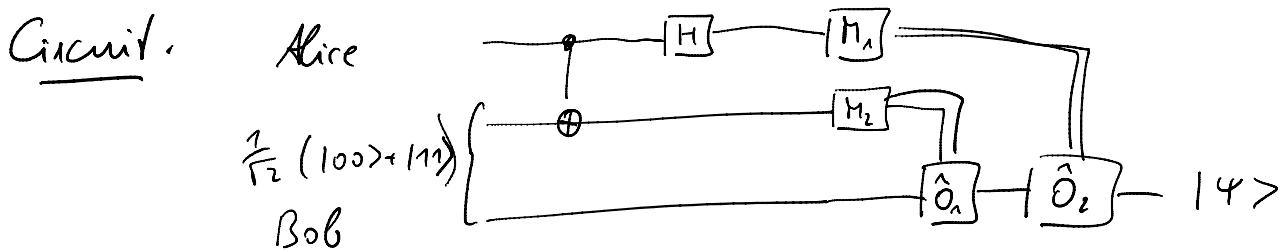
What are the options of Alice?

- *) transport state to Bob
- *) measure state, send classical information
- *) broadcast quantum information

Why is this difficult in quantum mechanics?

- *) quantum states cannot be measured reliably
- *) quantum states are fragile
- *) no-cloning theorem \rightarrow no broadcasting of quantum information

Resources: entangled pair of qubits & classical communication



- Steps:
- ① input (entangled pair)
 - ② CNOT
 - ③ Hadamard
 - ④ measurement
 - ⑤ conditional operations
 - ⑥ output

- Note:
- *) Alice has no information about $|ψ\rangle$ (and cannot obtain it)
 - *) state is always fully transferred

Teleportation Protocol

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① initial state: $(\alpha|0\rangle + \beta|1\rangle) \otimes \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) =$
 $\frac{1}{\sqrt{2}} (\alpha|000\rangle + \alpha|011\rangle + \beta|100\rangle + \beta|111\rangle)$

② CNOT₁₂: $\Rightarrow \frac{1}{\sqrt{2}} (\alpha|000\rangle + \alpha|011\rangle + \beta|110\rangle + \beta|101\rangle)$

③ Hadamard₁: $\Rightarrow \frac{1}{2} (\alpha|000\rangle + \alpha|100\rangle + \alpha|011\rangle + \alpha|111\rangle -$
 $- \beta|110\rangle + \beta|010\rangle - \beta|101\rangle + \beta|001\rangle)$

$$= \frac{1}{2} \left[\begin{array}{l} |00\rangle (\alpha|0\rangle + \beta|1\rangle) + \\ |10\rangle (\alpha|0\rangle - \beta|1\rangle) + \\ |01\rangle (\alpha|1\rangle + \beta|0\rangle) + \\ |11\rangle (\alpha|1\rangle - \beta|0\rangle) \end{array} \right]$$

④ measurement of qubits 1,2
 $M_1 \otimes M_2 \otimes \mathbb{1}$
 $P_{00} = P_{10} = P_{01} = P_{11} = \frac{1}{4}$

⑤ conditional manipulations on post measurement state $|4'\rangle$

$$\left. \begin{array}{l} |00\rangle: \mathbb{1}|4'\rangle \\ |10\rangle: \hat{Z}|4'\rangle \\ |01\rangle: \hat{X}|4'\rangle \\ |11\rangle: \hat{Z}\hat{X}|4'\rangle \end{array} \right\} = \alpha|0\rangle + \beta|1\rangle$$

\Rightarrow requires transfer of two bits of classical info to Bob to perform local operations that recover original state

Applications:
 1) Quantum error correction
 2) Quantum gates
 3) Quantum repeaters

Experiments: D. Bouwmeester et al Nature 380, 575 (1997) (first with photons)

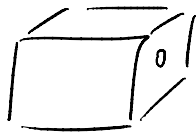
hallmark experiment for quantum information processing

tested with nuclear spins, ions, partly with superconducting qubits

1.12 Mixed States

Montag, 12. März 2012
09:21

How can we describe a statistical mixture of pure states, i.e. the state of a quantum system of which we have only partial information?



↑ ↑ ↑ ↑ ↑ ↑ e.g.

$ \psi\rangle$	prob
↑	0.5
↑	0.5

pure states $|\psi_i\rangle = \cos \frac{\theta_i}{2} |\uparrow\rangle + \sin \frac{\theta_i}{2} |\downarrow\rangle$ $\{\theta_i = 0, \delta\}$ with probabilities $\{p_1, p_2\}$.

→ $|\tilde{\psi}\rangle = p_1 |\psi_1\rangle + p_2 |\psi_2\rangle$? Is this the proper description?

Does not describe statistical mixture!

Main 'problem': coherence between basis states leads to interference!

Density matrix:

$$\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i| \quad \sum p_i = 1$$

$$\Rightarrow \rho_{ex} = 0.5 |\uparrow\rangle\langle\uparrow| + 0.5 |\downarrow\rangle\langle\downarrow|$$

$$|\uparrow\rangle\langle\uparrow| = \begin{pmatrix} 1 \\ 0 \end{pmatrix} (1 \ 0) = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$|\downarrow\rangle\langle\downarrow| = \begin{pmatrix} \cos \frac{\delta}{2} \\ \sin \frac{\delta}{2} \end{pmatrix} (\cos \frac{\delta}{2} \ , \ \sin \frac{\delta}{2}) = \begin{pmatrix} \cos^2 \frac{\delta}{2} & \cos \frac{\delta}{2} \sin \frac{\delta}{2} \\ \cos \frac{\delta}{2} \sin \frac{\delta}{2} & \sin^2 \frac{\delta}{2} \end{pmatrix}$$

$$\rho = \frac{1}{2} \begin{pmatrix} 1 + \cos^2 \frac{\delta}{2} & \cos \frac{\delta}{2} \sin \frac{\delta}{2} \\ \cos \frac{\delta}{2} \sin \frac{\delta}{2} & 1 + \sin^2 \frac{\delta}{2} \end{pmatrix}$$

General parametrization & properties

$$\rho = \begin{pmatrix} a & b+ic \\ b-ic & d \end{pmatrix}$$

→ ρ is hermitian: $\rho^\dagger = \rho$ (real eigenvalues $\hat{=}$ real valued probabilities)

→ $\text{Tr}[\rho] = 1$ ($\sum p_i = 1$)

→ positive semidefinite $\hat{=}$ all eigenvalues ≥ 0
(no negative probabilities)

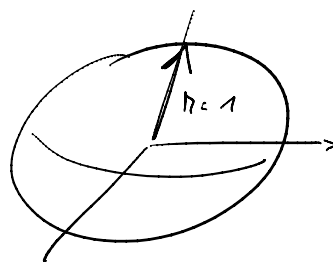
e.g. for $S = \bar{n}$: $\rho = \frac{1}{2}(|\uparrow\rangle\langle\uparrow| + |\downarrow\rangle\langle\downarrow|) = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$

$$p_\uparrow = p_\downarrow = 0.5$$

Bloch ball representation: $\rho = \frac{1}{2}(\mathbb{1} + \vec{n} \cdot \vec{\sigma})$ $|\vec{n}| \leq 1$

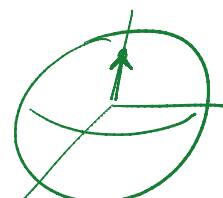
\vec{n} ... vector inside Bloch sphere

e.g. $\vec{n} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \hat{=} |1 \times 1|$

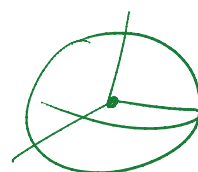


\Rightarrow pure states have $|\vec{n}| = 1$!

→ $\vec{n} = \begin{pmatrix} 0 \\ 0 \\ 0.5 \end{pmatrix}$? $\rho = \frac{1}{2}(\mathbb{1} + 0.5\sigma_z) = \begin{pmatrix} 3/4 & 0 \\ 0 & 1/4 \end{pmatrix}$



→ \vec{n} for $p_1 = p_2 = 0.5$? $\rho = \begin{pmatrix} 0.5 & 0 \\ 0 & 0.5 \end{pmatrix}$



\rightarrow totally mixed state, no information about state of system

Measurement

operator \hat{M} : $\langle \hat{M} \rangle = \text{Tr}[\hat{M}\rho]$... average value of \hat{M}

$$\text{Tr}[\hat{M}\rho] = \sum_k \langle k | \hat{M} \rho | k \rangle \dots \text{sum over basis states}$$

$$= \sum_{k,c} \langle k | \hat{M} | c \rangle \langle c | \rho | k \rangle =$$

$$= \sum_{k,c} \langle k | \hat{M} | c \rangle \sum_m p_m \underbrace{\langle c | m \rangle}_{\delta_{cm}} \underbrace{\langle m | k \rangle}_{\delta_{mk}} =$$

$$= \sum_m \langle m | \hat{M} | m \rangle p_m$$

↳ expectation value for pure state $|m\rangle$

$$(p_1=1, p_0=0 : \text{Tr}[\hat{M}\rho] = \langle 1 | \hat{M} | 1 \rangle \hat{=}$$

$$\langle \psi | \hat{M} | \psi \rangle = \langle \hat{M} \rangle_\psi$$

e.g.: determine coeffs of $\rho = \frac{1}{2}(1 + \vec{n} \cdot \vec{\sigma})$

$$\text{Tr}[\rho \sigma_{x,y,z}] = n_x, n_y, n_z$$

Unitary evolution:

$$\rho(t) = U \rho(0) U^\dagger$$

$$(\text{pure state: } |\psi(t)\rangle \langle \psi(t)| = U |\psi(0)\rangle \langle \psi(0)| U^\dagger)$$

1.13 Bloch Equations, T1 & T2 times

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Unitary Evolution: $H = -\mu \cdot \vec{\sigma} \cdot \vec{m}(t)$ $\vec{m}(t) = \begin{pmatrix} m_x(t) \\ m_y(t) \\ m_z(t) \end{pmatrix}$

(Liouville-) von Neumann equation:

$$\dot{\rho} = -\frac{i}{\hbar} [H, \rho] = -\frac{i}{\hbar} [H \cdot \rho - \rho \cdot H] \quad \rho = \frac{1}{2} (\mathbb{1} + \vec{\lambda} \cdot \vec{\sigma})$$

$$\dot{\rho} = \frac{i\mu}{\hbar} \left[\vec{\sigma} \cdot \vec{m} \cdot \frac{1}{2} (\mathbb{1} + \vec{\lambda} \cdot \vec{\sigma}) - \frac{1}{2} (\mathbb{1} + \vec{\lambda} \cdot \vec{\sigma}) \vec{\sigma} \cdot \vec{m} \right] =$$

$$= \frac{i\mu}{2\hbar} \left[(\vec{m} \cdot \vec{\sigma}) \cdot (\vec{\lambda} \cdot \vec{\sigma}) - (\vec{\lambda} \cdot \vec{\sigma}) (\vec{\sigma} \cdot \vec{m}) \right]$$

using $(\vec{a} \cdot \vec{\sigma})(\vec{b} \cdot \vec{\sigma}) = (\vec{a} \cdot \vec{b}) \mathbb{1} + i(\vec{a} \times \vec{b}) \cdot \vec{\sigma}$

$$= \frac{i\mu}{\hbar} (-i \vec{\lambda} \times \vec{m}) \cdot \vec{\sigma}$$

$$= \frac{\mu}{\hbar} (\vec{\lambda} \times \vec{m}) \cdot \vec{\sigma}$$

in terms of Bloch vector components:

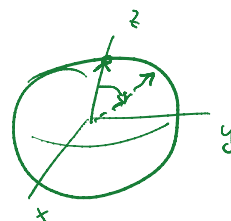
$$\dot{\vec{\lambda}} = T_1 \left[\dot{\rho} \cdot \vec{\sigma} \right] = \left(\frac{\mu}{\hbar} \right) (\vec{\lambda} \times \vec{m}) \quad \dots \text{ 'Bloch equations' }$$

e.g.: $\vec{\lambda}(0) = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$ $m(t) = m_x$

$$\dot{\lambda}_x = \gamma (\lambda_y m_z - \lambda_z m_y) = 0$$

$$\dot{\lambda}_y = \gamma (\lambda_z m_x - \lambda_x m_z) = \gamma \lambda_z m_x$$

$$\dot{\lambda}_z = \gamma (\lambda_x m_y - \lambda_y m_x) = -\gamma \lambda_y m_x$$



} coupled diff. eq.

Can one obtain a mixed state ($|\vec{\lambda}| < 1$) out of a pure state ($|\vec{\lambda}| = 1$)?

\Rightarrow introduce energy relaxation and dephasing!

Dissipative evolution:

$$\dot{\eta}_x = \gamma (\vec{\eta} \times \vec{m})_x - \frac{\eta_x}{T_2}$$

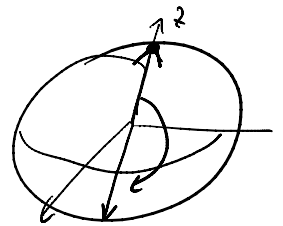
$$\dot{\eta}_y = \gamma (\vec{\eta} \times \vec{m})_y - \frac{\eta_y}{T_2}$$

$$\dot{\eta}_z = \gamma (\vec{\eta} \times \vec{m})_z - \frac{\eta_z + 1}{T_1}$$

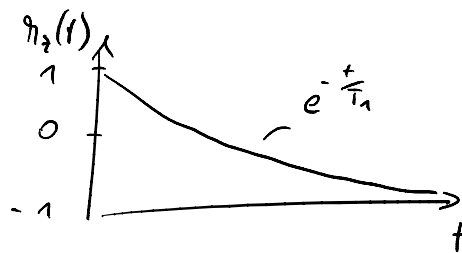
Energy relaxation: $\vec{m} = \vec{0}$, $T_2 = 0$, $\boxed{T_1 \neq 0}$

$$\dot{\eta}_z = -(\eta_z + 1)/T_1$$

$$\eta_z(t) = -1 + 2e^{-t/T_1} \text{ for } \eta_z(0) = 1$$



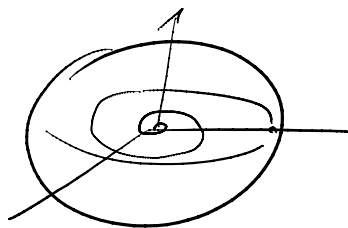
$\Rightarrow T_1$... longitudinal relaxation time, energy relaxation



Dephasing: $\vec{m} = m_z$, $T_1 = 0$, $\boxed{T_2 \neq 0}$

$$\left. \begin{aligned} \dot{\eta}_y &= -\gamma \eta_x m_z - \frac{\eta_y}{T_2} \\ \dot{\eta}_x &= \gamma \eta_y m_z - \frac{\eta_x}{T_2} \end{aligned} \right\} \eta_x(t), \eta_y(t) \propto e^{-t/T_2}$$

unitary dissipative



T_2 : transverse relaxation time, dephasing