

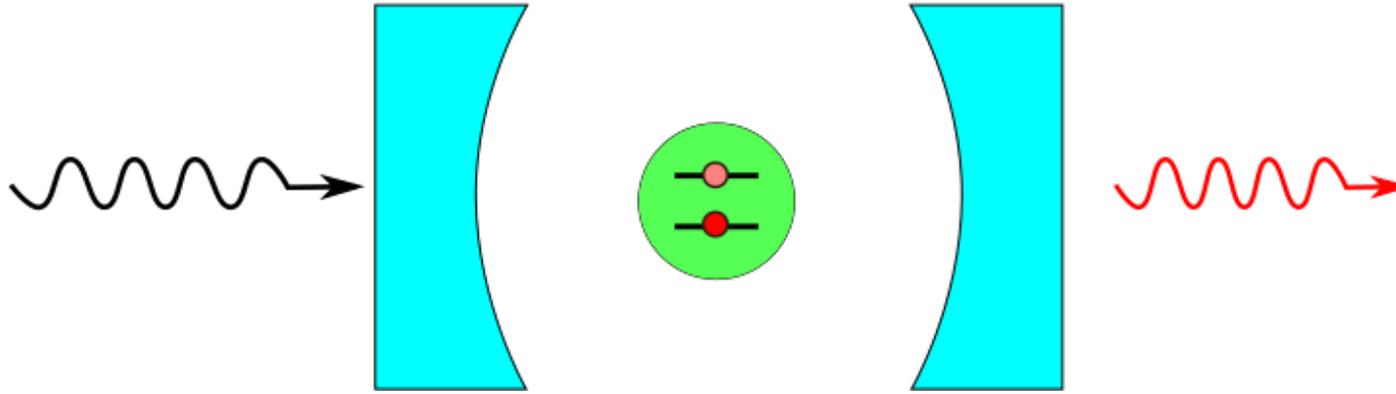
Superconducting Qubits Readout

Bruno Eckmann, Florian Schroeder
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Motivation

- Properties we want for a QM-measurement:
 - Projective
 - Non-demolition (QND)
 - High SNR
 - Faster than life- and coherence time (T_1 , T_2)
 - High fidelity

Setup

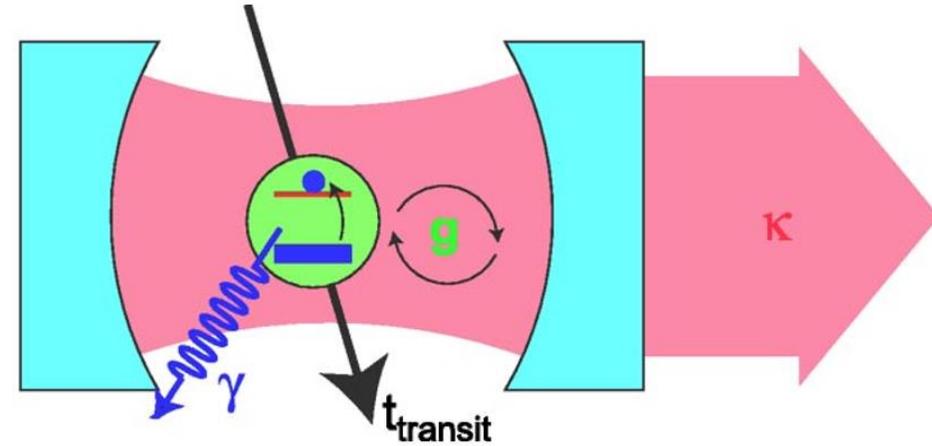


- Put qubit inside cavity
- Measure the transmitted / reflected radiation of the cavity
- Transform quantum qubit state to macroscopic observable
- This is analogous to e.g. Stern-Gerlach

CQED introduction

- Two-level atom inside cavity
- Jaynes-Cummings Hamiltonian:

$$H = \underbrace{\hbar\omega_r \left(a^\dagger a + \frac{1}{2} \right)}_{\text{Cavity field}} + \underbrace{\frac{\hbar\omega_q}{2} \sigma^z}_{\text{qubit}} + \underbrace{\hbar g (a^\dagger \sigma^- + \sigma^+ a)}_{\text{coupling}} + \underbrace{H_\kappa + H_\gamma}_{\text{environment}}$$



- Coupling strength:
- Detuning:
- Cavity and qubit to photon decay rate:

$$g = \frac{\mathcal{E}_{rms} d}{\hbar}$$

$$\Delta := \omega_q - \omega_r$$

$$\kappa, \gamma$$

A. Blais et al., PRA 69, 062320 (2004)

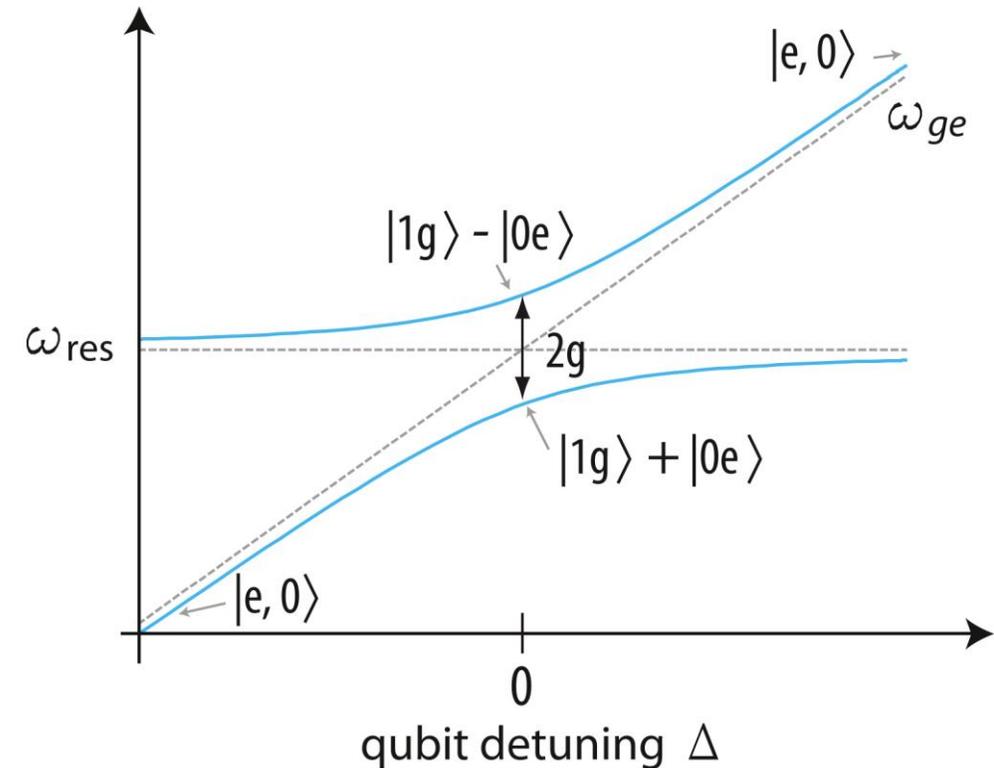
CQED introduction

- Neglecting damping, exact diagonalization of Jaynes-Cummings Hamiltonian yields «dressed» (entangled) states:

$$|+, n\rangle = \cos(\theta_n) |e, n\rangle + \sin(\theta_n) |g, n + 1\rangle$$
$$|-, n\rangle = -\sin(\theta_n) |e, n\rangle + \cos(\theta_n) |g, n + 1\rangle$$

$$E_{\pm, n} = (n + 1)\hbar\omega_r \pm \frac{\hbar}{2} \sqrt{4g^2(n + 1) + \Delta^2}$$

$$\tan(2\theta_n) = \frac{2g\sqrt{n + 1}}{\Delta}$$

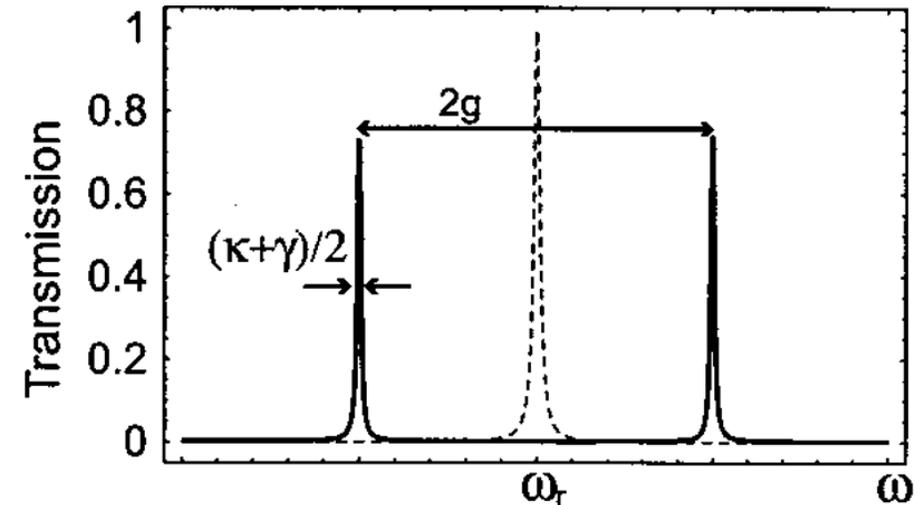
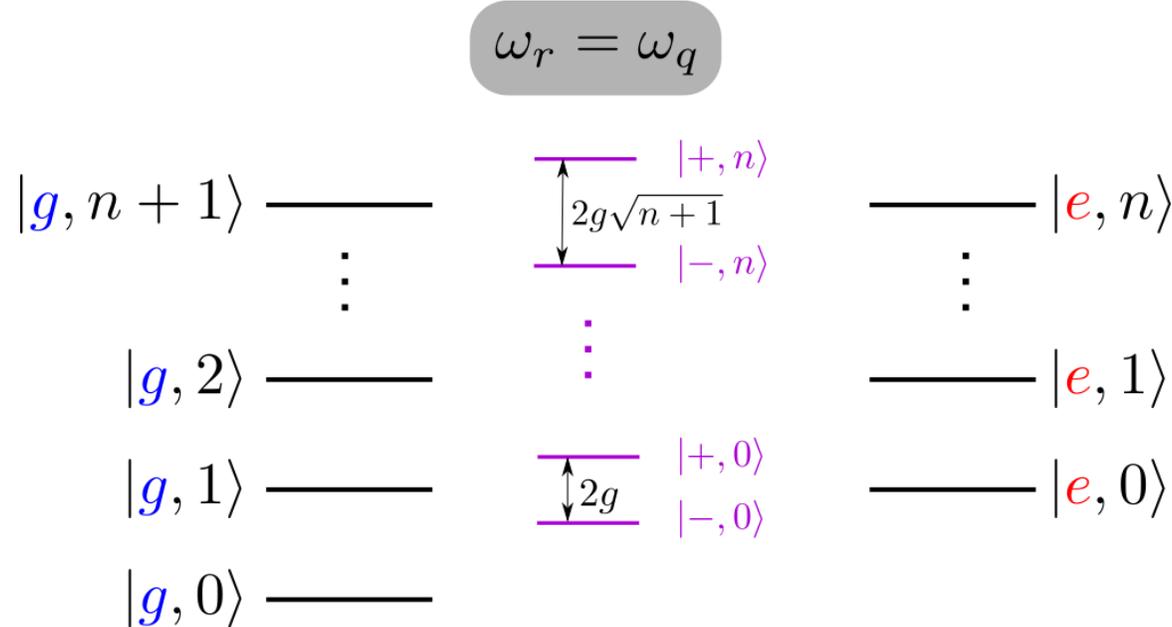


C. Eichler, Ph.D. thesis, ETH Zurich (2013)

CQED introduction

Case 1: zero detuning

- Strong coupling limit: $g \gg \kappa, \gamma$
- Energy spectrum:



A. Blais et al., PRA 69, 062320 (2004)

CQED introduction

Case 2: large detuning (dispersive regime)

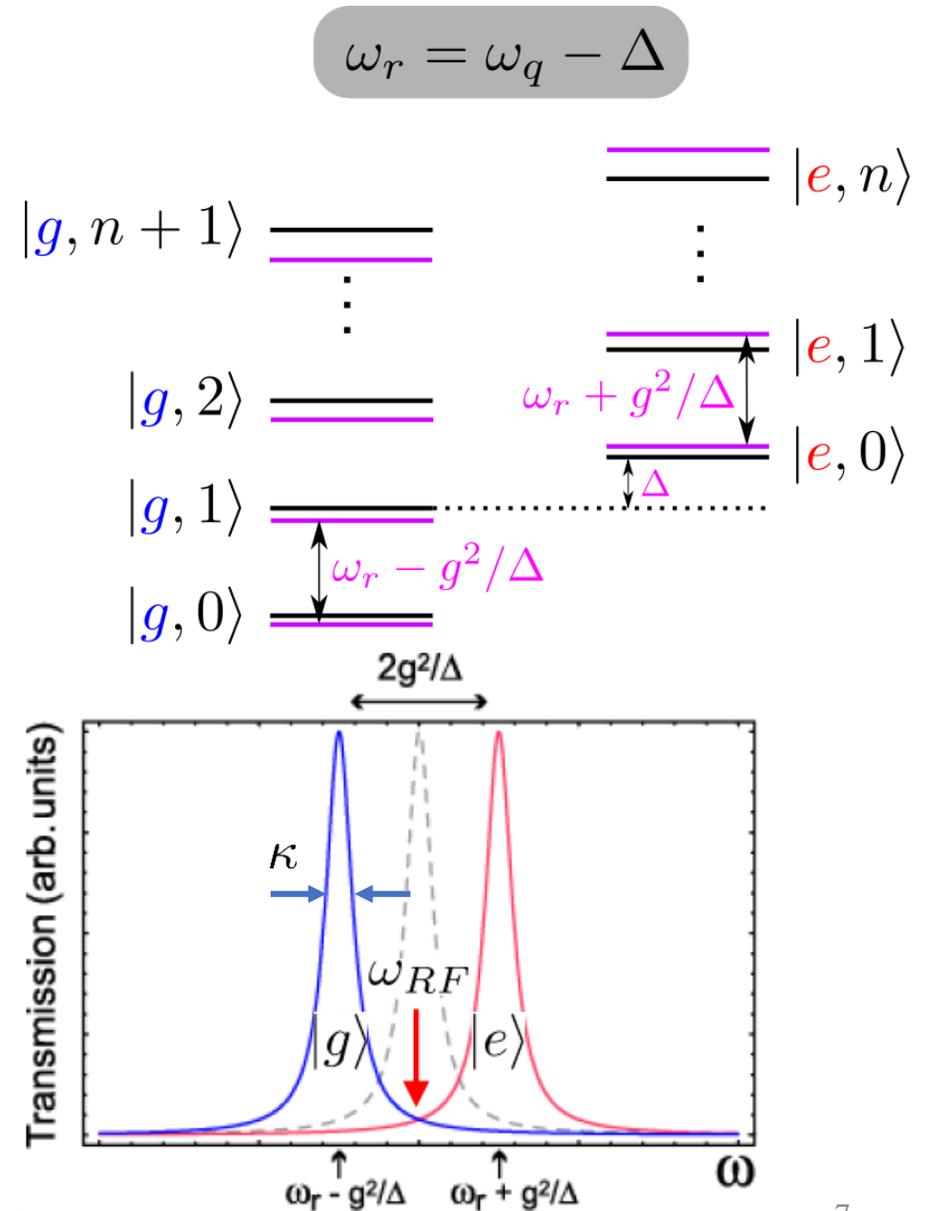
$$\frac{g}{\Delta} \ll 1 \Rightarrow \begin{aligned} |+, n\rangle &\rightarrow |e, n\rangle \\ |-, n\rangle &\rightarrow |g, n+1\rangle \end{aligned}$$

→ qubit states become eigenstates

→ qubit flip suppressed as $\sim \epsilon \frac{g}{\Delta}$

- Also qubit lifetime enhancement; $\frac{\gamma}{\kappa} \rightarrow 0$
- Schrieffer-Wolff transformation:

$$UHU^\dagger \approx \underbrace{\hbar \left(\omega_r + \frac{g^2}{\Delta} \sigma^z \right)}_{\omega_r(\sigma^z)} a^\dagger a + \frac{\hbar}{2} \left(\omega_q + \frac{g^2}{\Delta} \right) \sigma^z$$



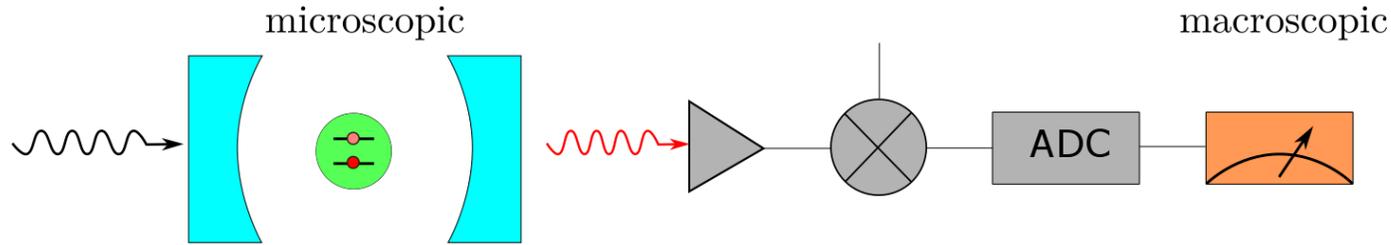
CQED introduction

Case 2: large detuning (dispersive regime)

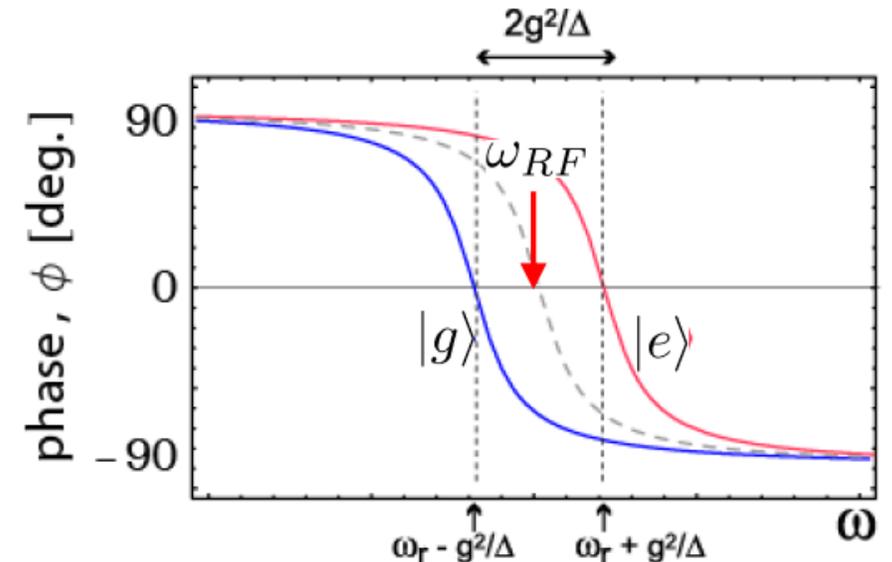
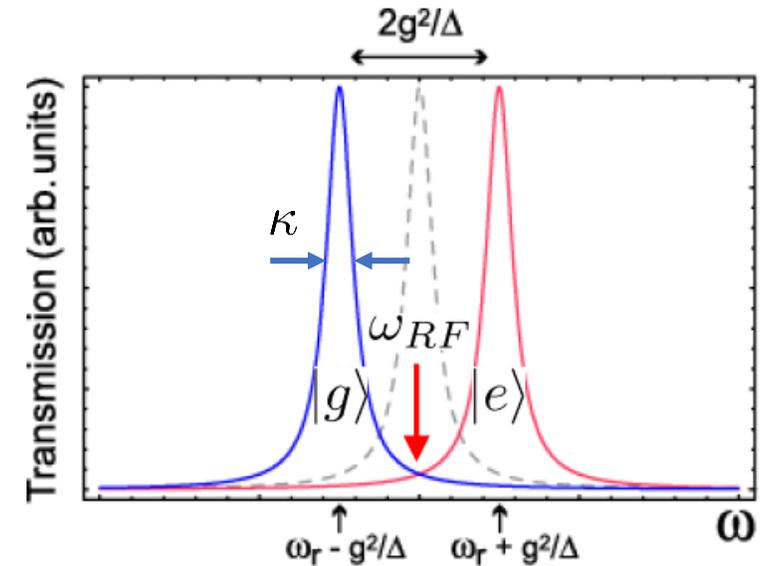
- Schrieffer-Wolff transformation is an approximate diagonalization of Hamiltonian
- Interaction Hamiltonian now commutes with qubit Hamiltonian
 - no qubit flip
 - **non-demolition measurement (QND)**

Dispersive Readout of Qubits

Measurement Strategies



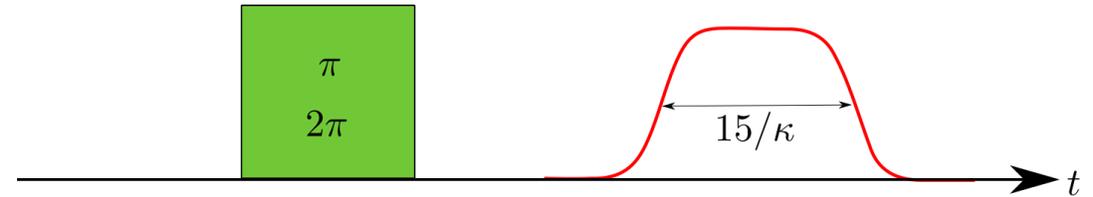
- The transmission spectrum presents a peak of width κ at $\omega_r - g^2/\Delta$ or $\omega_r + g^2/\Delta$, depending on the qubit state
- The phase jumps in the spectrum are shifted, depending on the qubit state
- Two measurement strategies
 1. Choose $\omega_d = \omega_r \pm g^2/\Delta$ and measure transmission amplitude.
 2. Choose $\omega_d = \omega_r$ and measure the phase of the transmitted waves.



Dispersive Readout of Qubits

Measuring **Transmission Amplitude** (Simulation)

- Qubit prepared in ground or excited state
- Measurement pulse duration $\sim 15/\kappa$
- Drive frequency is $\omega_d = \omega_r + g^2/\Delta$
- Hyperbolic tangent rise and fall

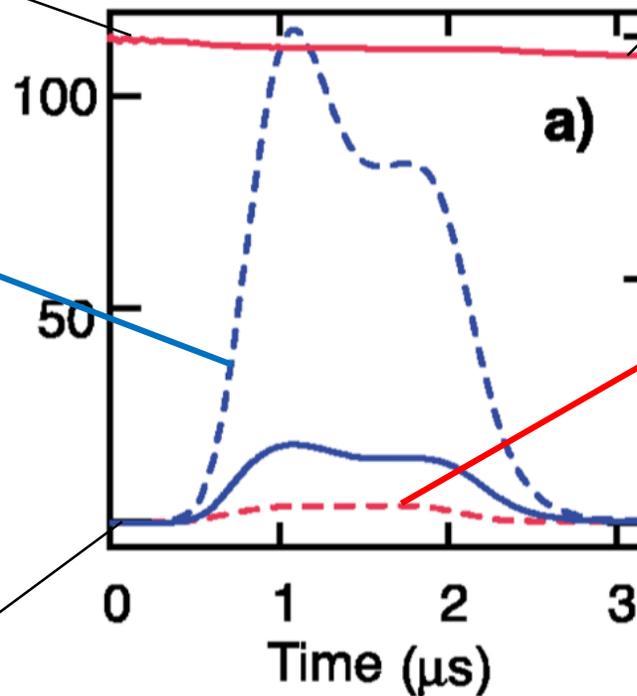


qubit initially in excited state

for a qubit in the ground state, most of the photons are reflected

qubit initially in ground state

Cavity Photon Number



Excited State Population

qubit in excited state after measurement

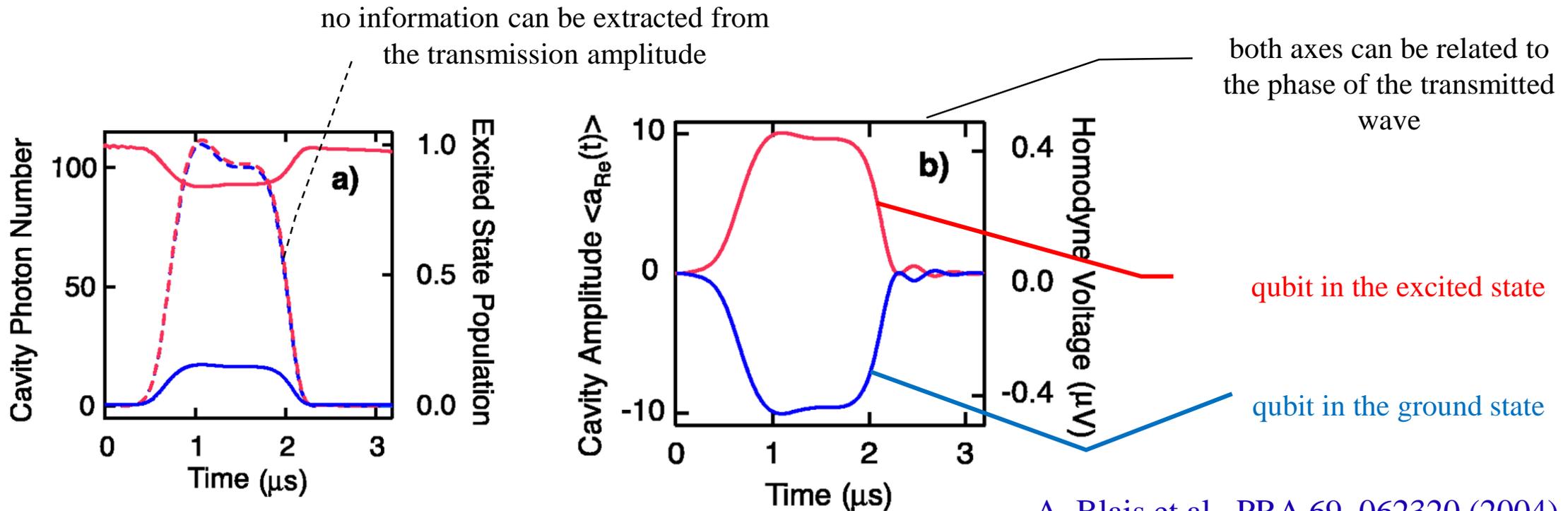
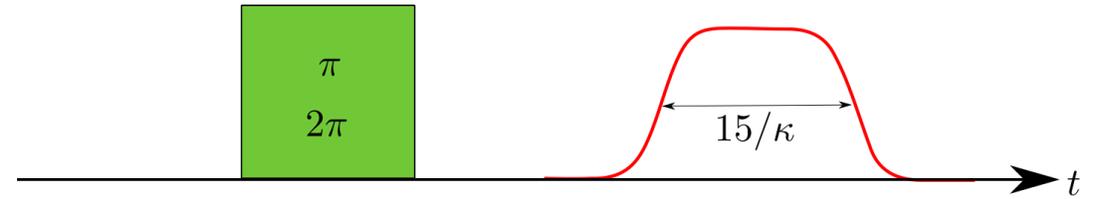
for a qubit in the excited state (red), most of the photons are transmitted

qubit in ground state after measurement

Dispersive Readout of Qubits

Measuring **Phase Shift** (Simulation)

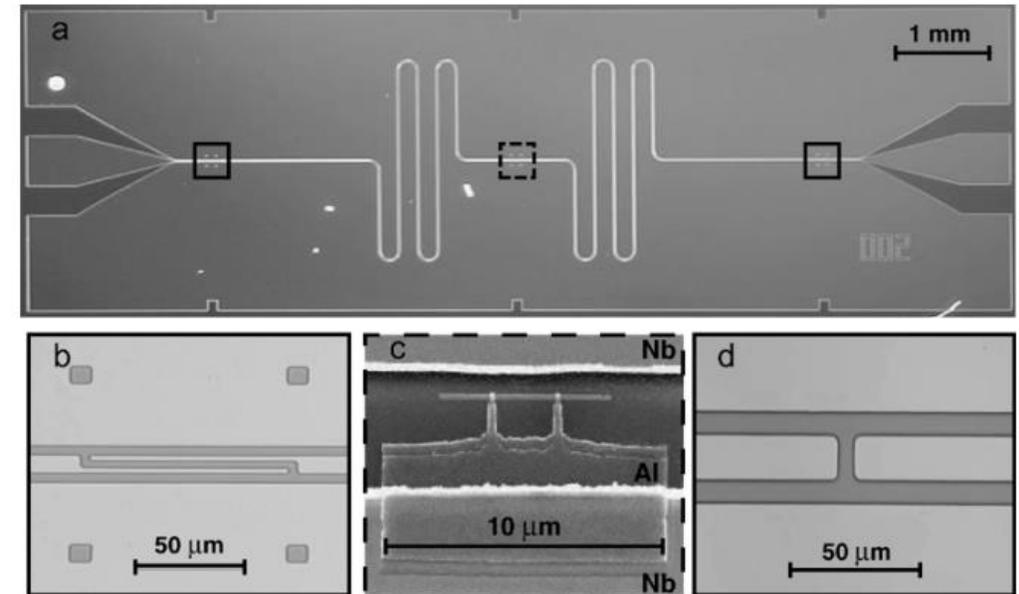
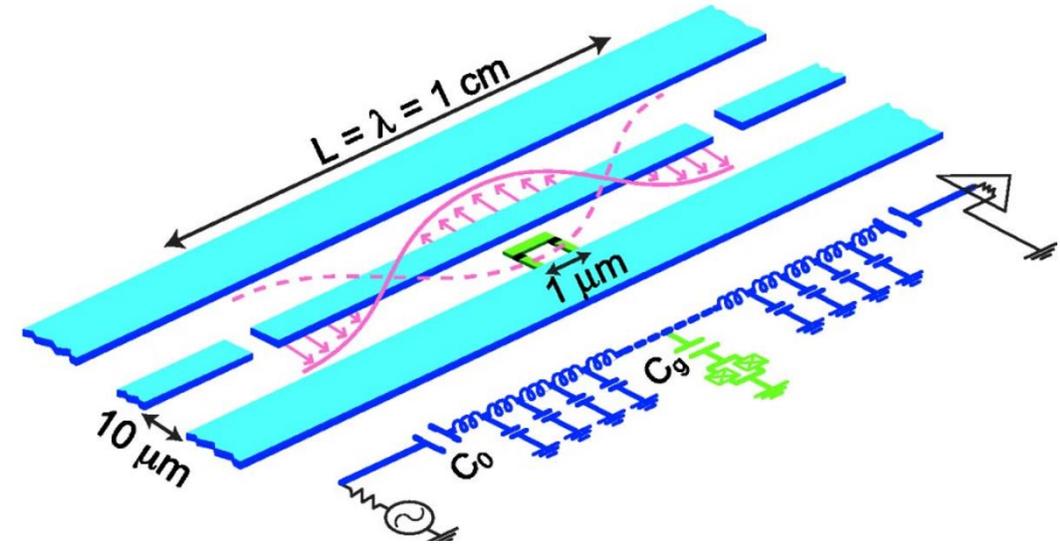
- Qubit prepared in ground or excited state
- Measurement pulse duration $\sim 15/\kappa$
- Drive frequency is $\omega_d = \omega_r$
- Hyperbolic tangent rise and fall



Cavity Quantum Electrodynamics for SC Circuits

Design Example

- 1D transmission line resonator: superconducting coplanar waveguide, geometry defines resonance frequency ω_r
- Small effective volume of the resonator allow **large coupling strength**, as the electric fields are some ~ 100 times larger than in 3D microwave resonators
- Resonator is coupled via **capacitive gaps** (Fig. b, d) to input and output. Their design defines κ
- Qubit (here two Cooper-pair boxes) is capacitively coupled to the resonator where the fields reach maxima for **strong interaction** g
- Qubit is **tuned** by external flux and **controlled** by gate voltage. Control detuning Δ and qubit state



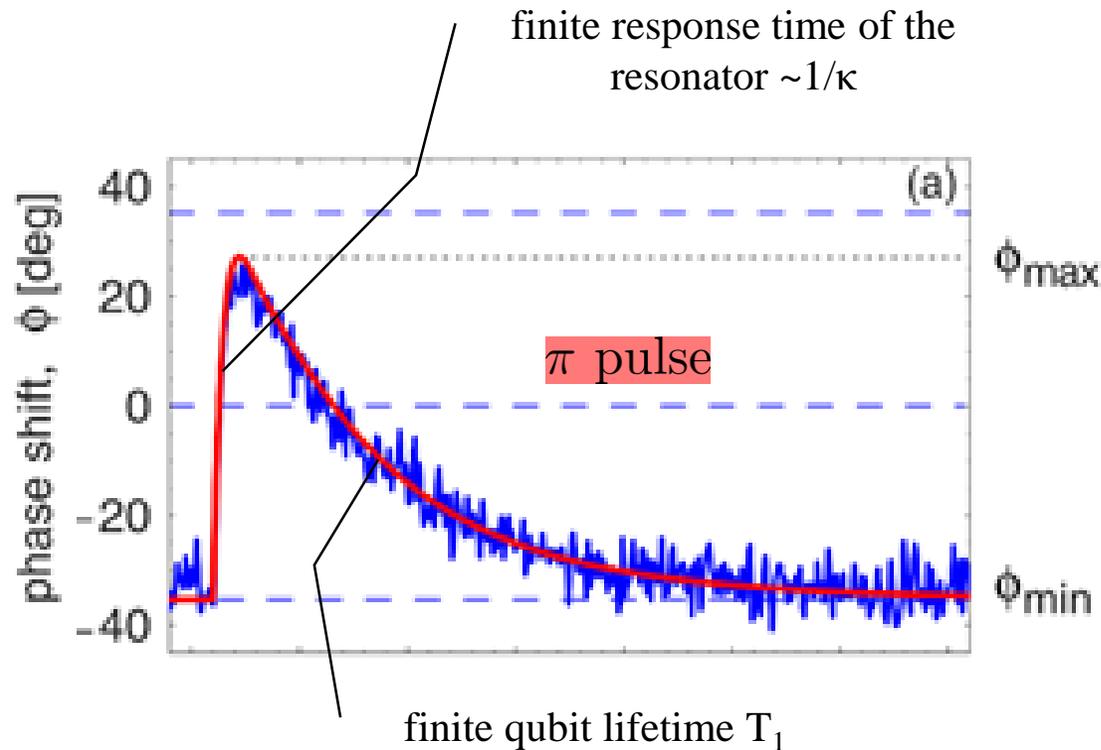
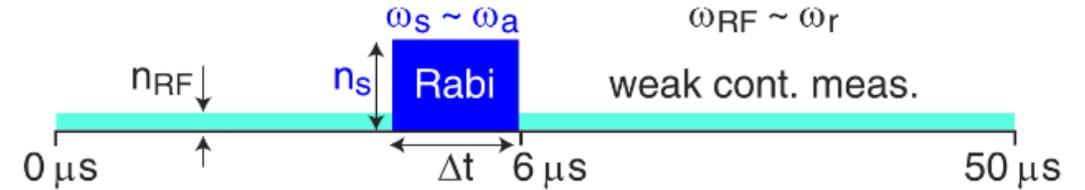
A. Blais et al., PRA 69, 062320 (2004)

L. Frunzio et al., IEEE Vol. 15, No. 2, 1051-8223 (2005)

Dispersive Readout of Qubits

Example: Rabi oscillations

- Weak continuous measurement of the phase of the transmitted wave
- Drive frequency close to resonator frequency
 $\omega_d \approx \omega_r$

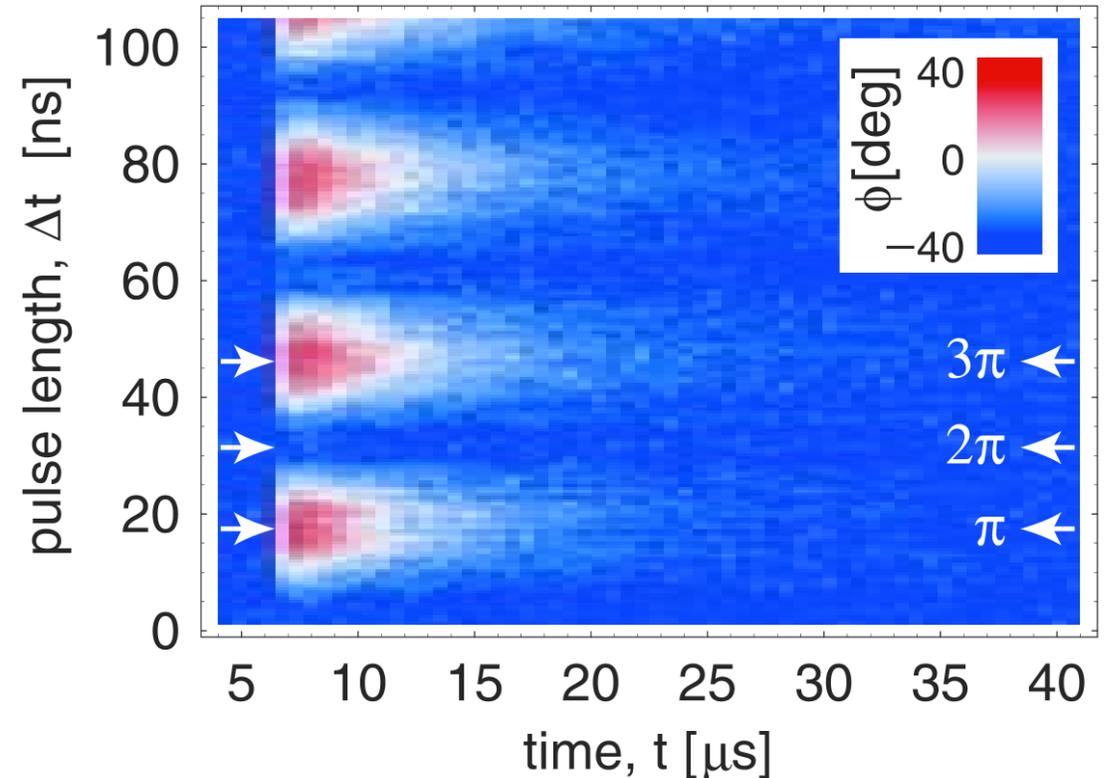


- Signal to Noise Ratio SNR=0.1
- Average over 5×10^4 times
- Projective character of the measurement is lost

Dispersive Readout of Qubits

Example: Rabi oscillations

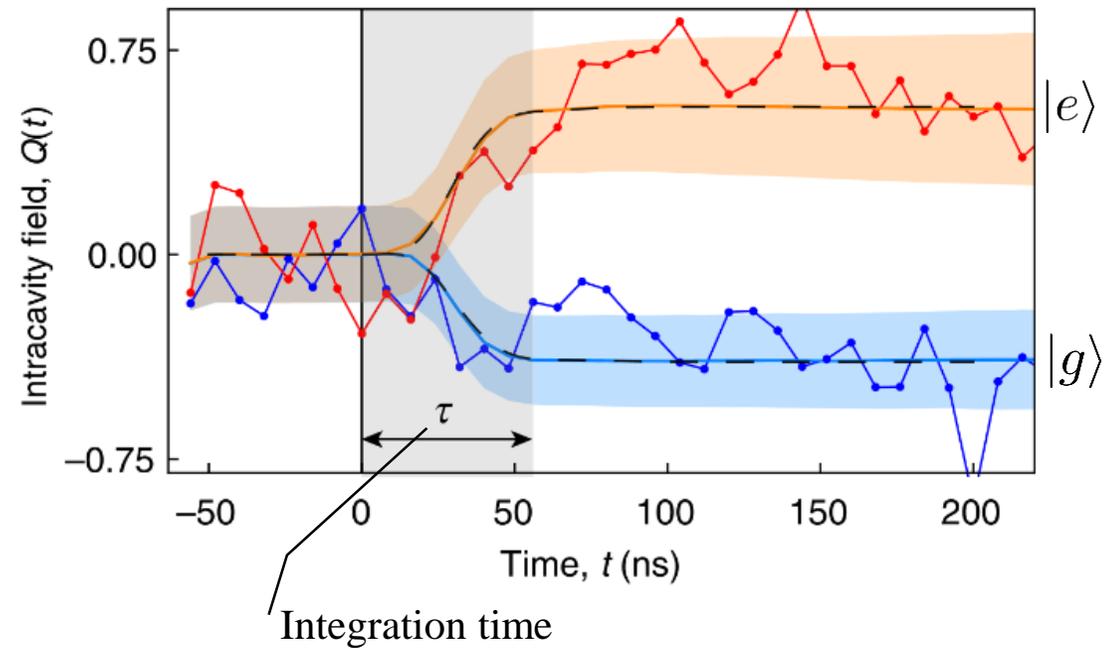
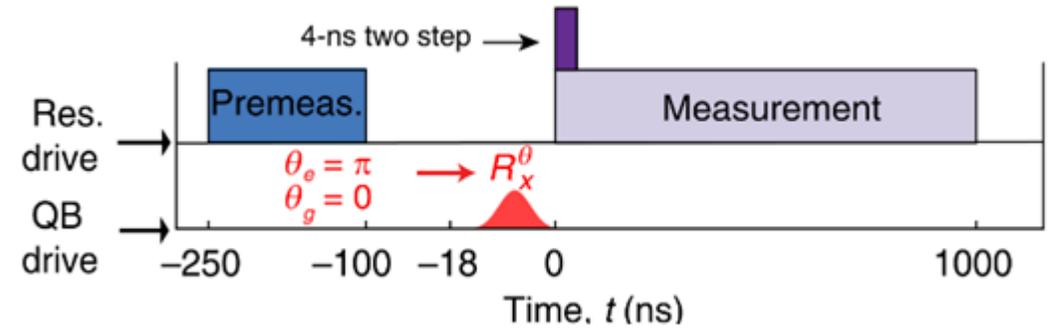
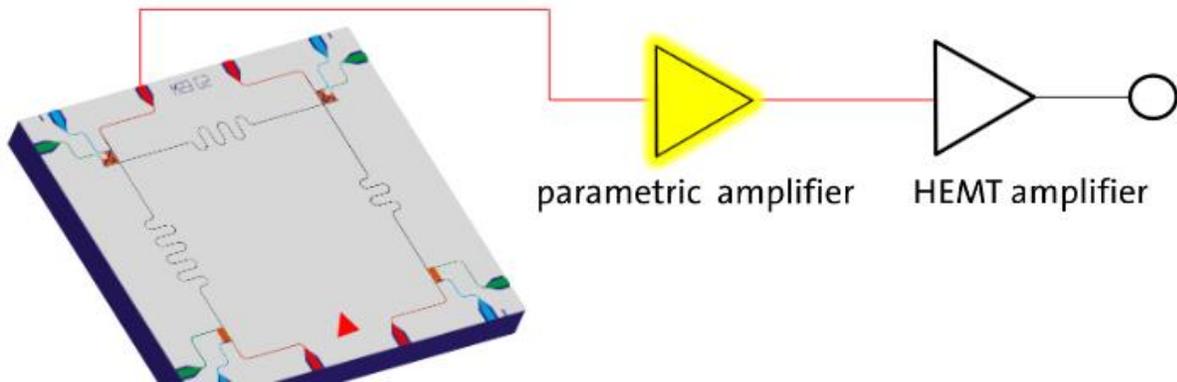
- Does the signal really correspond to the qubit state?
- Measurement duration of order T_1 : we need to average
- Why don't we just increase the number of photons?
Approximation of the Hamiltonian is no longer valid. We start driving qubit transitions!
- Ideally: **single shot** measurement
 - Reduce noise
 - Reduce measurement time
 - Increase qubit lifetime
 - Increase interaction strength



Recent progress

in dispersive qubit readout

- Additional tools (e.g. parametric amplifier, special readout tones, ...) were developed to increase **readout fidelity**, measuring **multiple qubits simultaneously** and decrease **qubit decay** during measurements
- **Single-shot dispersive readout**: One single measurement allows state determination with high fidelity
- Provides **projective measurement**



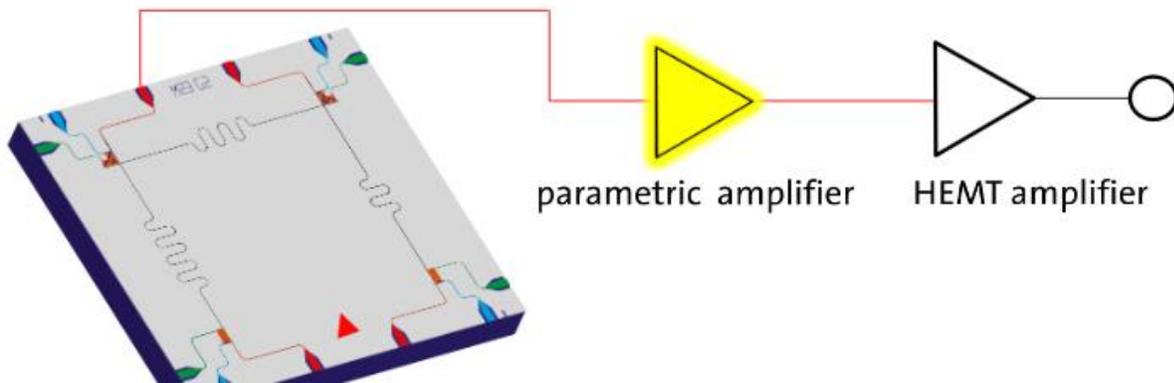
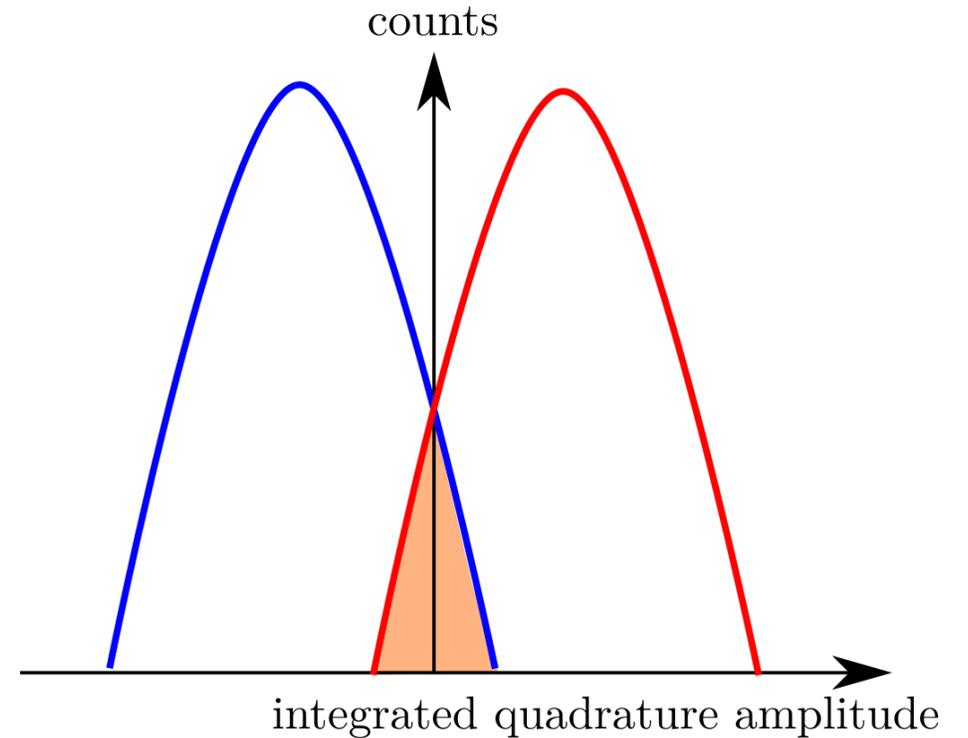
Recent progress

in dispersive qubit readout

- Extract fidelity from statistic

$$F = 1 - P(e|g) - P(g|e)$$

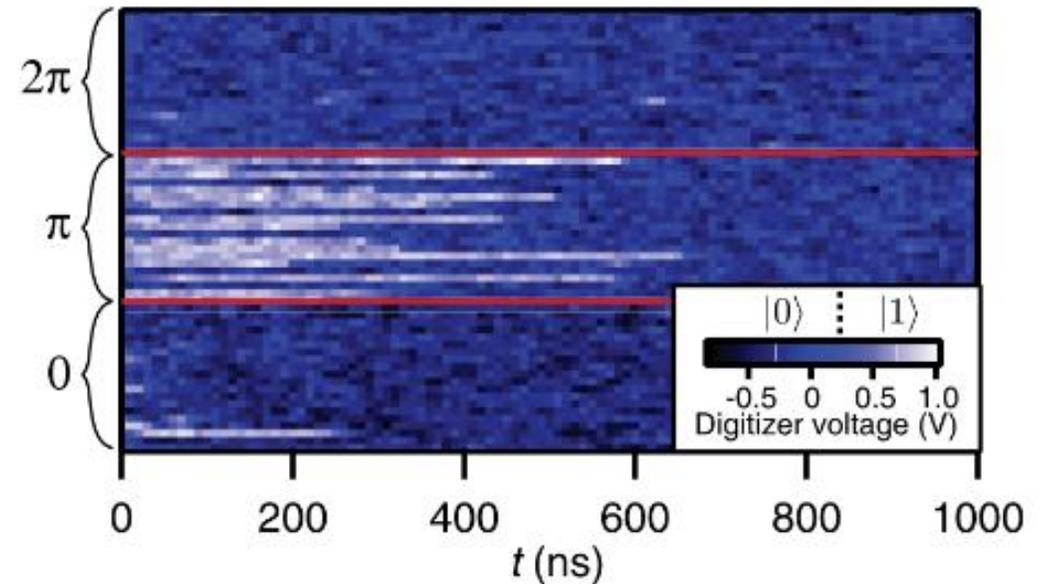
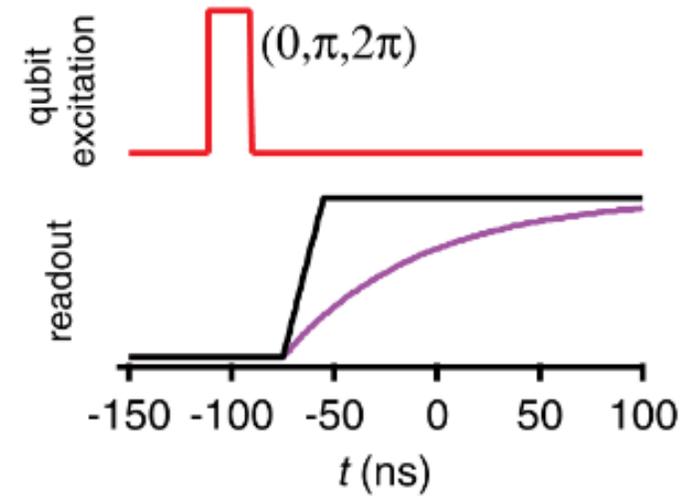
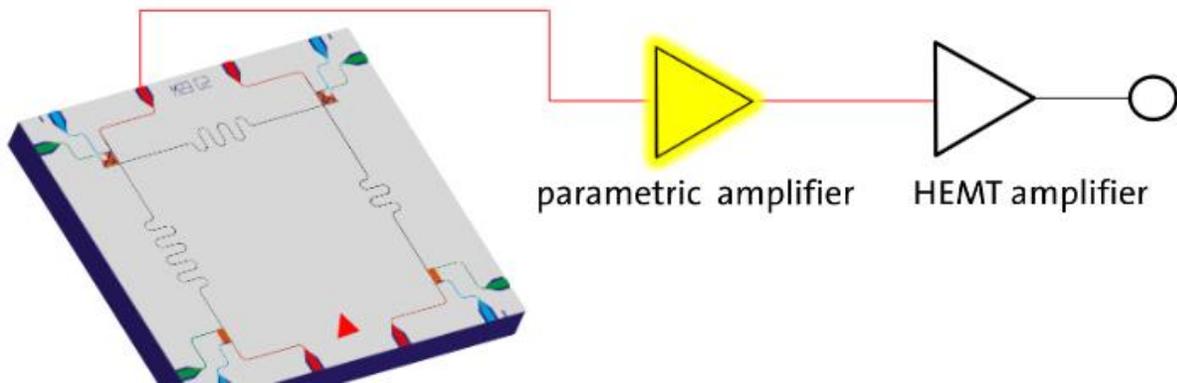
- Choose integration time properly
 - Too short: state trajectories are not clearly separated
 - Too long: spontaneous emission, thermal excitations, readout-induced transitions



Recent progress

in dispersive qubit readout

- Time scale of measurement much faster than T_1
- QND measurement scheme, allowing repetitions
- 20 single shot traces
- **Realtime** observation of **quantum jumps**
- Continuous error correction



References

- [1] A. Wallraff, D. I. Schuster, A. Blais, L. Frunzio, J. Majer, M. H. Devoret, S. M. Girvin, and R. J. Schoelkopf, *Approaching unit visibility for control of a superconducting qubit with dispersive readout*, Phys. Rev. Lett. **95**, 060501 (2005), URL <https://link.aps.org/doi/10.1103/PhysRevLett.95.060501>.
- [2] A. Blais, R.-S. Huang, A. Wallraff, S. M. Girvin, and R. J. Schoelkopf, *Cavity quantum electrodynamics for superconducting electrical circuits: An architecture for quantum computation*, Phys. Rev. A **69**, 062320 (2004), URL <https://link.aps.org/doi/10.1103/PhysRevA.69.062320>.
- [3] R. Vijay, D. H. Slichter, and I. Siddiqi, *Observation of quantum jumps in a superconducting artificial atom*, Phys. Rev. Lett. **106**, 110502 (2011), URL <https://link.aps.org/doi/10.1103/PhysRevLett.106.110502>.
- [4] T. Walter, P. Kurpiers, S. Gasparinetti, P. Magnard, A. Potočnik, Y. Salathé, M. Pechal, M. Mondal, M. Oppliger, C. Eichler, and A. Wallraff, *Rapid high-fidelity single-shot dispersive readout of superconducting qubits*, Phys. Rev. Applied **7**, 054020 (2017), URL <https://link.aps.org/doi/10.1103/PhysRevApplied.7.054020>.