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LETTER

Loophole-free Bell inequality violation using electron spins separated by 1.3 kilometres

B. Hensen^{1,2}, H. Bernien^{1,2}, A. E. Dréau^{1,2}, A. Reiserer^{1,2}, N. Kalb^{1,2}, M. S. Blok^{1,2}, J. Ruitenberg^{1,2}, R. F. L. Vermeulen^{1,2}, R. N. Schouten^{1,2}, C. Abellán³, W. Amaya³, V. Pruneri^{3,4}, M. W. Mitchell^{3,4}, M. Markham⁵, D. J. Twitchen⁵, D. Elkouss¹, S. Wehner¹, T. H. Taminiau^{1,2} & R. Hanson^{1,2}

More than 50 years ago¹, John Bell proved that no theory of nature that obeys locality and realism² can reproduce all the predictions of quantum theory: in any local-realist theory, the correlations between outcomes of measurements on distant particles satisfy an inequality that can be violated if the particles are entangled. Numerous Bell inequality tests have been reported³⁻¹³; however, all experiments reported so far required additional assumptions to obtain a contradiction with local realism resulting in

sufficiently separated such that locality prevents communication between the boxes during a trial, then the following inequality holds under local realism:

$$S = \left| \langle x \cdot y \rangle_{(0,0)} + \langle x \cdot y \rangle_{(0,1)} + \langle x \cdot y \rangle_{(1,0)} - \langle x \cdot y \rangle_{(1,1)} \right| \le 2 \qquad (1)$$

where $\langle x \cdot y \rangle_{(a,b)}$ denotes the expectation value of the product of *x* and *y* for input bits *a* and *b*. (A mathematical formulation of the concepts

Paper presentation

Which one is the Theory of Nature?

Local realism theory

- 1. Locality: Physical influences between systems can not propagate faster than light.
- 2. Realistic: The properties of a system have definite values which exists before and independend of possible measurements.

Supporter: Albert Einstein

Quantum Mechanics

1. Quantum entanglement

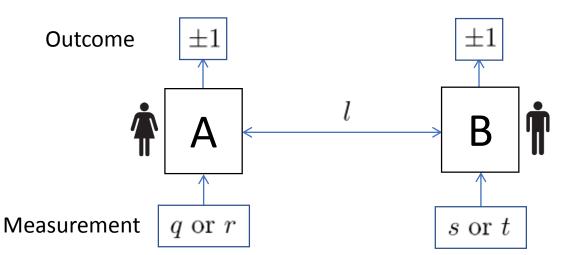
Supporter: Niels Bohr

(Bohr N. "Discussions with Einstein on Epistemological Problems in Atomic Physics")



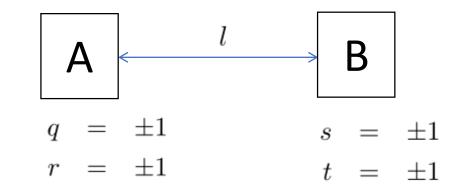
Josua Schär, Nicholas Meinhardt, Amit Patil

Bell test



- Prepared Systems A, B are separated in space one at Alice and one at Bob
- Both do a measurement on their system.
 Alice: {q,r}, Bob: {s,t} they choose randomly
- The measurement outcome can be +1 or -1 for each measurement
- Afterwards we study correlations between measurement and outcome

CHSH-Bell inequality in local realism theory



$$E\left(qs + rs + rt - qt\right) = \sum_{\substack{q,r,s,t=\pm 1 \\ q,r,s,t=\pm 1}} p\left(q,r,s,t\right) \left(qs + rs + rt - qt\right)$$

$$= \sum_{\substack{q,r,s,t=\pm 1 \\ q,r,s,t=\pm 1 \\ q,r,s,t=\pm 1 \\ q} p\left(q,r,s,t\right) \cdot 2$$

$$\leq \sum_{\substack{q,r,s,t=\pm 1 \\ q}} p\left(q,r,s,t\right) \cdot 2$$

$$= 2$$

$$E\left(qs + rs + rt - qt\right) \stackrel{\text{linearity}}{=} E\left(qs\right) + E\left(rs\right) + E\left(rt\right) - E\left(qt\right)$$
Concert from: Quantum Computation and Quantum

Concept from: Quantum Computation and Quantum Information by Michael A.Nielsen and Isaac L. Chuang

CHSH-Bell inequality in QM

$$E\left(qs + rs + rt - qt\right) \longrightarrow \langle qs + rs + rt - qt \rangle$$
$$\hat{Z}_{i} = \left|\uparrow\right\rangle \left\langle\uparrow\right|_{i} - \left|\downarrow\right\rangle \left\langle\downarrow\right|_{i}$$

 $\hat{X_i} = \left|\uparrow\right\rangle \left\langle\downarrow\right|_i + \left|\downarrow\right\rangle \left\langle\uparrow\right|_i$

$$\begin{split} \left| \Psi^{-} \right\rangle &= \frac{1}{\sqrt{2}} \left(\left| \uparrow \downarrow \right\rangle - \left| \downarrow \uparrow \right\rangle \right) \\ \left\langle qs \right\rangle &= \left\langle rs \right\rangle = \left\langle rt \right\rangle = \frac{1}{\sqrt{2}}, \quad \left\langle qt \right\rangle = -\frac{1}{\sqrt{2}} \end{split}$$

Concept from: Quantum Computation and Quantum Information by Michael A.Nielsen and Isaac L. Chuang

CHSH-Bell inequality

When we assume that the theory of nature is local-realistic

- 1. Locality: Physical influences between systems can not propagate faster than light.
- 2. Realistic: The properties of a System have definite values which exists before and independend of possible measurements.

We have the issue that QM contradicts our assumptions:

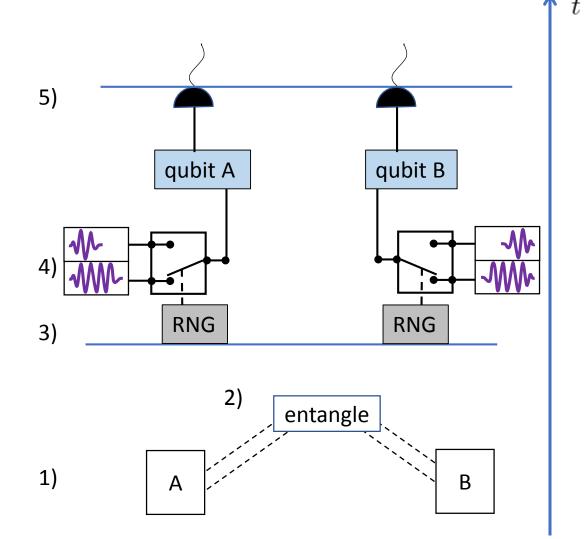
$$\frac{S_{Cl} \leq 2}{1} < S_{QM} = 2\sqrt{2}$$

If we have an experimental proof

Bell test protocol

- 1. separated 2-level systems (qubit)
- 2. Entanglement
- 3. Measurement decision
- 4. Measurement
- 5. Detection

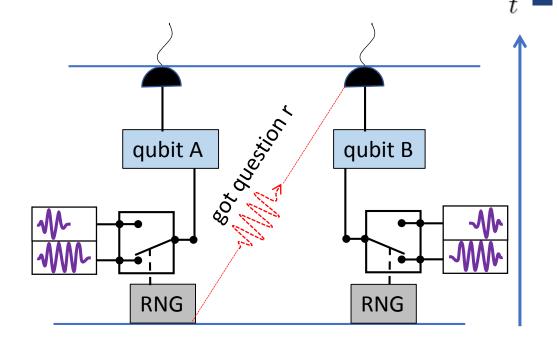
Almost ready to start but there are still some issues...



Loopholes

In an experiment different loopholes can arise.

- 1. Locality loophole: After Alice's RNG decides for a measurement axis to do a spin measurement on her qubit. A signaling photon could travel to Bob to inform his system about her question before Bob get's the measurement.
- 2. Detection loophole (fair-sampling assumption): Due to inefficient detection or loss some outcomes are not considered. We hope that the discarded outcomes have the same statistical distribution

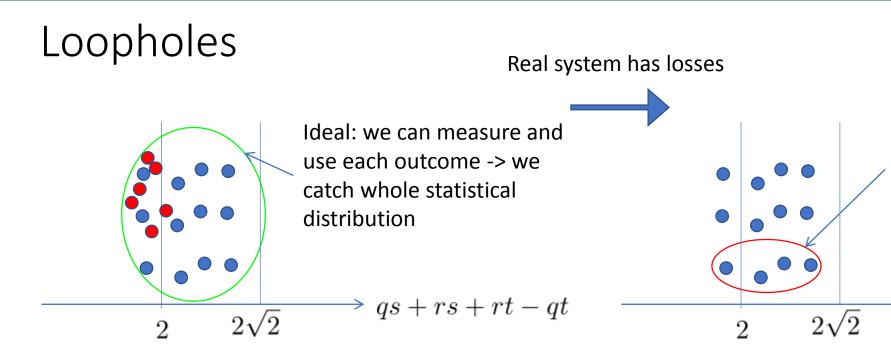


 $2\sqrt{2}$

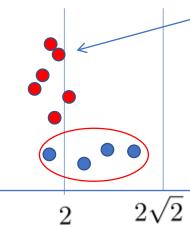
2

Assume we have a statistical distribution of 12 correlation measurement, goal is to violate 2

 $\Rightarrow qs + rs + rt - qt$



2. Detection loophole (fair-sampling assumption): Due to inefficient detection or loss some outcomes are not considered. We hope that the discarded outcomes have the same statistical distribution

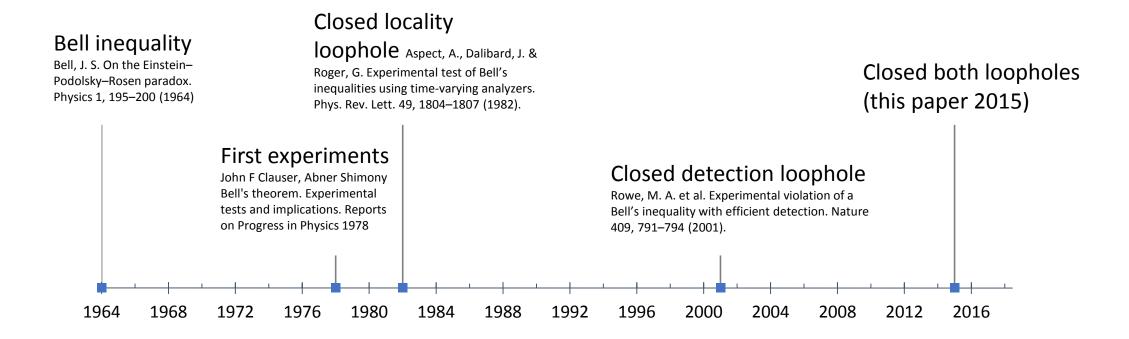


Reality: only subset of the distribution. If me make now the fair sampling assumption (not part of Bells inequality) then our subset is representative for the whole distribution

$$\Rightarrow qs + rs + rt - qt$$

Without the fair sampling assumption we can say nothing about the not considered outcomes and therefore we might not violate Bells inequality

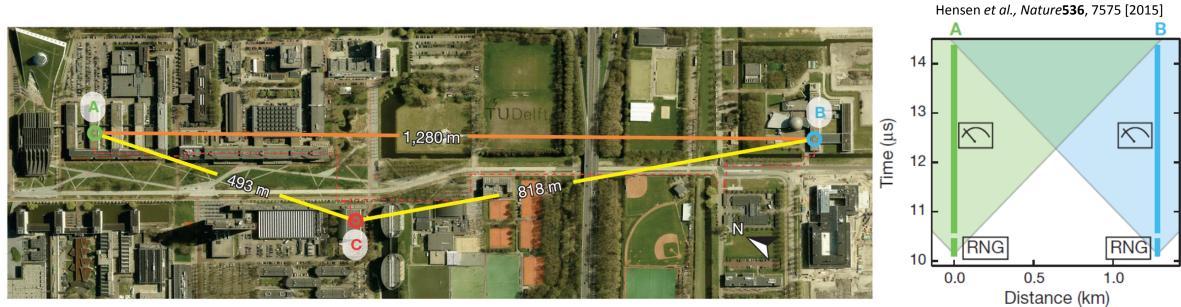
Historical overview



Setup

- Spatial separation of A and B
- Time window of 4.27 μs
- Fast operations

\Rightarrow Locality loophole closed \checkmark



Hensen et al., Nature**536**, 7575 [2015]

Implementation

- 1. Qubit: NV center
- 2. Initialize qubit
- 3. Entanglement swapping
- 4. Measurement decision
- 5. Mitigate decoherence

8

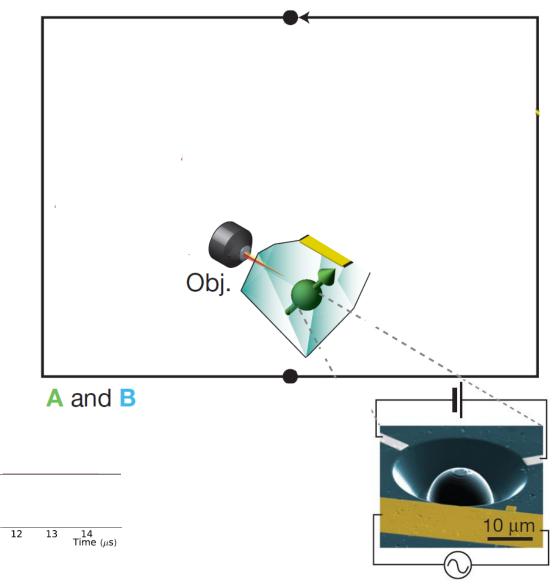
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9

10

11

6. Read-out



3

4

5

6

2

laser microwave

1

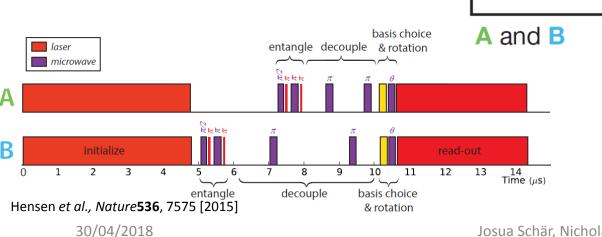
Α

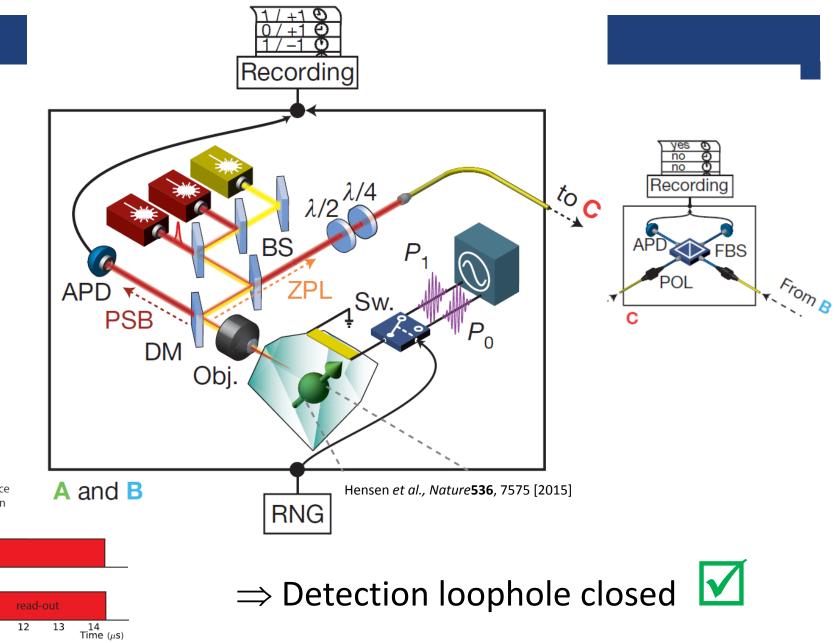
B

0

Implementation

- 1. Qubit: NV center
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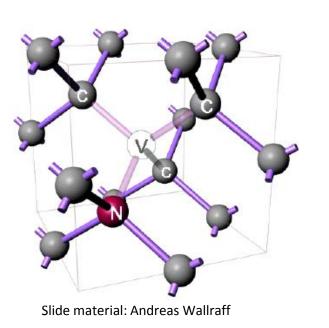




1. Qubit: NV center

CB

- Defect in diamond
- Negatively charged
- $6 e^{-}$ system $\rightarrow 2$ holes
- Total spin 1 \rightarrow triplett
- In band gap
- Like an ion trap

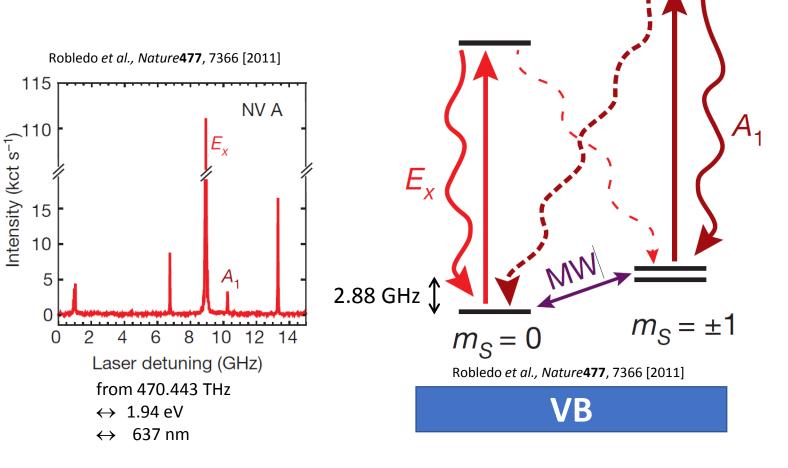


2.88 GHz
$$\updownarrow$$
 0.14 GHz \diamondsuit $m_S = \pm m_S = \pm m_S$

1. Qubit: NV center

CB

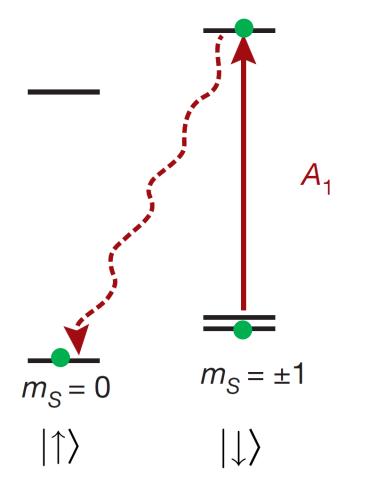
- Defect in diamond
- Negatively charged
- $6 e^{-}$ system $\rightarrow 2$ holes
- Total spin 1 \rightarrow triplett
- In band gap
- Like an ion trap
- Spin-selective optical transitions



2. Initialize qubit

- Initialize to $|\uparrow\rangle$
- Resonant A₁ excitation
- Fast decay to dark m_s = 0 state via spin mixing in excited manifold
- Fidelity: 99.8 % in 5 µs

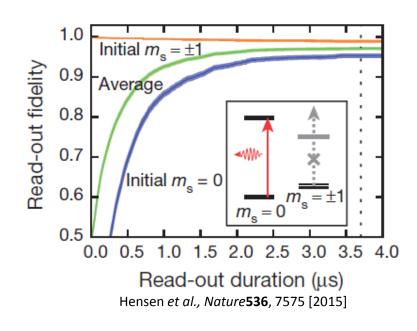
$$\begin{array}{rcl} |\uparrow\rangle &\equiv& m_{s}=0\\ |\downarrow\rangle &\equiv& m_{s}=-1 \end{array}$$

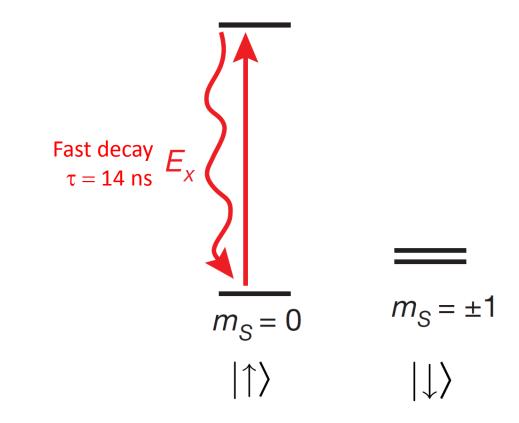


Assign:

5. Read-out

- Excite resonantly with E_x transition
- $|\uparrow\rangle \leftrightarrow$ bright, $|\downarrow\rangle \leftrightarrow$ dark
- Wait for detections during 3.7 µs
- If at least 1 photon measured $\rightarrow |\uparrow
 angle$

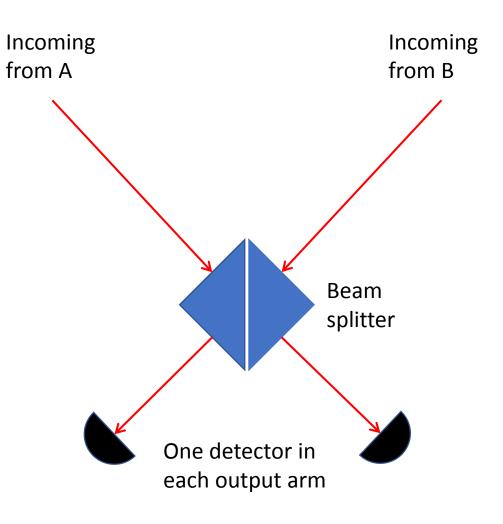




3. Entanglement swapping

- Use Hong Ou Mandel interference
- Overlap beams at beam splitter
- Photons are: indistinguishable → bunching
 (= both photons in only one output arm)

distinguishable \rightarrow coincidences (= one photon in each output arm)

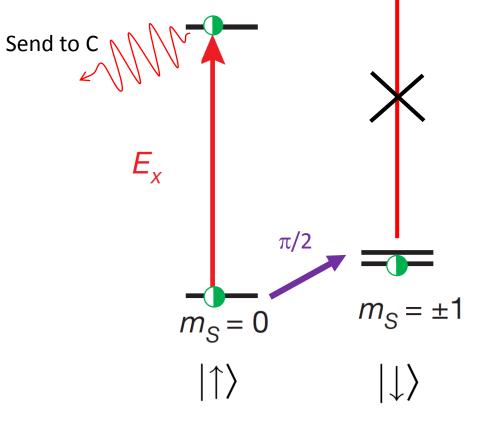


3. Entanglement swapping

- Procedure at each qubit:
- 1. Initialized to:
- 2. Excite resonantly at E_x transition \rightarrow photon and spin entangled early time bin (e)
- 3. π pulse
- 4. Again drive E_x transition \rightarrow another photon late time bin (I)

 $|\uparrow\rangle + |\downarrow\rangle$ $|\uparrow1_{e}\rangle + |\downarrow0_{e}\rangle$

 $|\downarrow> + |\uparrow>$ $|\downarrow_0|> + |\uparrow_1|>$

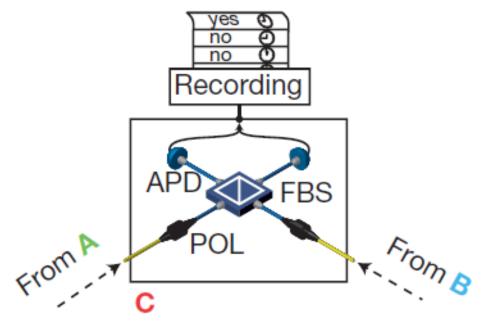


- 3. Entanglement swapping
- Measurements at C and projections ⇒ Scheme to herald successful entanglement

Early time bin From A and B: A⊗B: $|\mathbf{1}_{A}\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}\mathbf{\downarrow}_{B}\rangle|\mathbf{1}_{A}\mathbf{0}_{B}\rangle + |\mathbf{\downarrow}_{A}\mathbf{\uparrow}_{B}\rangle|\mathbf{0}_{A}\mathbf{1}_{B}\rangle + |\mathbf{\downarrow}_{A}\mathbf{\downarrow}_{B}\rangle|\mathbf{0}_{A}\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}\mathbf{\downarrow}_{B}\rangle|\mathbf{0}_{A}\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}\mathbf{\downarrow}_{B}\rangle|\mathbf{0}_{A}\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}\mathbf{\downarrow}_{B}\rangle|\mathbf{0}_{A}\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}\mathbf{\downarrow}_{B}\rangle|\mathbf{0}_{A}\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}\mathbf{\downarrow}_{B}\rangle|\mathbf{0}_{A}\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}\mathbf{\downarrow}_{B}\rangle|\mathbf{0}_{A}\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}\mathbf{1}_{B}\rangle|\mathbf{0}_{A}\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}\mathbf{1}_{B}\rangle|\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}\mathbf{1}_{B}\rangle + |\mathbf{1}_{A}$ $|\uparrow 1 > + |\downarrow 0 >$ Measuring one photon Spin flip $\mathbf{D}_{A}\mathbf{O}_{B} > + |\downarrow_{A}\uparrow_{B} > |\mathbf{O}_{A}\mathbf{1}_{B} > + |\uparrow_{A}\downarrow_{B} > |\mathbf{1}_{A}\mathbf{O}_{B} > |\mathbf{1}_{A}\mathbf{O}$ $|\downarrow 0>+|\uparrow 1>$ Late time bin (250 ns later) Measuring one photon Projection onto $|\uparrow_{a}\downarrow_{B}> \pm |\downarrow_{a}\uparrow_{B}>$ - (+): photon detection in different (same) detectors

- 3. Entanglement swapping
- Measurements at C and projections
- Robust against photon loss





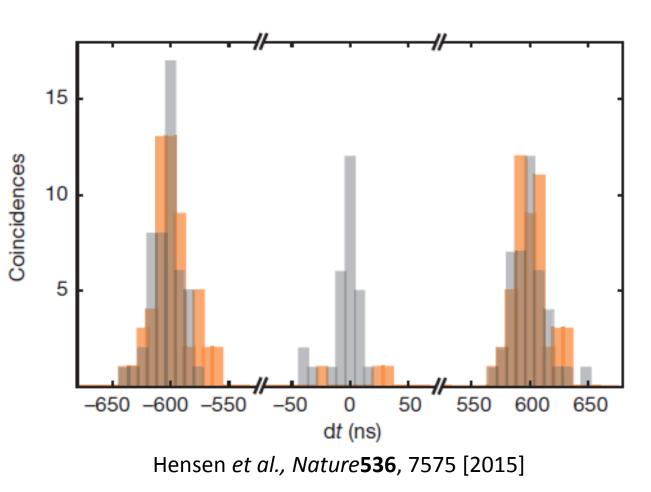
Hensen et al., Nature**536**, 7575 [2015]

Projection onto $|\uparrow_A \downarrow_B > \pm |\downarrow_A \uparrow_B >$ - (+): photon detection in different (same) detectors

Degree of indistinguishability

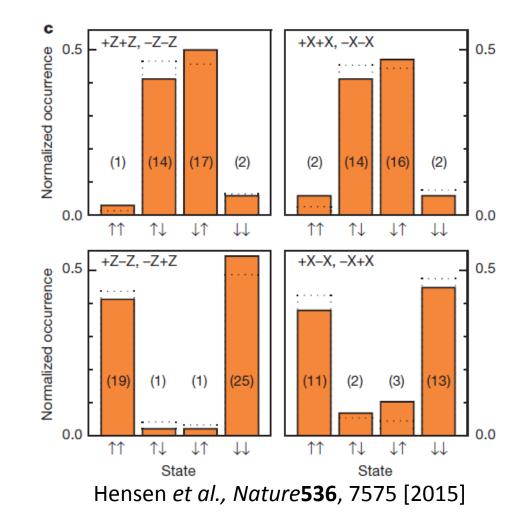
- HOM interferometry
- Dip at dt=0 as expected
- Fidelity of $|\Psi^angle$ 0.92 \pm 0.03
- Expect $S \sim 2.30 \pm 0.07$





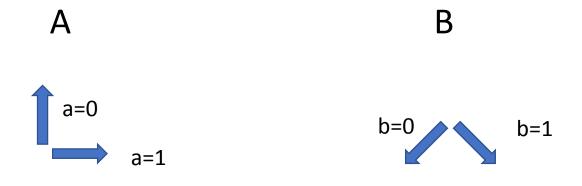
Final characterization with collinear axes

- Test performance of setup
- Randomly choose + or Z (X)
- Desired entangled state is generated $|\psi^{-}\rangle\!=\!(|\!\uparrow\!\downarrow\rangle\!-\!|\!\downarrow\!\uparrow\rangle)\big/\sqrt{2}$



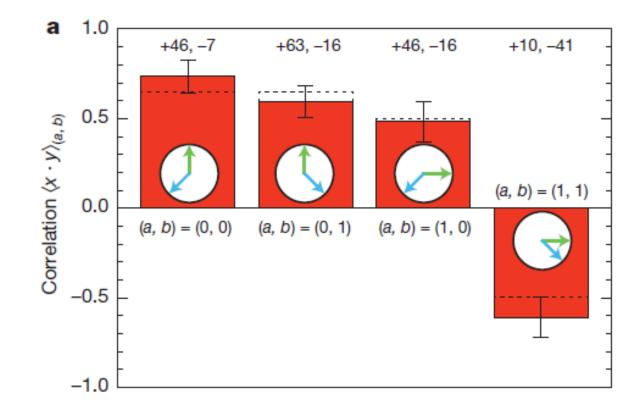
Characteristics

- Succes probability per entanglement generation attempt $\sim 6.4 \cdot 10^{-9}$
- Separation of entangled NV-centres two orders higher than before
- Opimization yields angles along z-Axis for read out bases



Results

- Only 245 successful trials during 220 h
- Arrows: green (A) and blue (B)
- $S = 2.42 \pm 0.20$
 - Ioophole-free violation of CHSH!
- P = 0.039
 - statistically significant rejection of nullhypothesis



Hensen et al., Nature536, 7575 [2015]

$$S = \left| \langle x \cdot y \rangle_{(0,0)} + \langle x \cdot y \rangle_{(0,1)} + \langle x \cdot y \rangle_{(1,0)} - \langle x \cdot y \rangle_{(1,1)} \right| \le 2$$

CHSH-Bell inequality

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We have the issue that QM contradicts our assumptions:

$$\frac{S_{Cl} \leq 2}{\swarrow} < S_{QM} = 2\sqrt{2}$$

experimental proof

Second experiment half a year later

- Modified using classical random numbers as input
- Larger time window gives better data rate

• Also $|\Psi^+>$ state, combine into single hypothesis test

Result: $S = 2.38 \pm 0.14$

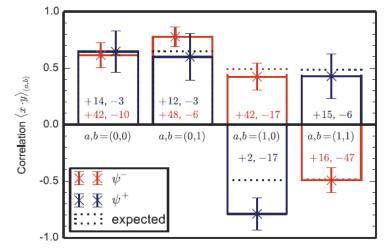


OPEN Loophole-free Bell test using electron spins in diamond: second experiment and additional analysis

Received: 11 April 2016 B. Accepted: 27 June 2016 R. Published: 11 August 2016 T.

B. Hensen^{1,2}, N. Kalb^{1,2}, M.S. Blok^{1,2}, A. E. Dréau^{1,2}, A. Reiserer^{1,2}, R. F. L. Vermeulen^{1,2}, R. N. Schouten^{1,2}, M. Markham³, D. J. Twitchen³, K. Goodenough¹, D. Elkouss¹, S. Wehner¹, T. H. Taminiau^{1,2} & R. Hanson^{1,2}

The recently reported violation of a Bell inequality using entangled electronic spins in diamonds (Hensen *et al.*, *Nature* 526, 682–686) provided the first loophole-free evidence against local-realist theories of nature. Here we report on data from a second Bell experiment using the same experimental setup with minor modifications. We find a violation of the CHSH-Bell inequality of 2.35 ± 0.18 , in agreement with the first run, yielding an overall value of $S = 2.38 \pm 0.14$. We calculate the resulting



Outlook

- Strictly speaking, Bell cannot exclude all local-realist theories due to free-will loophole
- Combination of event-ready scheme with higher entanglement rate might be used for:
- quantum key distribution
- randomness certification

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LETTER

Loophole-free Bell inequality violation using electron spins separated by 1.3 kilometres

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sufficiently separated such that locality prevents communication between the boxes during a trial, then the following inequality holds under local realism:

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where $\langle x \cdot y \rangle_{(a,b)}$ denotes the expectation value of the product of *x* and *y* for input bits *a* and *b*. (A mathematical formulation of the concepts

Thanks for your attention

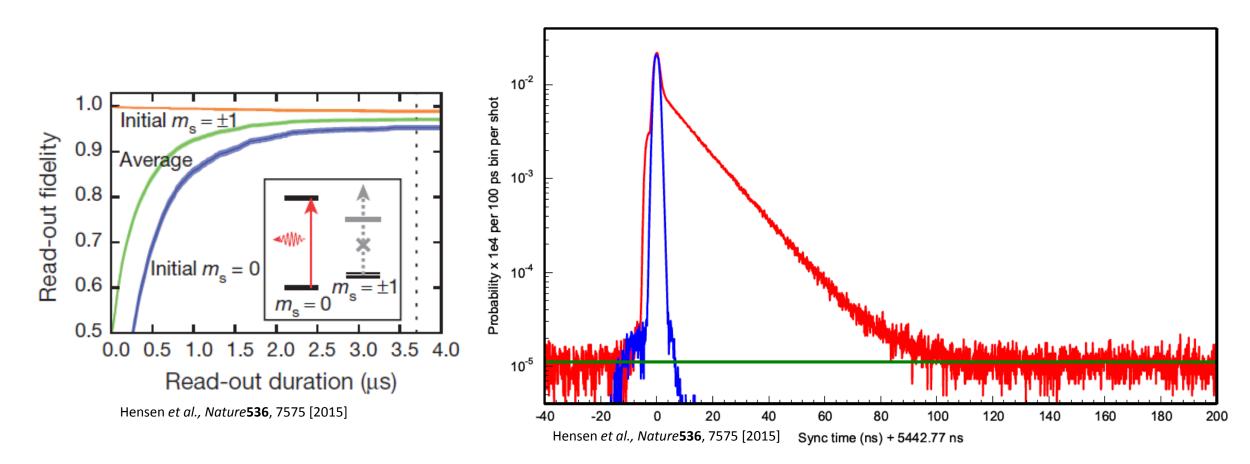
Additional slides

4. Measurement decision

- Quantum Random number generator
- Driven by process (including spontaneous emission) that is unpredictable both in quantum and classical treatment
- Excess predictability below 10^{-5}
- 160 ns

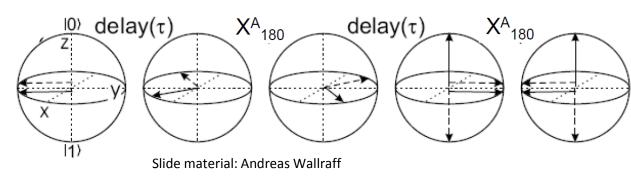
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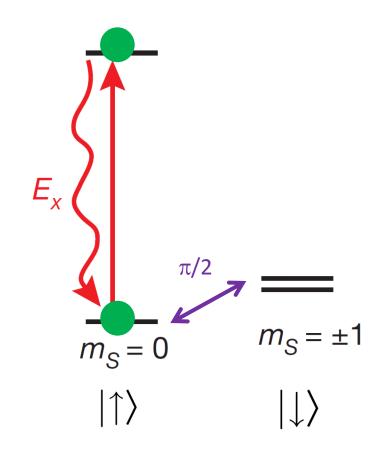
5. Read-out



6. Mitigate decoherence - refocusing

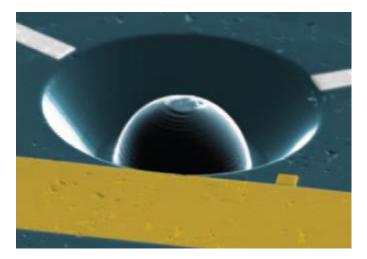
- Coherence limited by bath of ¹³C nuclear spins
- Dephasing time \sim few μ s
- Apply dynamical decoupling sequence: two MW $\pi/2$ pulses
- Probability to end up in initial state is > 99 %





7. Detection

- Photon mulitpliers
- Photon detector
- Solid immersion lens to enhance collection efficiency



Hensen et al., Nature**536**, 7575 [2015]

Statistical analysis of result

- Conventional analysis: no memory of devices, independent trials, absolute randomness of RNG → P=0.019
- Complete: arbitrary memory, partial predictability...
 → P=0.039

