

A SIX IONS



**SCHROEDINGER
CAT**

A SIX IONS SCHROEDINGER CAT

A presentation held by

Pierre Joris & Fabien Wildhaber

on the paper :

Creation of a six-atom 'Schrödinger' cat state,

D. Leibfried, E. Knill, S. Seidelin, J. Britton, R. B. Blakestad, J. Chiaverini, D. B. Hume, W. M. Itano, J. D. Jost, C. Langer, R. Ozeri, R. Reichle & D. J. Wineland
(Boulder group)

Nature **438**, 639-642 (2005)

Outline

- Implementation of a Six-ions Schrödinger Cat
 - Hyperfine structure
 - Trap
 - Laser cooling
 - Schrödinger cat
- Some Theory...
 - Criteria for complete entanglement
- Measurements and Results
 - Application of the criteria to measured data

Implementation of a Six-ions Schrödinger Cat

Implementation of a Six-ions Schrödinger Cat

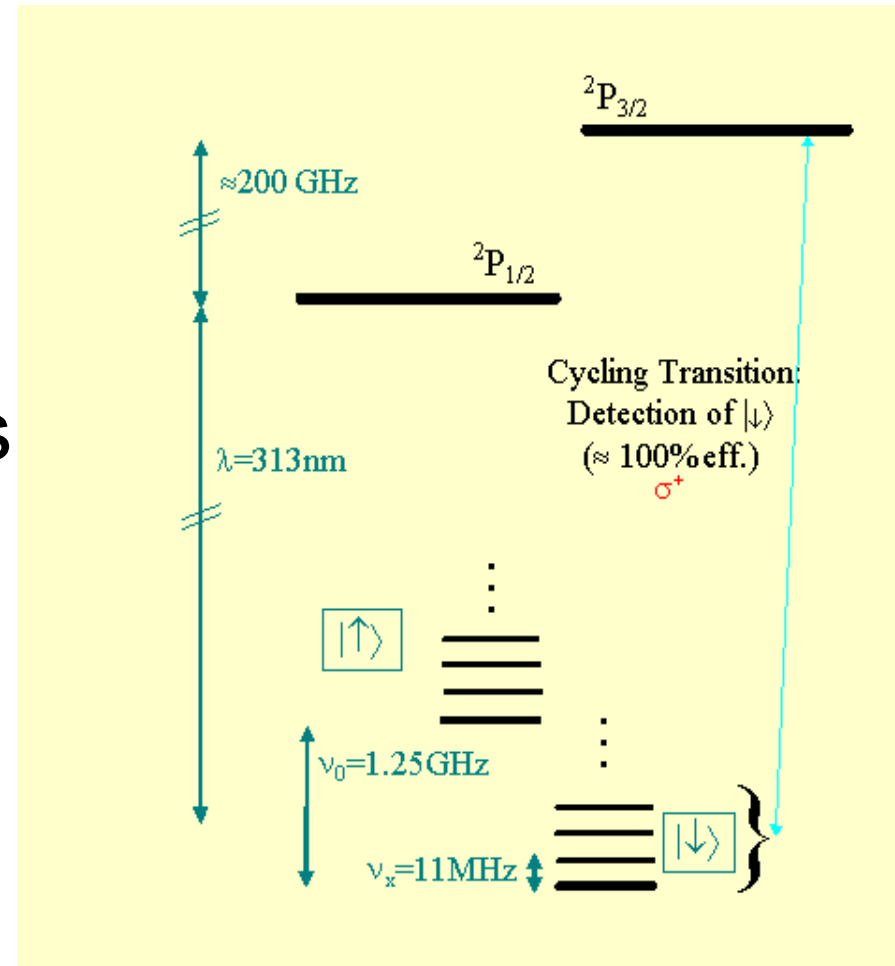
- Qubits = Hyperfine levels of ${}^9\text{Be}^+$ corresponding to $|\uparrow\rangle$ and $|\downarrow\rangle$
- Paul trap : ions trapped in a RF-field
- Cooling and initialization of the 'cat state' using laser pulses

Hyperfine Structure : from QM

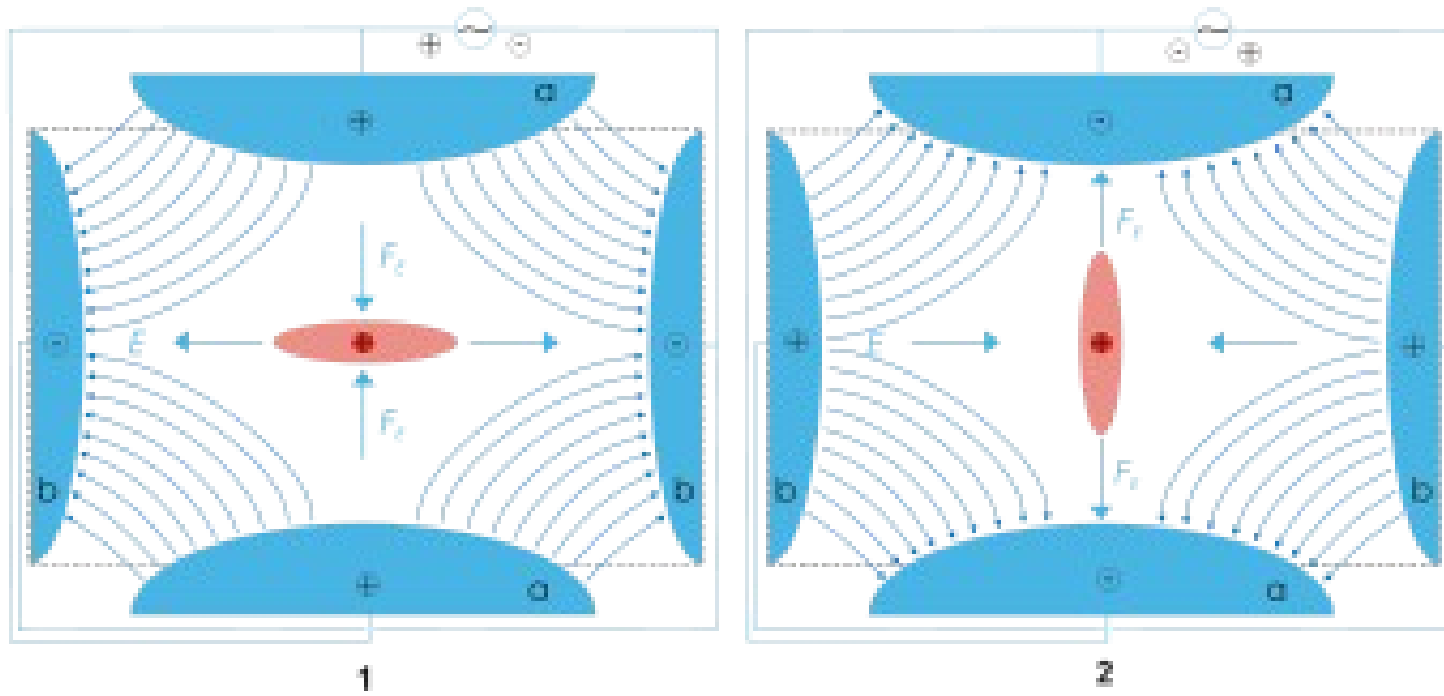
- Interaction between nuclear spin and dipolar moment of the electron
- F = total angular momentum
 m_F = projection on the z-axis
- Qubits states:

$$|\downarrow\rangle = |F = 2, m_F = -2\rangle$$

$$|\uparrow\rangle = |F = 1, m_F = -1\rangle$$



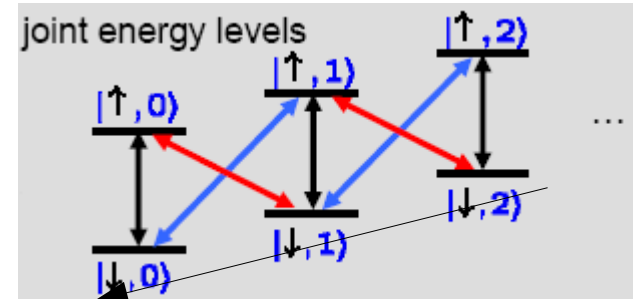
Paul Trap



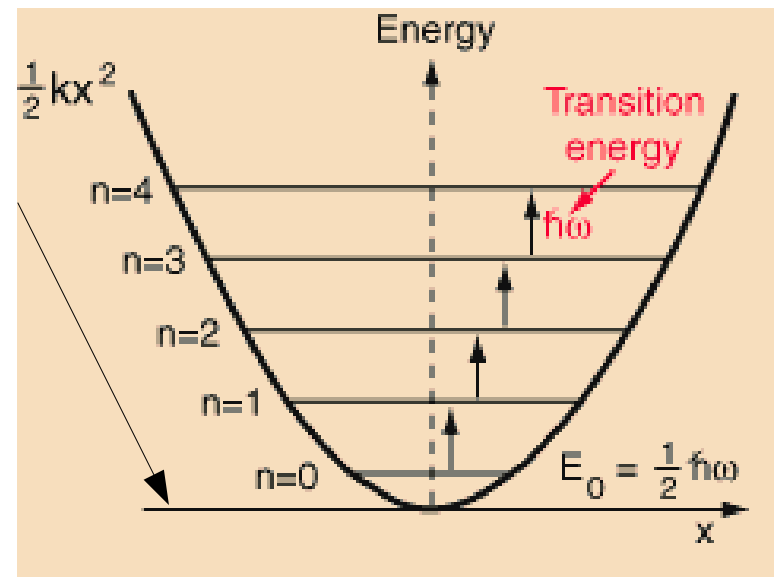
- Oscillating RF Field
=> harmonic oscillations of the ions
=> Confinement

Laser Cooling

- 'Red Sideband' transition stimulated by laser:
... $2 \rightarrow 1 \rightarrow 0$



- Decay by spontaneous emission: $\uparrow \rightarrow \downarrow$
- Each ion tends to the ground state of the harmonic oscillator $|\downarrow, 0\rangle$



- Notation for N ions: $|\uparrow\downarrow\dots\uparrow\rangle$
for the ground state: $|\downarrow\downarrow\dots\downarrow\rangle \equiv |N, \downarrow\rangle$

Schrödinger Cat States

- Schrödinger cat state =
Equal superposition of 2 **maximally different**
quantum states

$$|\Psi\rangle = \frac{1}{\sqrt{2}} ($$



+



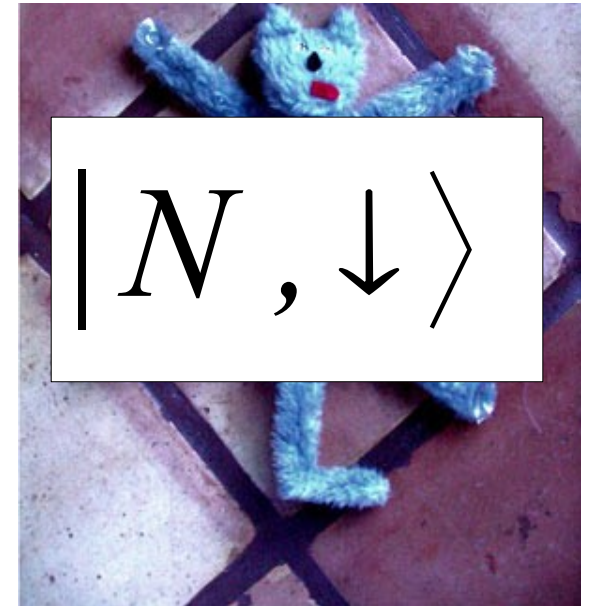
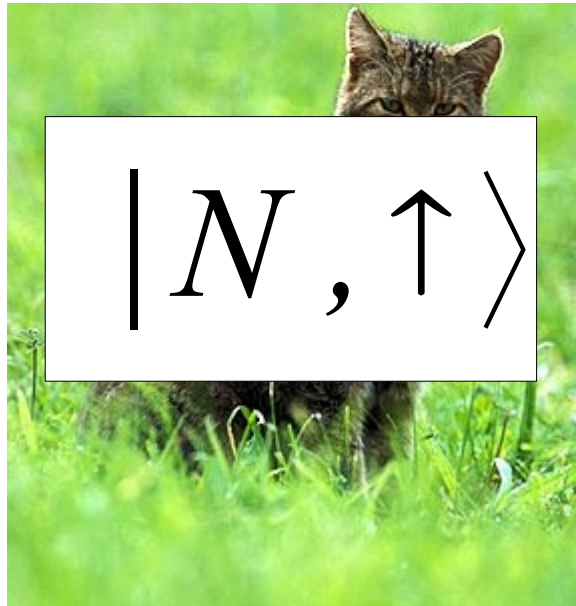
)

Schrödinger Cat States (2)

- Generated by rotation in the phase space of the ground state $|N, \downarrow\rangle = |\downarrow, \downarrow, \dots, \downarrow\rangle$
- Implemented by action of two detuned laser beams stimulating a Raman transition (phase gate U_N)

$$|N \text{ Cat}\rangle = U_N |N, \downarrow\rangle =$$

$$= \frac{1}{\sqrt{2}} \left(\begin{array}{c} |N, \uparrow\rangle \\ \text{[Cat Image]} \end{array} + \begin{array}{c} |N, \downarrow\rangle \\ \text{[Toilet Paper Roll Image]} \end{array} \right)$$



Some
Theory...

Some Theory ...

- Proofs of complete entanglement
 - Fidelity
 - Witness operator
 - Condition over density matrix



Proofs of entangled state :

Fidelity

- Fidelity F is the simplest measure of how close the generated state $|\Psi_N\rangle$ is to $|N\text{Cat}\rangle$

$$F_{N\text{Cat}} \equiv |\langle \Psi_N | N\text{Cat} \rangle|^2$$

if maximally entangled cat state $F_{N\text{Cat}} = 1$

- Not sufficient to fully characterize entanglement

Proofs of entangled state : The Witness Operators

- Type of operators, for which
 $\langle W \rangle < 0 \Rightarrow$ complete entanglement (i.e.
each ion entangled with all others)
- In this paper we use

$$W \equiv 1 - 2 |NCat\rangle\langle NCat| \quad (1)$$

that has the important property :

$$\langle W \rangle \leq 1 - 4 |C_{\downarrow N; \uparrow N}| \quad (2)$$

where $C_{\downarrow N; \uparrow N} = \langle \Psi_N | \downarrow, N \rangle \langle \uparrow, N | \Psi_N \rangle$

is one component of the density matrix ($2^6 \times 2^6$)

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Consequence of (1) and (2) :

$$\boxed{|C_{\downarrow N; \uparrow N}| > 0.25 \Rightarrow \text{complete entanglement}} \quad (A)$$

Proofs of entangled state: condition over density matrix



- Another condition for complete entanglement:

$$2|C_{\downarrow N; \uparrow N}| > \max(P_j + P_{\bar{j}}) \Rightarrow \text{compl. entanglement} \quad (\text{B})$$

where

$P_j = |\langle \Psi_N | j \rangle|^2$ is the probability to measure $|j\rangle$,

with $|j\rangle$ a basis vector different from $|N, \uparrow\rangle, |N, \downarrow\rangle$

i.e. $|j\rangle = |\downarrow\downarrow\downarrow\downarrow\downarrow\uparrow\rangle, \dots, |\downarrow\downarrow\uparrow\downarrow\downarrow\uparrow\rangle, \dots, |\uparrow\uparrow\uparrow\uparrow\uparrow\downarrow\rangle$
 $|\bar{j}\rangle = |\uparrow\uparrow\uparrow\uparrow\uparrow\downarrow\rangle, \dots, |\uparrow\uparrow\downarrow\uparrow\uparrow\downarrow\rangle, \dots, |\downarrow\downarrow\downarrow\downarrow\downarrow\uparrow\rangle$

Measurements and Results

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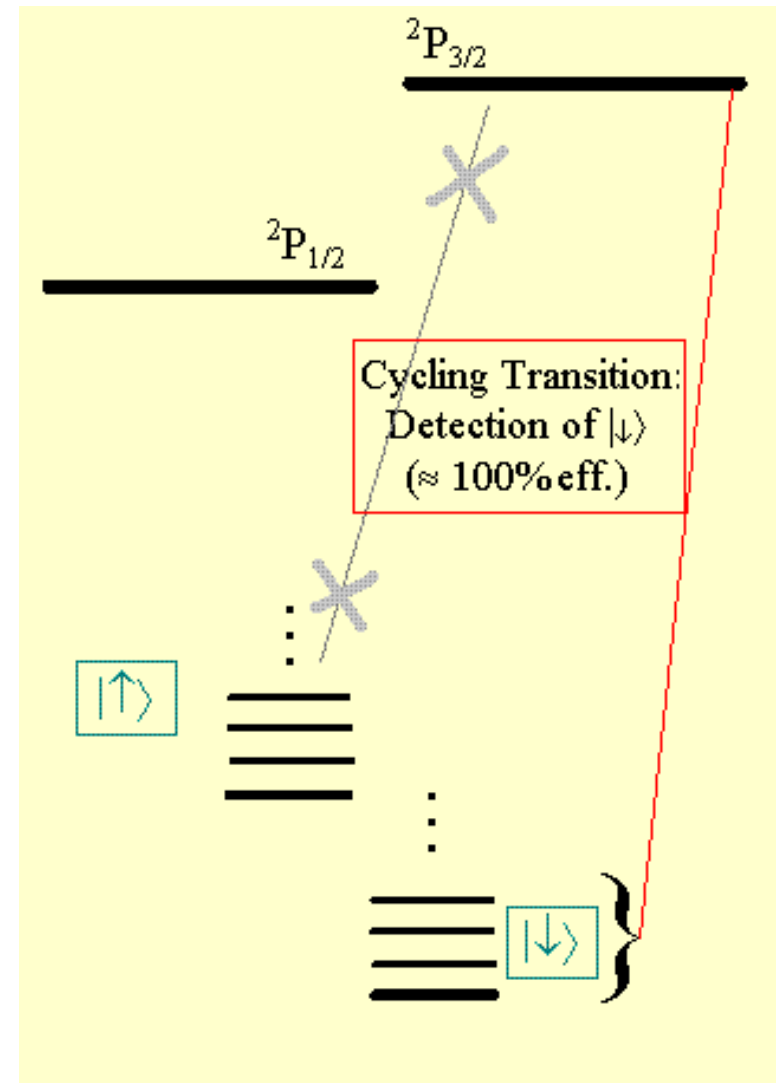
- Determination of population by fluorescence
- Calibration
- Determination of the density matrix components $C_{\downarrow N; \uparrow N}$
- Proof of the complete entanglement of the generated states



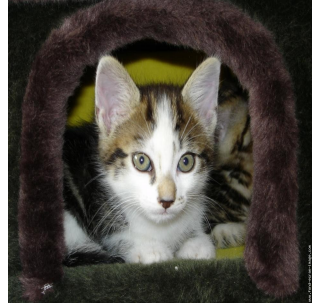
Determination of population by fluorescence and calibration

- Main observable to which we have access is the fluorescence of the down states (no fluorescence in the up state)
- Calibration carried out by statistical interpretation of fluorescence measurement of the known states

$$|N, \downarrow\rangle, N=0, \dots, 6$$



Determination of the $C_{\downarrow N; \uparrow N}$'s



- We need the $C_{\downarrow N; \uparrow N}$ to verify entanglement conditions A and B
- By use of 'decoding' phase gate $U_{N, \Phi}$ (implemented by laser pulses) on the generated state $|\Psi_N\rangle$

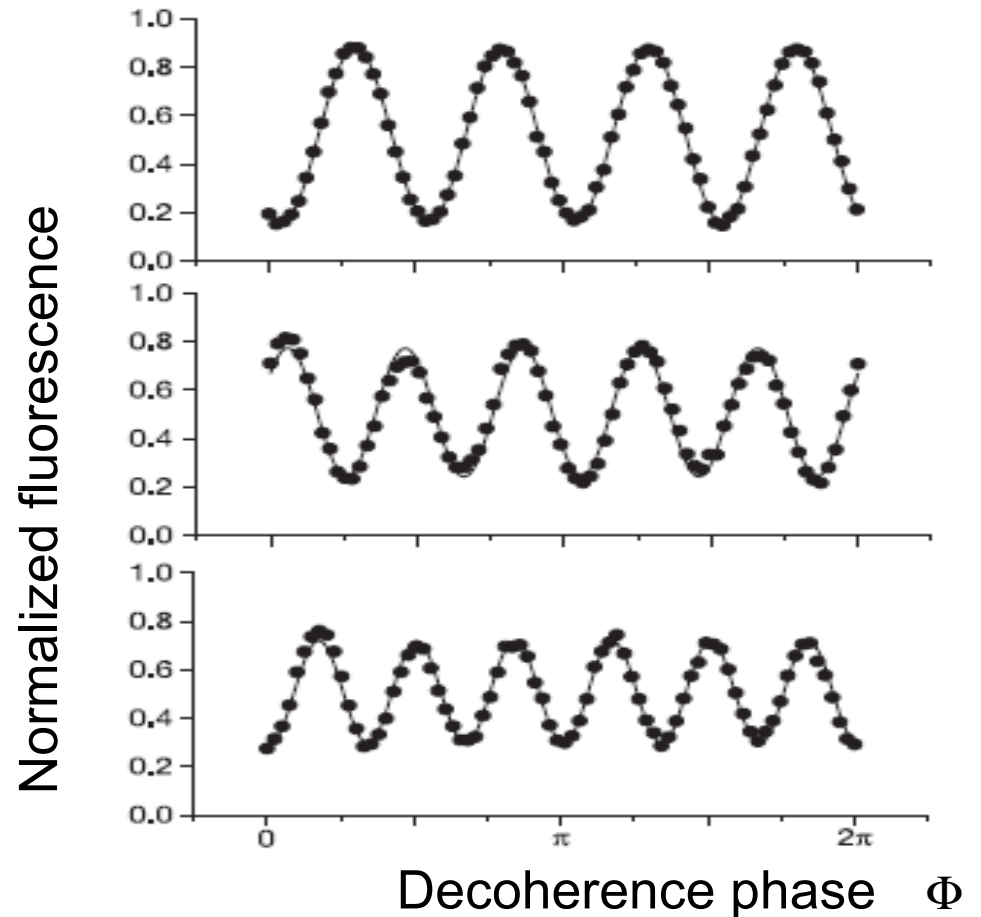
we get a new **decoded** state:

$$\begin{aligned} |\Psi_\Phi\rangle &= U_{N, \Phi} |N \text{ Cat}\rangle \\ &= -i \sin\left(\frac{N}{2} \Phi\right) |\downarrow, N\rangle + i^{N+1+\xi} \cos\left(\frac{N}{2} \Phi\right) |\uparrow, N\rangle \end{aligned}$$

Determination of the $C_{\downarrow N; \uparrow N}$'s (2)

- Fluorescence measurement over decoded states

=> Sinusoidal oscillations as function of ϕ (the parameter of $U_{N, \phi}$)



Determination of the $C_{\downarrow N; \uparrow N}$'s (3)

- We have

$$C_{\downarrow N; \uparrow N} \geq \frac{1}{2} A / A_{max}$$

A_{max} : Amplitude in perfect
cat state case

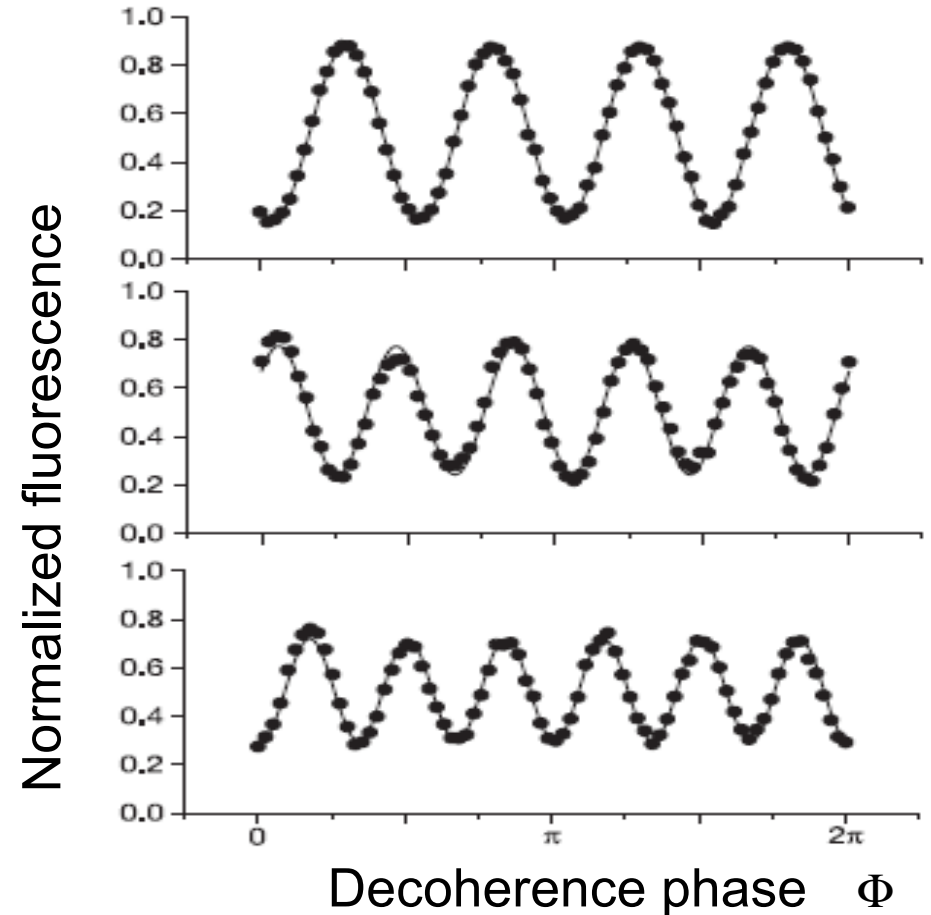
A : Effective amplitude

- Experimentally:

$$|C_{\downarrow 4; \uparrow 4}| \geq 0.349$$

$$|C_{\downarrow 5; \uparrow 5}| \geq 0.264$$

$$|C_{\downarrow 6; \uparrow 6}| \geq 0.210$$



Proof of entanglement (N=4,5)

- Recalling condition (A)

$$|C_{\downarrow N; \uparrow N}| > 0.25 \quad \Rightarrow \text{compl. entanglement}$$

- And since experimentally we got :
 $|C_{\downarrow 4; \uparrow 4}| \geq 0.349$
 $|C_{\downarrow 5; \uparrow 5}| \geq 0.264$
 $|C_{\downarrow 6; \uparrow 6}| \geq 0.210$

we conclude that the states with N=4,5 are **definitely completely entangled**
(for N=6 it is not sure at this point)

Proof of entanglement for N=6

- We rewrite now condition B:

$$|C_{\downarrow N; \uparrow N}| > \max(P_j) \Rightarrow \text{compl. entanglement}$$

where the P_j are determined by fluorescence measurements over the state $|\Psi_N\rangle$

- Experimentally we have for N=6:

$$|C_{\downarrow 6; \uparrow 6}| = 0.210 > \max(P_j) = 0.119$$

- Conclusion : our $|\Psi_6\rangle$ state is **proved to be entangled too**

References

- D. Leibfried *et al.*, *Creation of a six-atom 'Schrödinger' cat state*, Nature **438**, 639-642 (2005)
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QUESTIONS

