2.6 Quantum Measurement

Quantum measurement is done by having a closed quantum system interact in a controlled way with an external system from which the state of the quantum system under measurement can be recovered.

- example to be discussed: dispersive measurement in cavity QED

2.6.1 The quantum measurement postulate

QM postulate: **quantum measurement** is described by a set of operators **$\{M_m\}$** acting on the state space of the system. The **probability p of a measurement result m** occurring when the state ψ is measured is

$$p(m) = \langle \psi | M_m^\dagger M_m | \psi \rangle$$

the **state of the system after the measurement** is

$$|\psi'\rangle = \frac{M_m |\psi\rangle}{\sqrt{p(m)}}$$

completeness: the sum over all measurement outcomes has to be unity

$$1 = \sum_m p(m) = \sum_m \langle \psi | M_m^\dagger M_m | \psi \rangle$$

2.6.2 Example: projective measurement of a qubit in state ψ in its computational basis

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$$

measurement operators:

$$M_0 = |0\rangle\langle 0| = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \; ; \quad M_1 = |1\rangle\langle 1| = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

measurement probabilities:

$$p(0) = \langle\psi|M_0^\dagger M_0|\psi\rangle = \alpha^*\alpha\langle 0|0\rangle = |\alpha|^2$$

$$p(1) = \langle\psi|M_1^\dagger M_1|\psi\rangle = \beta^*\beta\langle 1|1\rangle = |\beta|^2$$

state after measurement:

$$\frac{M_0|\psi\rangle}{\sqrt{p(0)}} = \frac{\alpha|0\rangle}{\sqrt{|\alpha|^2}} = \frac{\alpha}{|\alpha|}|0\rangle$$

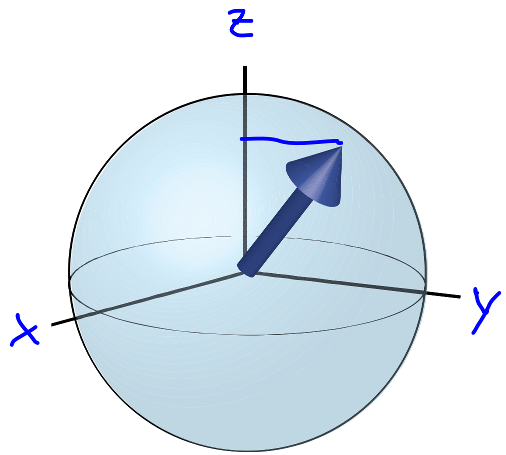
$$\frac{M_1|\psi\rangle}{\sqrt{p(1)}} = \frac{\beta|1\rangle}{\sqrt{|\beta|^2}} = \frac{\alpha}{|\alpha|}|1\rangle$$

measuring the state again after a first measurement yields the same state as the initial measurement with unit probability

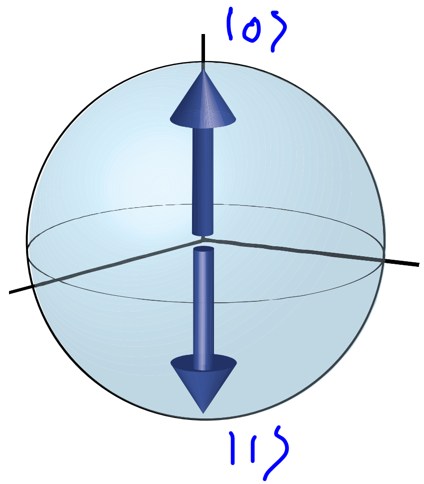
2.6.3 Interpretation of the Action of a Projective Measurement

One way to determine the state of a qubit is to measure the projection of its state vector along a given axis, say the z-axis.

On the Bloch sphere this corresponds to the following operation:



After a projective measurement is completed the qubit will be in either one of its computational basis states.



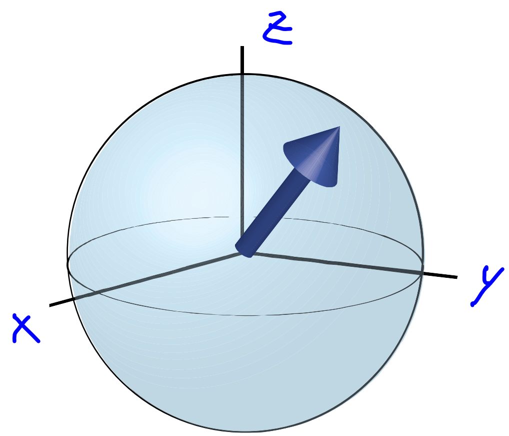
In a repeated measurement the projected state will be measured with certainty.

Information content in a single qubit state

- infinite number of qubit states

- but single measurement reveals only 0 or 1 with probabilities $|\alpha|^2$ or $|\beta|^2$
- measurement will collapse state vector on basis state
- to determine α and β an infinite number of measurements has to be made

But if not measured the qubit contains 'hidden' information about α and β.



2.7 Multiple Qubits

2.7.1 Two Qubits

2 classical bits with states:

bit 1	bit 2
0	0
0	1
1	0
1	1

- 2^n different states (here n=2)
- but only one is realized at any given time

2 qubits with quantum states:

qubit 1	qubit 2

$|00\rangle$
$|01\rangle$
$|10\rangle$
$|11\rangle$

- 2^n basis states (n=2)
- can be realized simultaneously
- quantum parallelism

2^n complex coefficients describe quantum state

$$|\psi\rangle = \alpha_{00}|00\rangle + \alpha_{01}|01\rangle + \alpha_{10}|10\rangle + \alpha_{11}|11\rangle$$

normalization condition

$$\sum_{ij} |\alpha_{ij}|^2 = 1$$

2.7.2 Composite quantum systems

QM postulate: The state space of a composite systems is the tensor product of the state spaces of the component physical systems. If the component systems have states ψ_i the composite system state is

$$|\psi\rangle = |\psi_1\rangle \otimes |\psi_2\rangle \otimes \dots \otimes |\psi_n\rangle$$

This is a product state of the individual systems.

example:

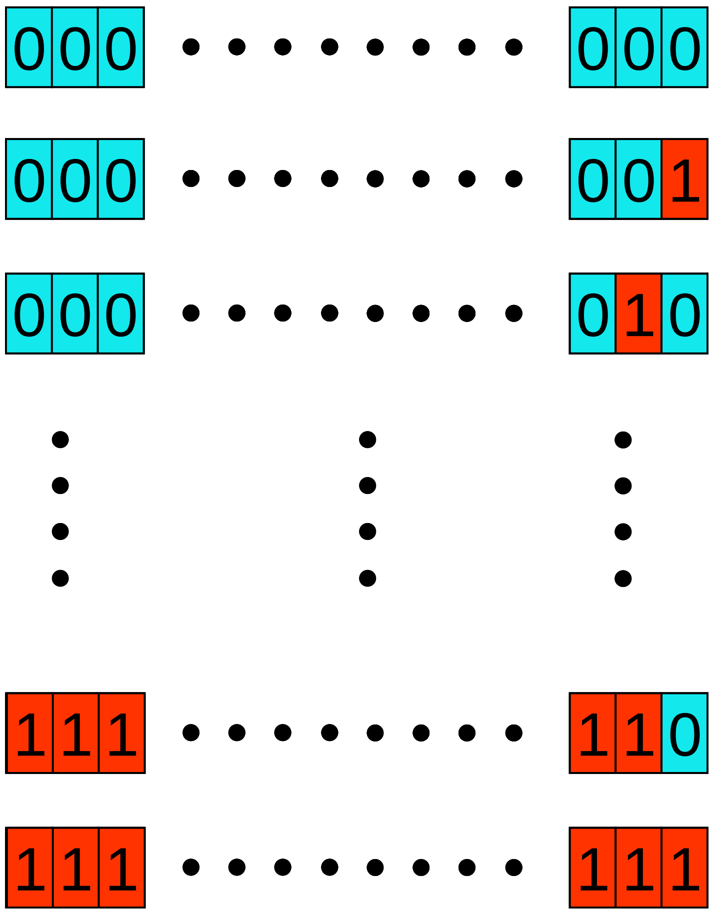
$$|\psi_1\rangle = \alpha_1|0\rangle + \beta_1|1\rangle$$
$$|\psi_2\rangle = \alpha_2|0\rangle + \beta_2|1\rangle$$

$$\rightarrow \quad |\psi\rangle = |\psi_1\rangle \otimes |\psi_2\rangle = |\psi_1 \psi_2\rangle$$
$$= \alpha_1\alpha_2|00\rangle + \alpha_1\beta_2|01\rangle + \beta_1\alpha_2|10\rangle + \beta_1\beta_2|11\rangle$$

exercise: Write down the state vector (matrix representation) of two qubits, i.e. the tensor product, in the computational basis. Write down the basis vectors of the composite system.

there is no generalization of Bloch sphere picture to many qubits

2.7.3 A register of N quantum bits



classical register:

- has 2^N possible configurations
- but can store only 1 number

quantum register:

- has 2^N possible basis states
- can store superpositions of all numbers simultaneously

Goal: Try to process superposition of numbers simultaneously in a quantum computer.

- But what is needed to construct a quantum computer and how would it be operated?

2.7.4 Information content in multiple qubits

- 2^n complex coefficients describe the state of a composite quantum system with n qubits

- Imagine to have 500 qubits, then 2^{500} complex coefficients describe their state.

- How to store this state?
    - 2^{500} is larger than the number of atoms in the universe.
    - It is impossible in classical bits.
    - This is also why it is hard to simulate quantum systems on classical computers.

- A quantum computer would be much more efficient than a classical computer at simulating quantum systems.

- Make use of the information that can be stored in qubits for quantum information processing!

2.7.5 Entanglement

Definition: An **entangled state** of a composite system is a state that cannot be written as a product state of the component systems.

example: an entangled 2-qubit state (one of the Bell states)

$$|\psi\rangle = \frac{1}{\sqrt{2}}\left(|00\rangle + |11\rangle\right)$$

What is special about this state? Try to write it as a product state!

$$|\psi_1\rangle = \alpha_1|0\rangle + \beta_1|1\rangle \; ; \; |\psi_2\rangle = \alpha_2|0\rangle + \beta_2|1\rangle$$

$$|\psi_1\psi_2\rangle = \alpha_1\alpha_2|00\rangle + \alpha_1\beta_2|01\rangle + \beta_1\alpha_2|10\rangle + \beta_1\beta_2|11\rangle$$

$$|\psi\rangle \overset{!}{=} |\psi_1\psi_2\rangle \Rightarrow \alpha_1\alpha_2 = \frac{1}{\sqrt{2}} \wedge \beta_1\beta_2 = \frac{1}{\sqrt{2}} \Rightarrow \alpha_1\beta_2 \neq 0 \wedge \alpha_2\beta_1 \neq 0 \,!$$

It is not possible! This state is special, it is entangled!

Use this property as a resource in quantum information processing:

- super dense coding
- teleportation
- error correction

2.7.5 Measurement of a single qubit in an entangled state

$$|\psi\rangle = \frac{1}{\sqrt{2}}\left(|00\rangle + |11\rangle\right)$$

measurement of first qubit:

$$p_1(0) = \langle\psi|(M_0 \otimes I)^\dagger (M_0 \otimes I)|\psi\rangle = \frac{1}{\sqrt{2}}\langle 00|\frac{1}{\sqrt{2}}|00\rangle = \frac{1}{2}$$

post measurement state:

$$|\psi'\rangle = \frac{(M_0 \otimes I)|\psi\rangle}{\sqrt{p_1(0)}} = \frac{\frac{1}{\sqrt{2}}|00\rangle}{\frac{1}{\sqrt{2}}} = |00\rangle$$

measurement of qubit two will then result with certainty in the same result:

$$p_2(0) = \langle\psi'|(I \otimes M_0)^\dagger (I \otimes M_0)|\psi'\rangle = 1$$

The two measurement results are **correlated**!

- Correlations in quantum systems can be stronger than correlations in classical systems.
- This can be generally proven using the **Bell inequalities** which will be discussed later.
- Make use of such correlations as a **resource** for information processing
  - super dense coding, teleportation, error corrections

2.7.6 Super Dense Coding

task: Try to transmit two bits of classical information between Alice (A) and Bob (B) using only one qubit.

- As Alice and Bob are living in a quantum world they are allowed to use one pair of entangled qubits that they have prepared ahead of time.

protocol:
A) Alice and Bob each have one qubit of an entangled pair in their possession

$$|\psi\rangle = 1/\sqrt{2}\,(|00\rangle + |11\rangle)$$

B) Alice does a quantum operation on her qubit depending on which 2 classical bits she wants to communicate
C) Alice sends her qubit to Bob
D) Bob does one measurement on the entangled pair

shared entanglement $\quad \frac{1}{\sqrt{2}}|00\rangle + \frac{1}{\sqrt{2}}|11\rangle$

local operations $\quad X_1 ; Y_1 ; Z_1 ; I_1$

send Alices qubit to Bob

Bob measures

bits to be transferred:	Alice's operation	resulting 2-qubit state	Bob's measurement
00	I_1	$I_1 \lvert\psi\rangle = \frac{1}{\sqrt{2}}\left(\lvert 00\rangle + \lvert 11\rangle\right)$	measure in Bell basis
01	Z_1	$Z_1 \lvert\psi\rangle = \frac{1}{\sqrt{2}}\left(\lvert 00\rangle - \lvert 11\rangle\right)$	
10	X_1	$X_1 \lvert\psi\rangle = \frac{1}{\sqrt{2}}\left(\lvert 10\rangle + \lvert 01\rangle\right)$	
11	iY_1	$iY_1 \lvert\psi\rangle = \frac{1}{\sqrt{2}}\left(\lvert 10\rangle - \lvert 01\rangle\right)$	

- all these states are entangled (try!)
- they are called the Bell states

comments:

- two qubits are involved in protocol BUT Alice only interacts with one and sends only one along her quantum communications channel
- two bits cannot be communicated sending a single classical bit along a classical communications channel

original proposal of super dense coding: Charles H. Bennett and Stephen J. Wiesner, Communication via one- and two-particle operators on Einstein-Podolsky-Rosen states, **Phys. Rev. Lett. 69, 2881(1992)**

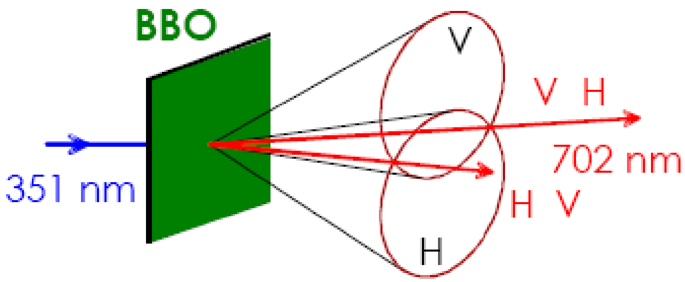
2.7.7 Experimental demonstration of super dense coding using photons

Generating polarization entangled photon pairs using **Parametric Down Conversion**:

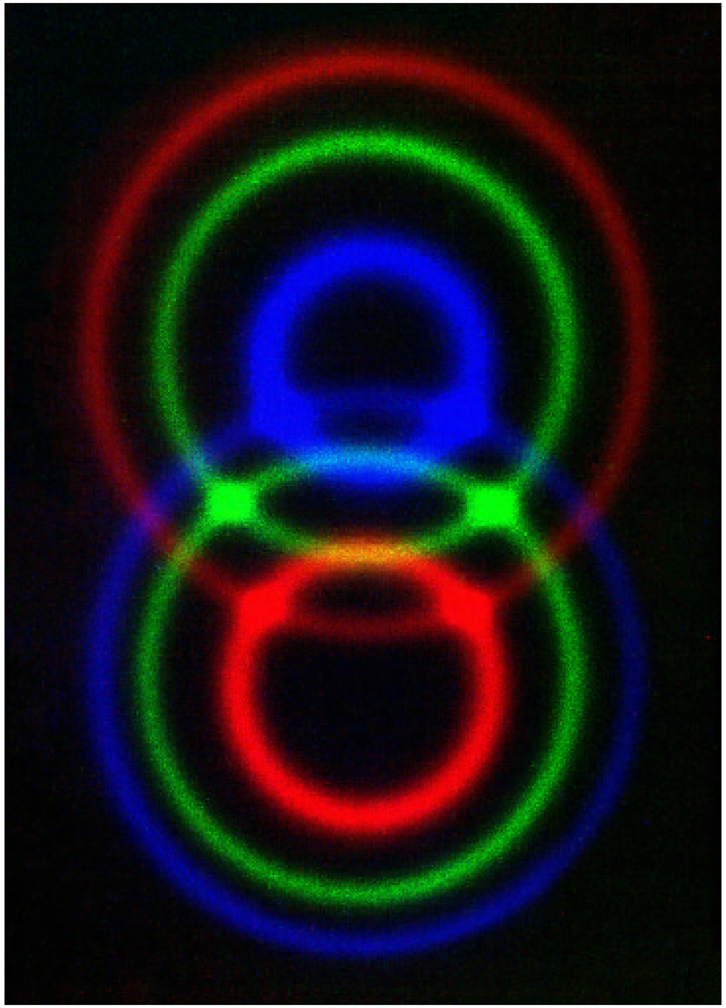
parametric down-conversion

- 1 UV-photon ➔ 2 "red" photons
- conservation of
  - energy $\omega_p = \omega_s + \omega_i$
  - momentum $\vec{k}_p = \vec{k}_s + \vec{k}_i$
- Polarisationskorrelationen (typ II)

optically nonlinear medium: BBO (BaB_2O_4) beta barium borate

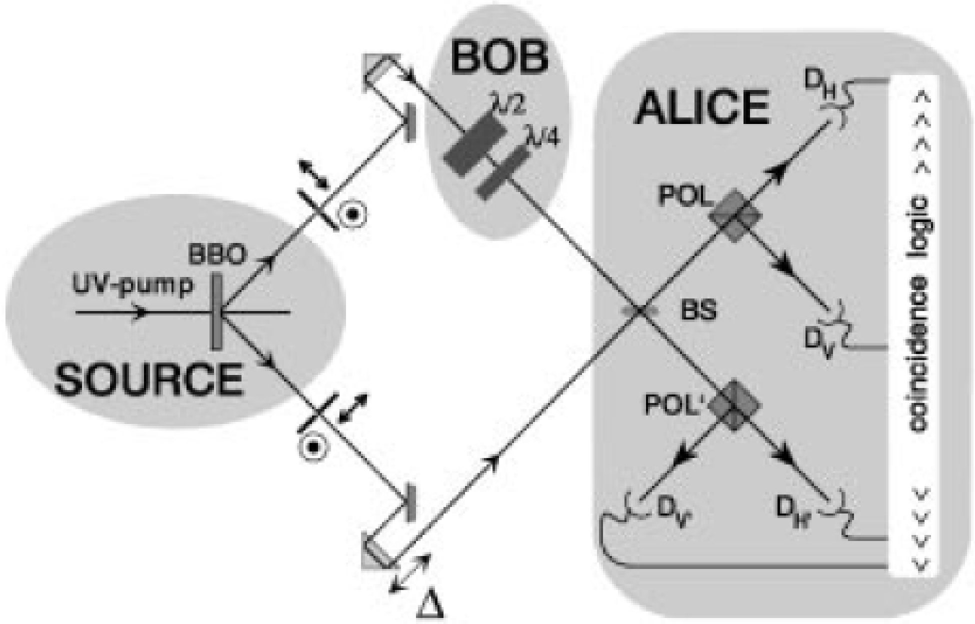


$$|\Psi^-\rangle = \frac{1}{\sqrt{2}}\left(|H\rangle|V\rangle - |V\rangle|H\rangle\right)$$



state manipulation

Bell state measurement



$$\Psi^- = \frac{1}{\sqrt{2}} (|HV\rangle - |VH\rangle) \quad \text{asym.}$$

$$\Psi^+ = \frac{1}{\sqrt{2}} (|HV\rangle + |VH\rangle)$$

$$\Phi^+ = \frac{1}{\sqrt{2}} (|HH\rangle + |VV\rangle)$$

$$\Phi^- = \frac{1}{\sqrt{2}} (|HH\rangle - |VV\rangle)$$

sym.

H = horizontal polarization
V = vertical polarization

Klaus Mattle, Harald Weinfurter, Paul G. Kwiat, and Anton Zeilinger, Dense coding in experimental quantum communication, Phys. Rev. Lett.76, 4656 (1996)

2.8 Two Qubit Quantum Logic Gates

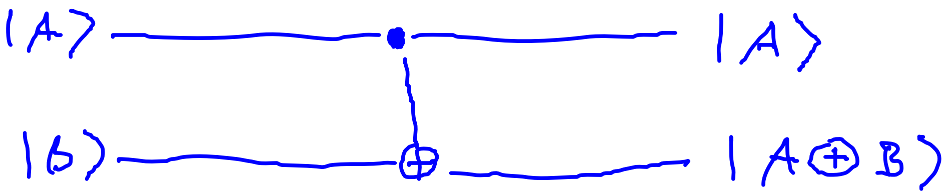
2.8.1 The controlled NOT gate (CNOT)

function:

$$|00\rangle \rightarrow |00\rangle$$
$$|01\rangle \rightarrow |01\rangle$$
$$|10\rangle \rightarrow |11\rangle$$
$$|11\rangle \rightarrow |10\rangle$$

$$|A,B\rangle \rightarrow |A, A \oplus B\rangle$$ addition mod 2 of basis states

CNOT circuit:



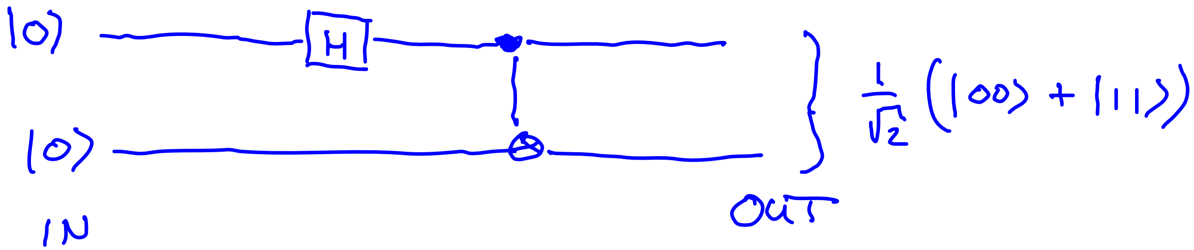
control qubit

target qubit

comparison with classical gates:
- XOR is not reversible
- CNOT is reversible (unitary)

Universality of controlled NOT:
Any multi qubit logic gate can be composed of CNOT gates and single qubit gates X,Y,Z.

2.8.2 **Application of CNOT: generation of entangled states (Bell states)**



$$|00\rangle \xrightarrow{H_1} \frac{1}{\sqrt{2}}\left(|00\rangle + |10\rangle\right) \xrightarrow{CNOT} \frac{1}{\sqrt{2}}\left(|00\rangle + |11\rangle\right)$$

$$|01\rangle \xrightarrow{H_1} \frac{1}{\sqrt{2}}\left(|01\rangle + |11\rangle\right) \xrightarrow{CNOT} \frac{1}{\sqrt{2}}\left(|01\rangle + |10\rangle\right)$$

$$|10\rangle \xrightarrow{H_1} \frac{1}{\sqrt{2}}\left(|00\rangle - |10\rangle\right) \xrightarrow{CNOT} \frac{1}{\sqrt{2}}\left(|00\rangle - |11\rangle\right)$$

$$|11\rangle \xrightarrow{H_1} \frac{1}{\sqrt{2}}\left(|01\rangle - |11\rangle\right) \xrightarrow{CNOT} \frac{1}{\sqrt{2}}\left(|01\rangle - |10\rangle\right)$$

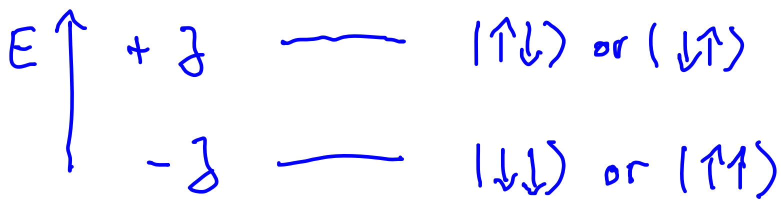
exercise: Write down the unitary matrix representations of the CNOT in the computational basis with qubit 1 being the control qubit. Write down the matrix in the same basis with qubit 2 being the control bit.

2.8.3 Implementation of CNOT using the Ising interaction

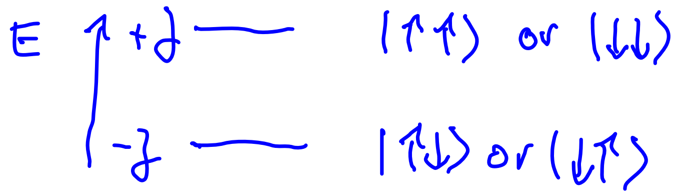
Ising interaction: $H = -\sum_{ij} J_{ij} \hat{Z}_i \hat{Z}_j$ pair wise spin interaction

generic two-qubit interaction: $H = -J \hat{Z}_1 \hat{Z}_2$

J > 0: ferromagnetic coupling



J < 0: anti-ferrom. coupling



2-qubit unitary evolution: $C(\gamma) = e^{-i \frac{\gamma}{2} \hat{Z}_1 \hat{Z}_2}$

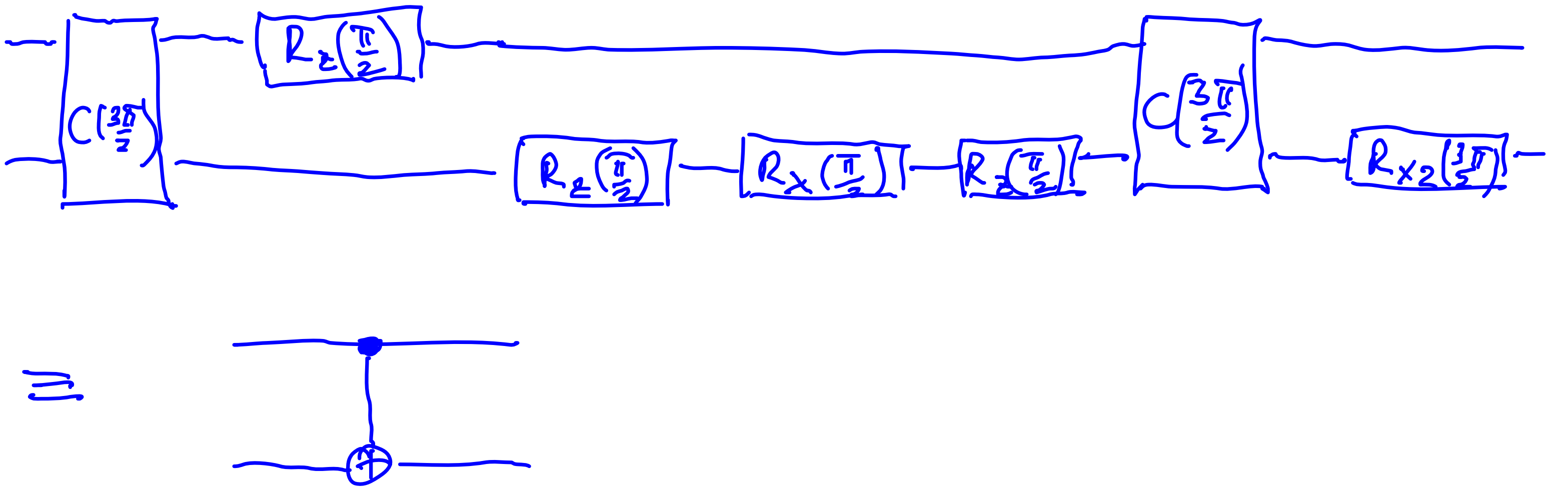
BUT this does not realize a CNOT gate yet. Additionally, single qubit operations on each of the qubits are required to realize a CNOT gate.

CNOT realization with the Ising-type interaction

CNOT - unitary:

$$C_{NOT} = e^{-i\frac{3\pi}{4}} R_{x_2}\left(\frac{3\pi}{2}\right) C\left(\frac{3\pi}{2}\right) R_{z_2}\left(\frac{\pi}{2}\right) R_{x_2}\left(\frac{\pi}{2}\right) R_{z_2}\left(\frac{\pi}{2}\right) R_{z_1}\left(\frac{\pi}{2}\right) C\left(\frac{3\pi}{2}\right)$$

circuit representation:



Any physical two-qubit interaction that can produce entanglement can be turned into a universal two-qubit gate (such as the CNOT gate) when it is augmented by arbitrary single qubit operations.

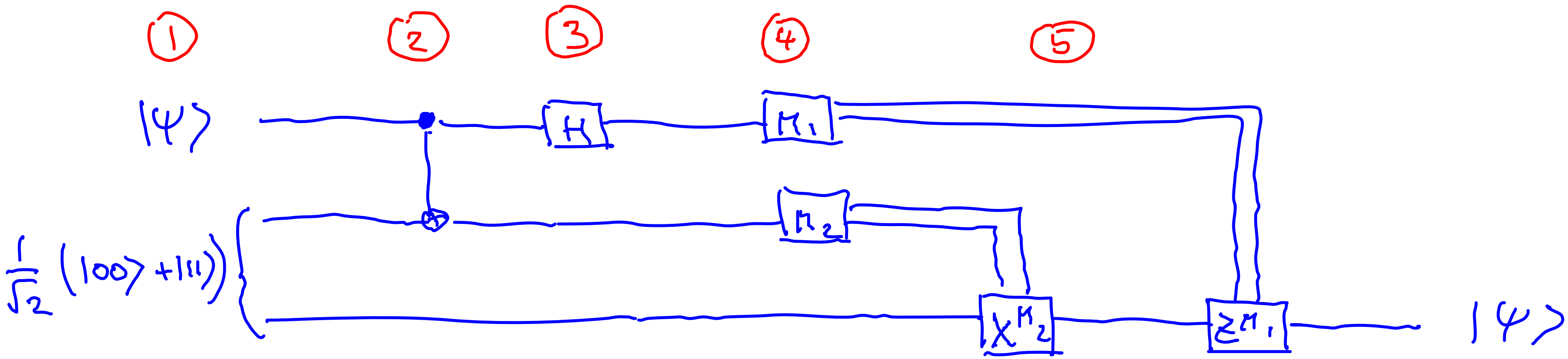
Bremner et al., *Phys. Rev. Lett.* **89**, 247902 (2002)

2.9 Quantum Teleportation

Task: Alice wants to transfer an unknown quantum state ψ to Bob only using **one entangled pair** of qubits and **classical information** as a resource.

note:
- Alice does not know the state to be transmitted
- Even if she knew it the classical amount of information that she would need to send would be infinite.

The **teleportation circuit**:



original article: Bennett, C. H. et al., Teleporting an unknown quantum state via dual classical and Einstein-Podolsky-Rosen channels, *Phys. Rev. Lett.* **70**, 1895-1899 (1993)

2.9.1 How does it work?

(1) $$|\psi\rangle \otimes \frac{1}{\sqrt{2}}\left(|00\rangle + |11\rangle\right) = \frac{1}{\sqrt{2}}\left(\alpha|000\rangle + \alpha|011\rangle + \beta|100\rangle + \beta|111\rangle\right)$$

CNOT between qubit to be teleported and one bit of the entangled pair:

(2) $$\xrightarrow{CNOT_{12}} \frac{1}{\sqrt{2}}\left(\alpha|000\rangle + \alpha|011\rangle + \beta|110\rangle + \beta|101\rangle\right)$$

Hadamard on qubit to be teleported:

(3) $$\xrightarrow{H_1} \frac{1}{2}\Big[|00\rangle(\alpha|0\rangle + \beta|1\rangle) + |10\rangle(\alpha|0\rangle - \beta|1\rangle) + |01\rangle(\alpha|1\rangle + \beta|0\rangle) + |11\rangle(\alpha|1\rangle - \beta|0\rangle)\Big]$$

measurement of qubit 1 and 2, classical information transfer and single bit manipulation on target qubit 3:

(4) $$\xrightarrow{M_1 \otimes M_2}$$

(5)

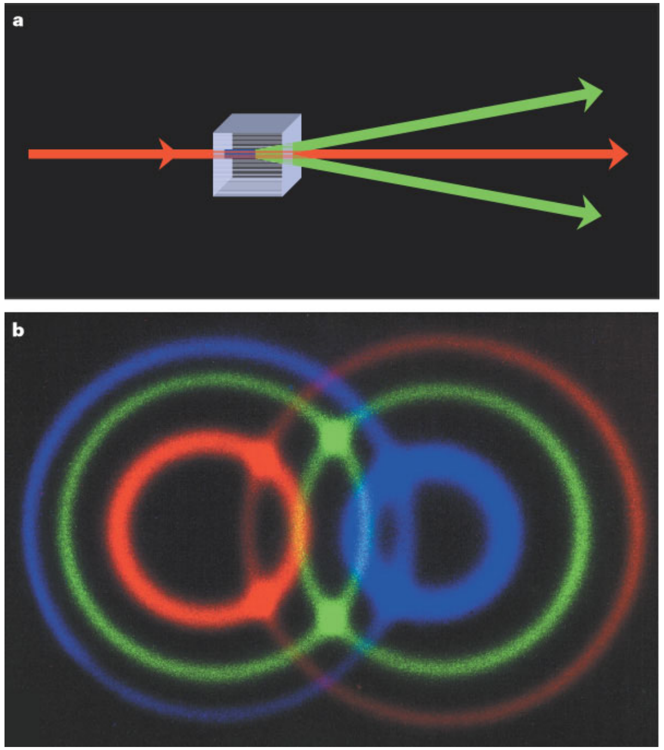
$$p_{00} = \frac{1}{4}\ ; \quad |\psi_3\rangle = \alpha|0\rangle + \beta|1\rangle \xrightarrow{I} |\psi\rangle$$

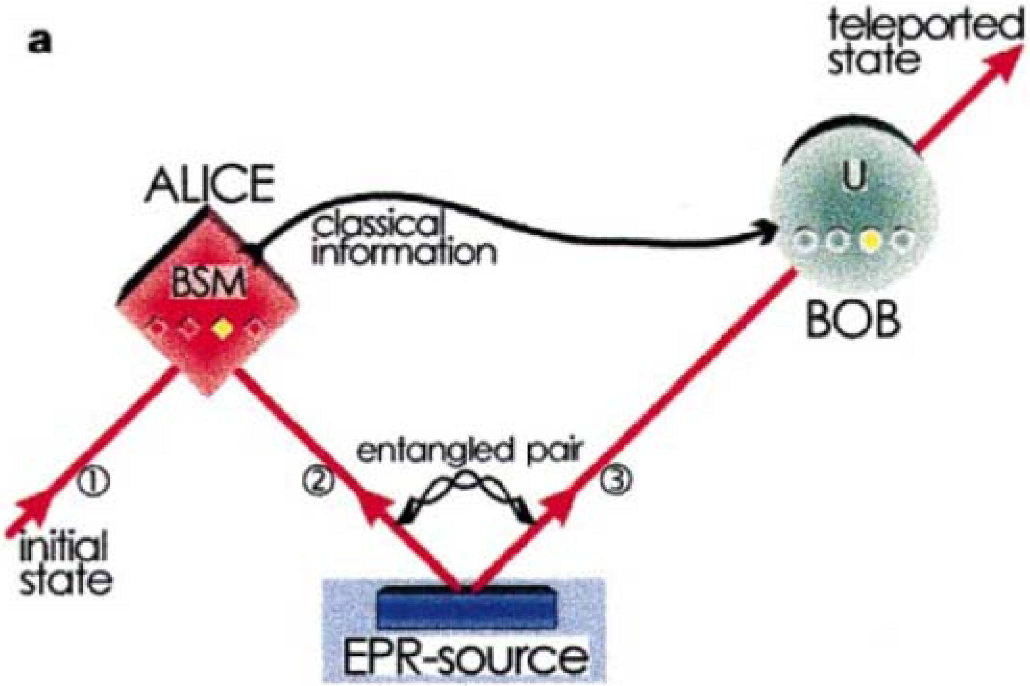
$$p_{10} = \frac{1}{4}\ ; \quad |\psi_3\rangle = \alpha|0\rangle - \beta|1\rangle \xrightarrow{Z} |\psi\rangle$$

$$p_{01} = \frac{1}{4}\ ; \quad |\psi_3\rangle = \alpha|1\rangle + \beta|0\rangle \xrightarrow{X} |\psi\rangle$$

$$p_{11} = \frac{1}{4}\ ; \quad |\psi_3\rangle = \alpha|1\rangle - \beta|0\rangle \xrightarrow{XZ} |\psi\rangle$$

2.9.2 (One) Experimental Realization of Teleportation using Photon Polarization:





- parametric down conversion (PDC) source of entangled photons
- qubits are polarization encoded

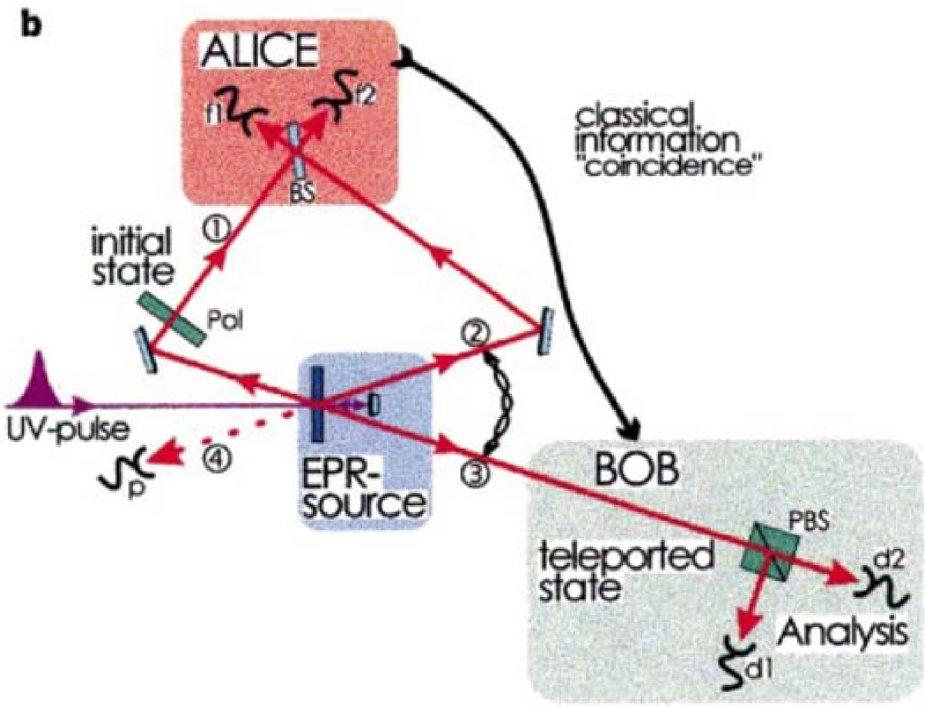
Dik Bouwmeester, Jian-Wei Pan, Klaus Mattle, Manfred Eibl, Harald Weinfurter, Anton Zeilinger, Experimental quantum teleportation *Nature* **390**, 575 (1997)

Experimental Implementation

start with states

$$|\psi_1\rangle = \alpha |H\rangle + \beta |V\rangle$$

$$|\psi_{23}\rangle = \frac{1}{\sqrt{2}} \left(|HV\rangle - |VH\rangle \right)$$



combine photon to be teleported (1) and one photon of entangled pair (2) on a 50/50 beam splitter (BS) and measure (at Alice) resulting state in Bell basis.

analyze resulting teleported state of photon (3) using polarizing beam splitters (PBS) single photon detectors

- polarizing beam splitters (PBS) as detectors of teleported states teleportation papers for you to present:

Experimental Realization of Teleporting an Unknown Pure Quantum State via Dual Classical and Einstein-Podolsky-Rosen Channels
D. Boschi, S. Branca, F. De Martini, L. Hardy, and S. Popescu
Phys. Rev. Lett. **80**, 1121 (1998) [PROLA Link]

Unconditional Quantum Teleportation
A. Furusawa, J. L. Sørensen, S. L. Braunstein, C. A. Fuchs, H. J. Kimble, and E. S. Polzik
Science 23 October 1998 282: 706-709 [DOI: 10.1126/science.282.5389.706] (in Research Articles)
Abstract » Full Text » PDF »

Complete quantum teleportation using nuclear magnetic resonance
M. A. Nielsen, E. Knill, R. Laflamme
Nature 396, 52 - 55 (05 Nov 1998) Letters to Editor
Abstract | Full Text | PDF | Rights and permissions | Save this link

Deterministic quantum teleportation of atomic qubits
M. D. Barrett, J. Chiaverini, T. Schaetz, J. Britton, W. M. Itano, J. D. Jost, E. Knill, C. Langer, D. Leibfried, R. Ozeri, D. J. Wineland
Nature 429, 737 - 739 (17 Jun 2004) Letters to Editor
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Deterministic quantum teleportation with atoms
M. Riebe, H. Haeffner, C. F. Roos, W. Haensel, J. Benhelm, G. P. T. Lancaster, T. W. Koerber, C. Becher, F. Schmidt-Kaler, D. F. V. James, R. Blatt
Nature 429, 734 - 737 (17 Jun 2004) Letters to Editor
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Quantum teleportation between light and matter
Jacob F. Sherson, Hanna Krauter, Rasmus K. Olsson, Brian Julsgaard, Klemens Hammerer, Ignacio Cirac, Eugene S. Polzik
Nature 443, 557 - 560 (05 Oct 2006) Letters to Editor
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