

Multiparticle Entanglement

Andrés Vargas Lugo Cantú

Andreas Wanner

(our names are entangled?)

les financiers

- US National Security Agency (NSA)
- Advanced Research and Development Activity (ARDA)
- Department of Defence Multidisciplinary University Research Initiative (MURI)
- NIST

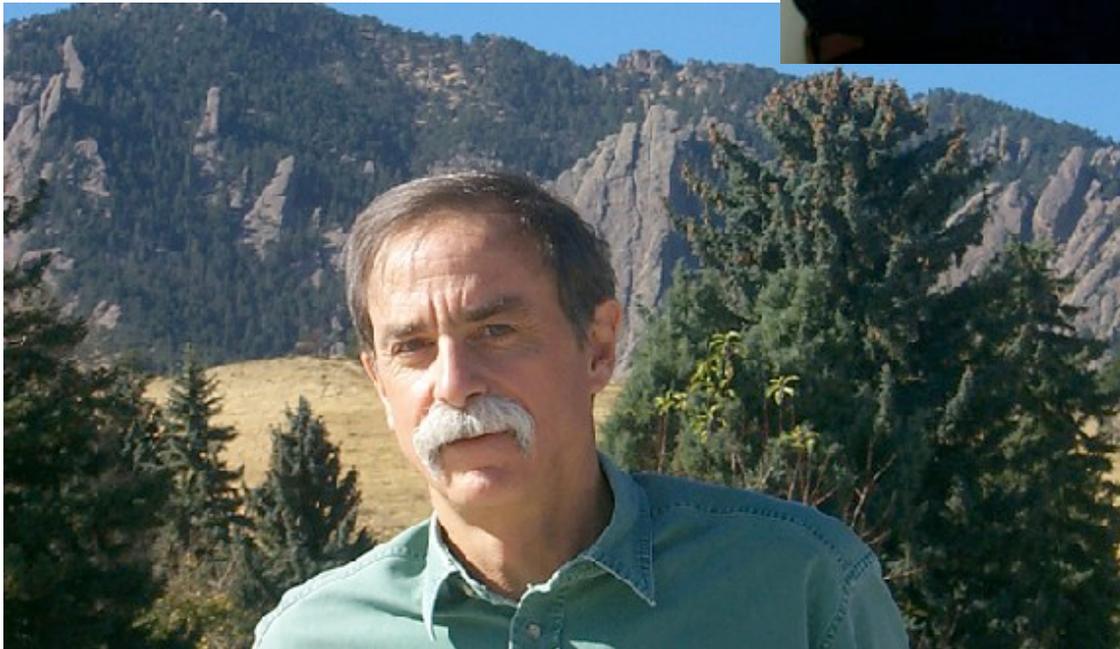
les financiers

- Austrian Science Fund (FWF)
- European Commission (AQUTE, STREP project MICROTRAP)
- IARPA, the CIFAR JFA, NSERC
- Institut für Quanteninformation GmbH
- This material is based upon work supported in part by the U. S. Army Research Office.

Entangled states of trapped atomic ions

Rainer Blatt^{1,2} & David Wineland³

Der Würdigungspreis
im Bereich
Naturwissenschaft ging
an den Innsbrucker
Physiker Rainer **Blatt**



David **Wineland**, a member
of the NIST Ion Storage
Group

14-Qubit Entanglement: Creation and Coherence

Thomas Monz,¹ Philipp Schindler,¹ Julio T. Barreiro,¹ Michael Chwalla,¹ Daniel Nigg,¹ William A. Coish,^{2,3}
Maximilian Harlander,¹ Wolfgang Hänsel,⁴ Markus Hennrich,^{1,*} and Rainer Blatt^{1,4}

¹*Institut für Experimentalphysik, Universität Innsbruck, Technikerstr. 25, A-6020 Innsbruck, Austria*

²*Institute for Quantum Computing and Department of Physics and Astronomy, University of Waterloo,
Waterloo, ON, N2L 3G1, Canada*

³*Department of Physics, McGill University, Montreal, Quebec, Canada H3A 2T8*

⁴*Institut für Quantenoptik und Quanteninformation, Österreichische Akademie der Wissenschaften,
Otto-Hittmair-Platz 1, A-6020 Innsbruck, Austria*

(Received 30 September 2010; published 31 March 2011)

We report the creation of Greenberger-Horne-Zeilinger states with up to 14 qubits. By investigating the coherence of up to 8 ions over time, we observe a decay proportional to the square of the number of qubits. The observed decay agrees with a theoretical model which assumes a system affected by correlated, Gaussian phase noise. This model holds for the majority of current experimental systems developed towards quantum computation and quantum metrology.

Creation of a six-atom 'Schrödinger cat' state

D. Leibfried¹, E. Knill¹, S. Seidelin¹, J. Britton¹, R. B. Blakestad¹, J. Chiaverini^{1,†}, D. B. Hume¹, W. M. Itano¹,
J. D. Jost¹, C. Langer¹, R. Ozeri¹, R. Reichle¹ & D. J. Wineland¹

Among the classes of highly entangled states of multiple quantum systems, the so-called 'Schrödinger cat' states are particularly useful. Cat states are equal superpositions of two maximally different quantum states. They are a fundamental resource in fault-tolerant quantum computing¹⁻³ and quantum communication, where they can enable protocols such as open-destination teleportation⁴ and secret sharing⁵. They play a role in fundamental tests of quantum mechanics⁶ and enable improved signal-to-noise

$\tilde{S}_z|\uparrow\rangle = \frac{1}{2}|\uparrow\rangle$ and $\tilde{S}_z|\downarrow\rangle = -\frac{1}{2}|\downarrow\rangle$ (for simplicity we set $\hbar = 1$). We define $|\uparrow, N\rangle \equiv |\uparrow\rangle_1|\uparrow\rangle_2\dots|\uparrow\rangle_N$ and $|\downarrow, N\rangle \equiv |\downarrow\rangle_1|\downarrow\rangle_2\dots|\downarrow\rangle_N$. In this notation, prototypical cat states of N qubits can be written as:

$$|N \text{ Cat}\rangle = \frac{1}{\sqrt{2}}(|\uparrow, N\rangle + e^{i\theta}|\downarrow, N\rangle) \quad (1)$$

To generate such states we initially prepare the ions in state $|\downarrow, N\rangle$ and then apply the following unitary operation to transform the

Cat states

- Cat states: „equal superpositions of two maximally different states“.
- For more than two qubits (subsystems):

GHZ-states (Greenberger-Horne-Zeilinger

$$|\Psi\rangle = 1/\sqrt{2} (|0\dots 0\rangle + |1\dots 1\rangle)$$

here: $|N \text{ Cat}\rangle = 1/\sqrt{2} (|\uparrow\dots\uparrow\rangle + e^{i\theta}|\downarrow\dots\downarrow\rangle)$

Ion traps



String of ions

· · · · ·
in a linear Paul trap

6 Qubits

${}^9\text{Be}^+$ ions

14 Qubits

${}^{40}\text{Ca}^+$ ions

Ingredients

- atomic ions
 - confined in electromagnetic traps
 - manipulated with laser beams -> pumping
- Centre-of-mass (COM) frequencies?
 - Axial COM f. between $\omega_{\text{COM}}/2\pi = 2.6\text{MHz}$ and 3.4MHz
 - Radial COM f. $\sim 8\text{ MHz}$
- Operations U_N using two-photon stimulated Raman transitions
 - Laser Pulses with a certain freq., duration, intensity & phase; same for all qubits.

Preparation of entangled state

- Start with $|\downarrow, N\rangle$
- Apply unitary operation:

$$U_N = (\exp[i\pi/2 J_x] \exp[i\xi\pi/2 J_z]) (\exp[i\pi/2 J_z^2]) (\exp[i\pi/2 J_x])$$

with $\xi=1$ if N is odd and $\xi=0$ otherwise

- Goal: $|N \text{ Cat}\rangle = 1/\sqrt{2} (|\uparrow \dots \uparrow\rangle + e^{i\theta} |\downarrow \dots \downarrow\rangle)$
- Measure Entanglement!

We need a measure

- Fidelity: $F = |\langle \Psi_N | N \text{ Cat} \rangle|^2$

- Useful:

$$F = \frac{1}{2}(P_{\uparrow N} + P_{\downarrow N}) + |C_{\downarrow N; \uparrow N}|$$

- That is not all there is:

for $N > 2$ there is no single measure that quantifies entanglement!

- Only comparable if $|\Phi\rangle \xrightarrow{\text{LOCC}} |\Psi\rangle$

Measure (cont.)

- Witness Operator: $W = 1 - 2|N \text{ Cat}\rangle\langle N \text{ Cat}|$
- $\langle W \rangle = 1 - 2 * \text{Fidelity}$
- If $\langle W \rangle < 0$ significantly \Rightarrow entanglement
- then states can be purified by LOCC

Measure (cont.)

- ‚depolarization‘ method: Using LOCC, density matrix gets transformed, then:

N-Particle entanglement if:

$$2 |C_{\downarrow N; \uparrow N}| > \max_j (P_j + P_{j'})$$

- The most important information resides in the magnitude of coherence $C_{\downarrow N; \uparrow N}$

Results

From amplitude of parity-oscillations:

$$|C_{\downarrow 4; \uparrow 4}| \geq 0.349(2) \quad \checkmark$$

$$|C_{\downarrow 5; \uparrow 5}| \geq 0.264(2) \quad \checkmark$$

$$|C_{\downarrow 6; \uparrow 6}| \geq 0.210(2) \quad \times$$

from poissonian fits:

$$F_{4\text{Cat}} \geq 0.76(1) \quad \checkmark$$

$$F_{5\text{Cat}} \geq 0.60(2) \quad \checkmark$$

$$F_{6\text{Cat}} \geq 0.509(4)$$

$$\langle W_4 \rangle \leq 0.51(2) \quad \checkmark$$

$$\langle W_5 \rangle \leq 0.20(2) \quad \checkmark$$

$$\langle W_6 \rangle \leq 0.018(8) \quad \sim$$

Results

$$\langle W_6 \rangle \leq 0.018(8) \sim$$

Use 'depolarization' method:

$$2 |C_{\downarrow N; \uparrow N}| > \max_j (P_j + P_{j'})$$

knowing: $\max_j (P_j + P_{j'}) \leq 2 \max(P_j)$

For the |6 Cat \rangle state:

$$\begin{aligned} |C_{\downarrow 6; \uparrow 6}| &\geq 0.210(2) \\ &\geq \max_j (P_{\downarrow j} \mid j \in \{1, 2, 3, 4, 5\}) = 0.119(9) \end{aligned}$$



Take Home Message

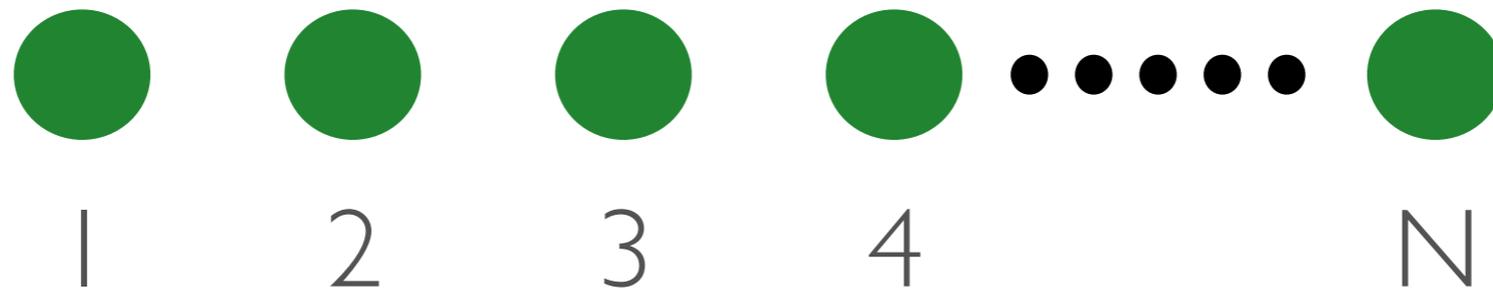
- Entanglement of N particles:
(theoretical) problem to quantify if $N > 2$

DECOHERENCE IN MULTIPARTICLE SYSTEMS

Andres Vargas
Andreas Wanner

EXPERIMENT DESCRIPTION

N $^{40}\text{Ca}^+$ Ions



$$D_{5/2}(m = -1/2) \equiv |0\rangle$$

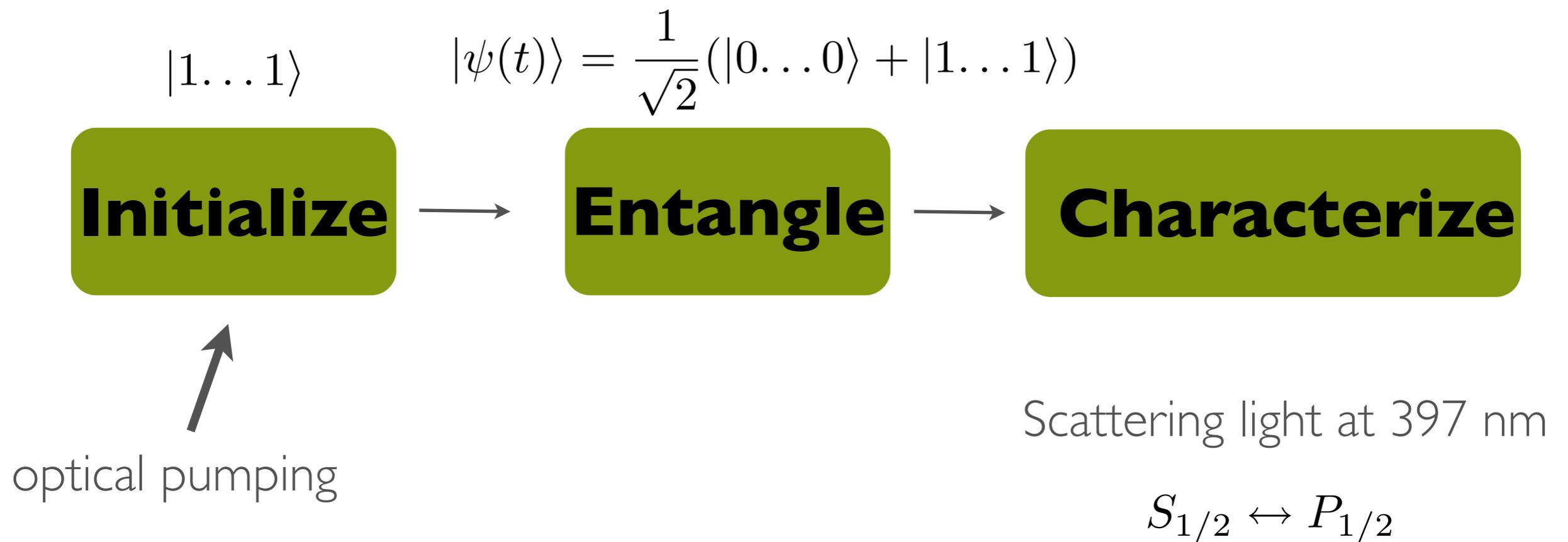
$$S_{1/2}(m = -1/2) \equiv |1\rangle$$

$$|\tilde{0}\rangle = |0\dots 0\rangle$$

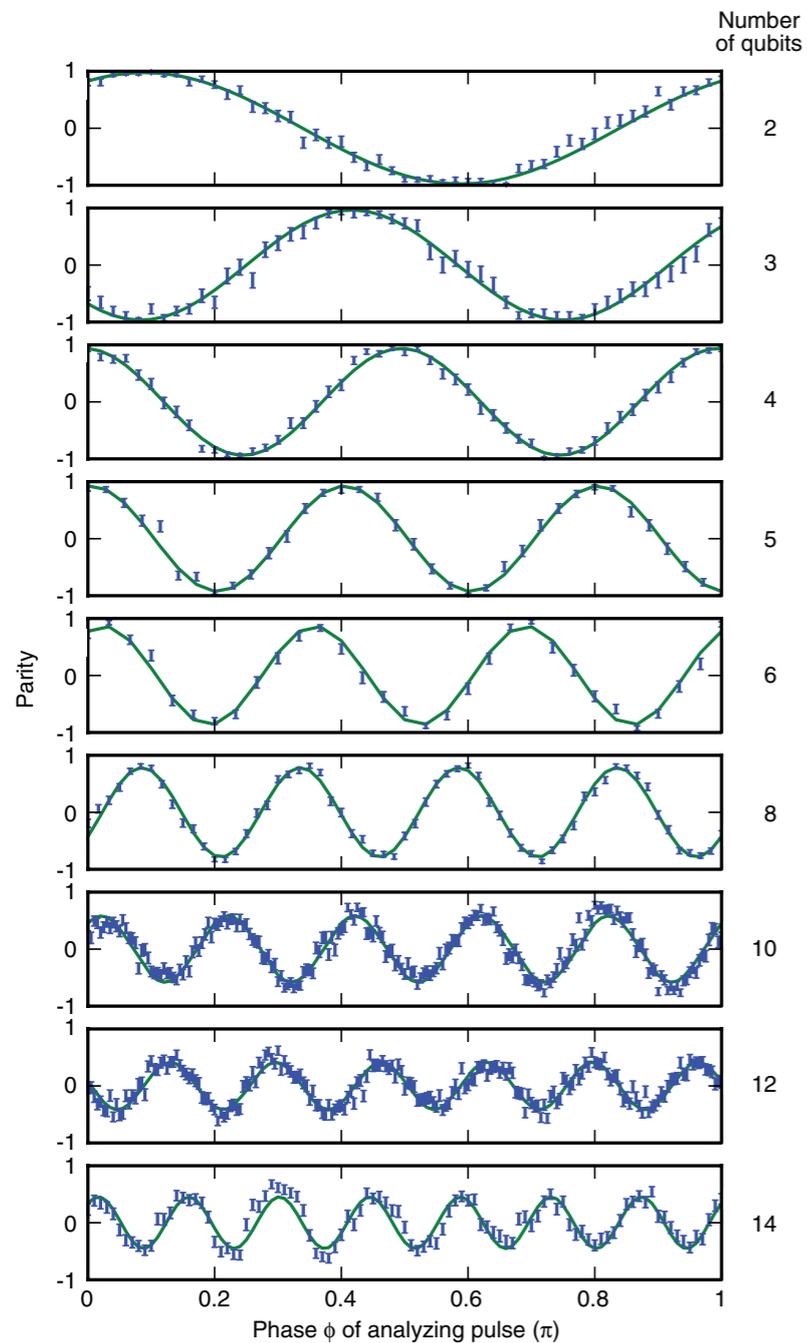
$$|\tilde{1}\rangle = |1\dots 1\rangle$$

$$|\psi(t)\rangle = \frac{1}{\sqrt{2}}(|0\dots 0\rangle + |1\dots 1\rangle)$$

EXPERIMENT DESCRIPTION



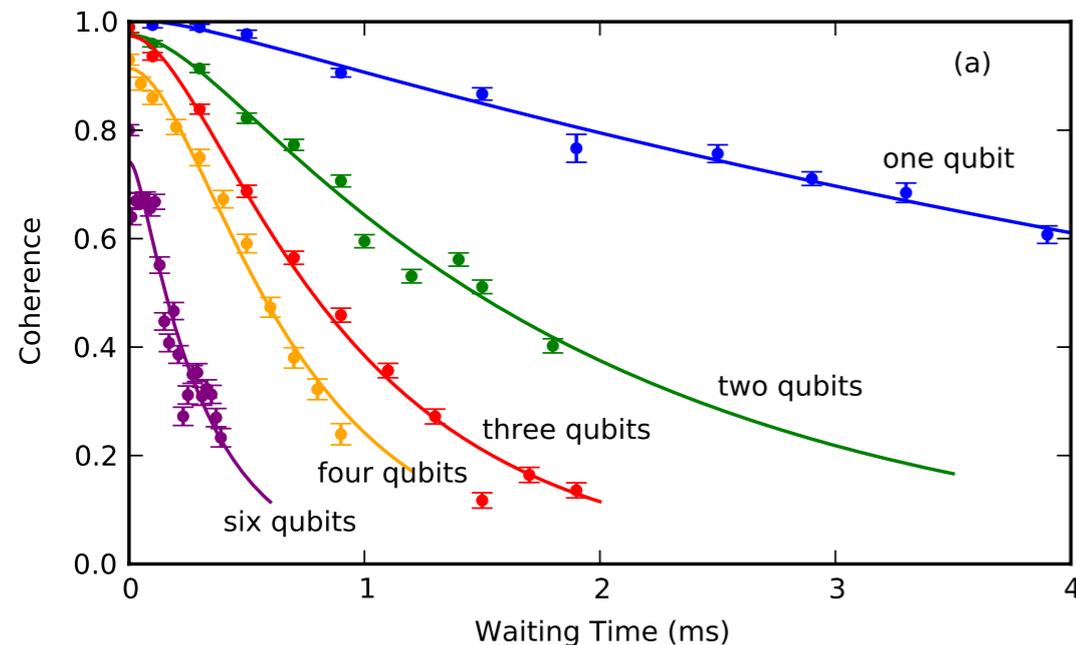
COHERENCE AS A FUNCTION OF QBIT NUMBER



After the GHZ state is generated the qubit are rotated by

The amplitude of the oscillations is the coherence

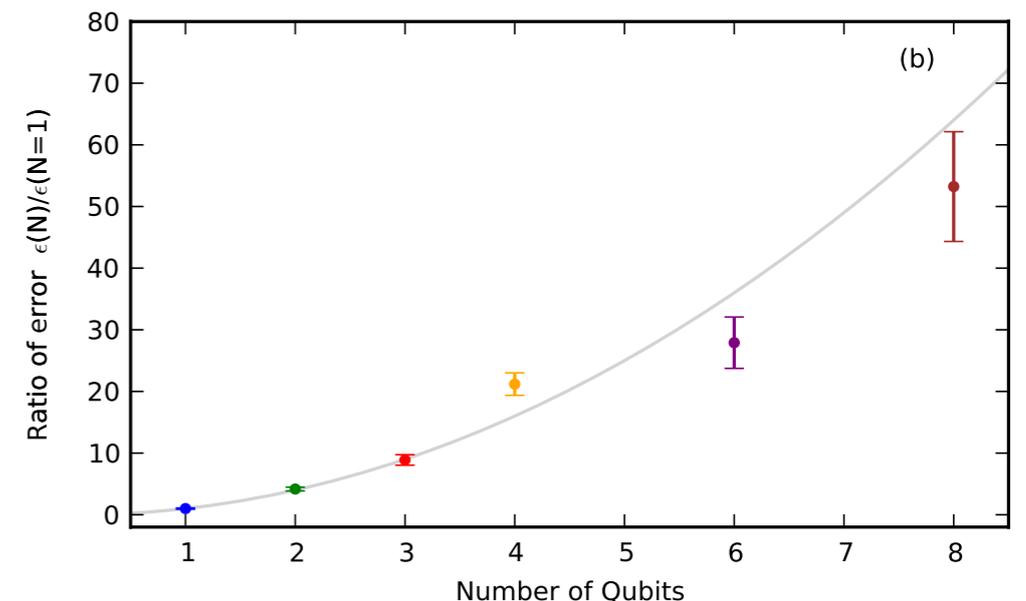
DECOHERENCE AS A FUNCTION OF TIME



A time delay is introduced between characterization and creation of state

Coherence time decreases and error increases as number of Q-bit increases

$$\epsilon(N) = N^2$$



MOST RELEVANT RESULTS

For 1 Qbit with a coherence time of 95 ms is found

When using a base

$$|00001111\rangle + |11110000\rangle$$

a 324 ms coherence time is
obtained, why??

THE MODEL FOR A SINGLE QBIT

$$H = \frac{1}{2}\sigma_z\omega_0 + \sum_k b_k^\dagger b_k + \sum_k \sigma_z (g_k b_k^\dagger + g_k^* b_k)$$

$$U(t) = \exp\left\{\sigma_z \frac{1}{2} \sum_k (b_k^\dagger \xi_k(t) - b_k \xi_k^*(t))\right\}$$

$$\sigma_z |0\rangle = -1|0\rangle$$

$$\sigma_z |1\rangle = 1|1\rangle$$

$$\xi_k(t) = 2g_k \frac{1 - e^{i\omega_k t}}{\omega_k}$$

No energy exchange or spin flip

Displacement operator in
Quantum optics

$$D(\alpha) = \exp\{\alpha a^\dagger - \alpha^* a\}$$

$$D(\alpha)|0\rangle = |\alpha\rangle$$

$$D(1/2\sigma_z\xi_k) = \exp\{1/2\sigma_z\xi_k b^\dagger - 1/2\sigma_z\xi_k^* b\}$$

$$U(t)|0\rangle \otimes |\Psi\rangle = \Pi_k D(-1/2\xi_k(t))$$

$$|\Psi\rangle = (c_0|0\rangle + c_1|1\rangle) \otimes |0_k\rangle \quad \text{Entanglement}$$

$$U(t)|\Psi\rangle = c_0|0\rangle| - 1/2\xi_k(t)\rangle + c_1|1\rangle| + 1/2\xi_k(t)\rangle$$

Generalization to 2 Qbits

$$H_{int} = \sum_k \sigma_z^a (g_k^a b_k^\dagger + g_k^{*a} b_k) + \sigma_z^b (g_k^b b_k^\dagger + g_k^{*b} b_k)$$

$$U(t) = D(1/2\sigma_z^a \xi_k^a(t) + 1/2\sigma_z^b \xi_k^b(t))$$

$$|\Phi^{(-)}\rangle = (c_{10}|1_a, 0_b\rangle + c_{01}|0_a, 1_b\rangle) \otimes |0_k\rangle$$

$$|\Phi^{(+)}\rangle = (c_{00}|0_a, 0_b\rangle + c_{11}|1_a, 1_b\rangle) \otimes |0_k\rangle$$

States with 2 different Qbit configuration couple differently to the field

$$U(t)|\Phi^{(+)}\rangle = (c_{00}|0_a, 0_b\rangle| - 1/2(\xi_k^a + \xi_k^b)\rangle + c_{11}|1_a, 1_b\rangle|1/2(\xi_k^a + \xi_k^b)\rangle$$

$$U(t)|\Phi^{(-)}\rangle = (c_{10}|1_a, 0_b\rangle| + 1/2(\xi_k^a - \xi_k^b)\rangle + c_{01}|0_a, 1_b\rangle| - 1/2(\xi_k^a - \xi_k^b)\rangle$$

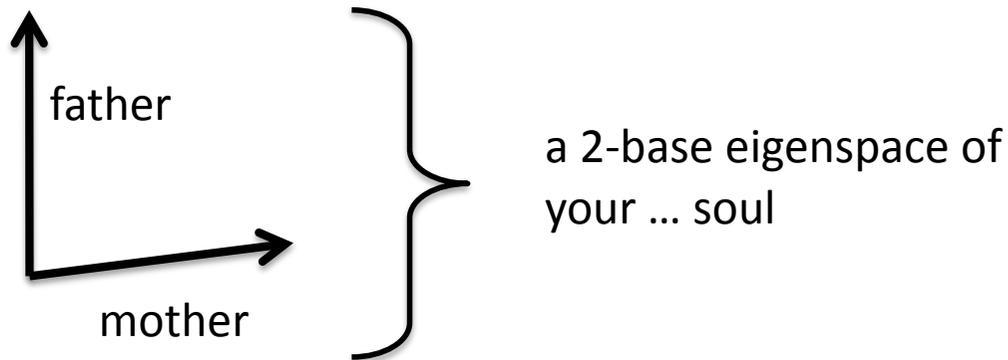
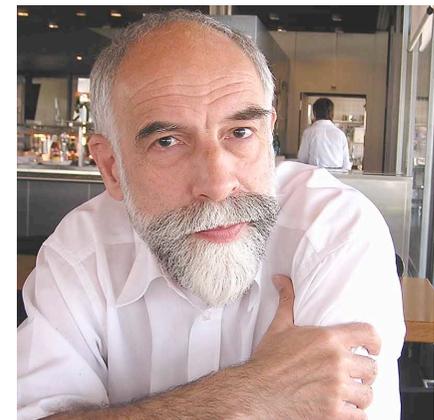
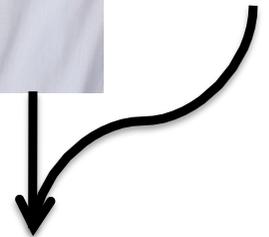
If $\xi_k^a = \xi_k^b$ $|\Phi^{-}\rangle$ Doesn't couple to the field

general *pros* and *cons* regarding
QC with iontraps

- electron and nuclear spins (Spin $\frac{1}{2}$ particles)
inherently only have two states (*good!*)
- *problem:*
the center of mass oscillations (phonons)
have short coherence time
- *possible Solution:* strong interaction through
chemical bonds -> NMR

Thanks go out to...

Anna Wanner



soul \cong $\frac{5}{8}$ (genome \otimes culture \otimes gods further ingredients)